

An Asymptotic Model of Electroporation-Mediated Molecular Delivery in Skeletal Muscle Tissue

Jonathan Preston Cranford

CASIS Workshop

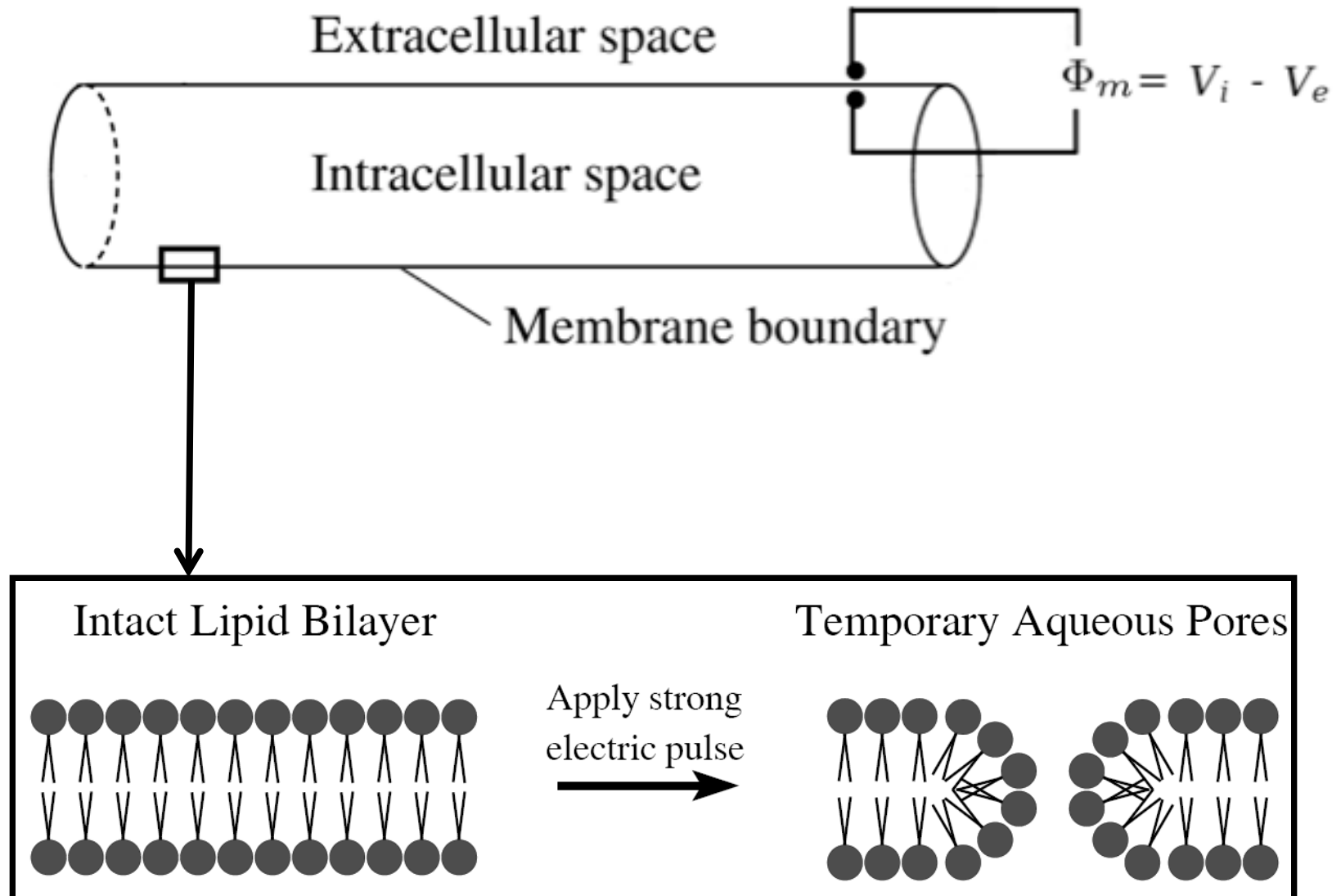
Livermore, CA

Wednesday May 18, 2016

The statements and opinions expressed herein are the expression of the authors, who are solely responsible for its content.

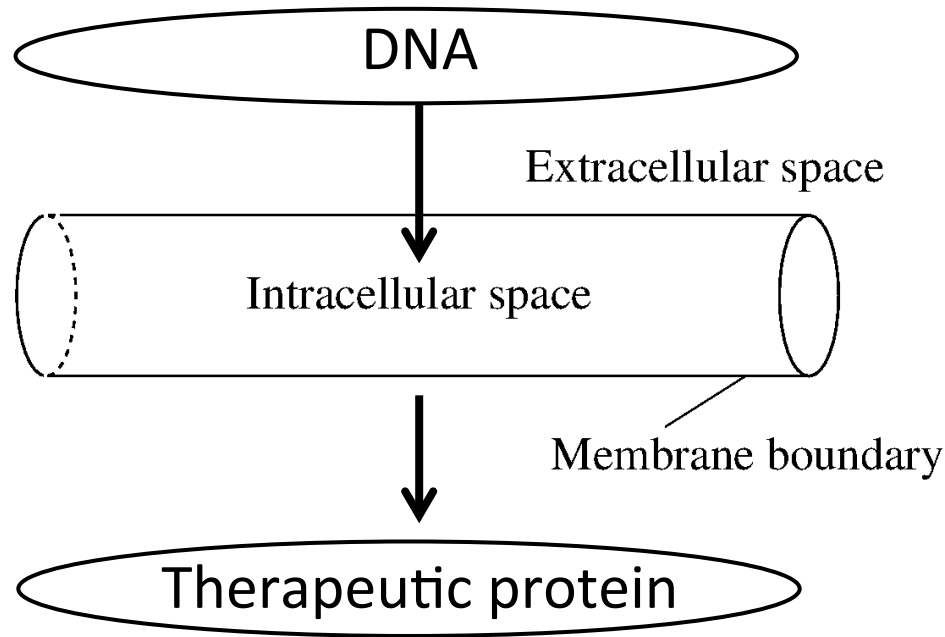
Copyright © 2016 by Jonathan P. Cranford, all rights reserved

Concept of Electroporation (EP)



Clinical Significance and Experiments

Gene therapy



EP enhances uptake and expression several orders of mag.!

Challenge: mechanism not well understood

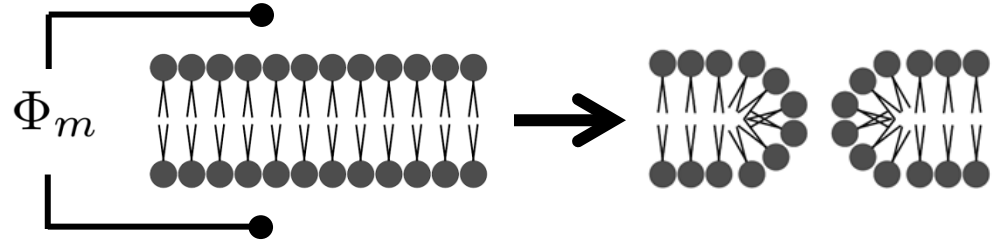
- EP depends on type of tissue and molecule
- Sensitive to pulsing protocol

Theoretical Studies

Four components

1. Membrane level

Creation of pores in response
to Φ_m

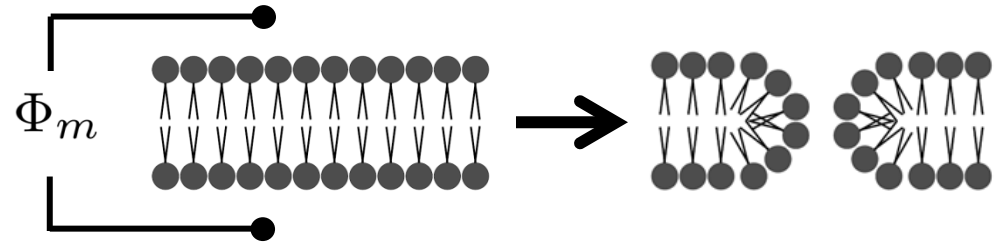


Theoretical Studies

Four components

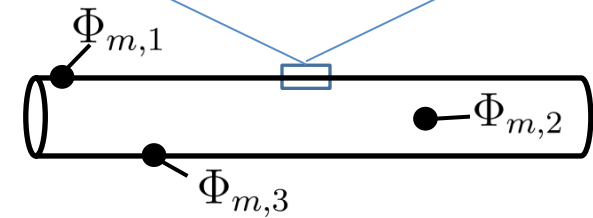
1. Membrane level

Creation of pores in response to Φ_m



2. Cellular level

Distribution of Φ_m due to microscale geometry of individual cells

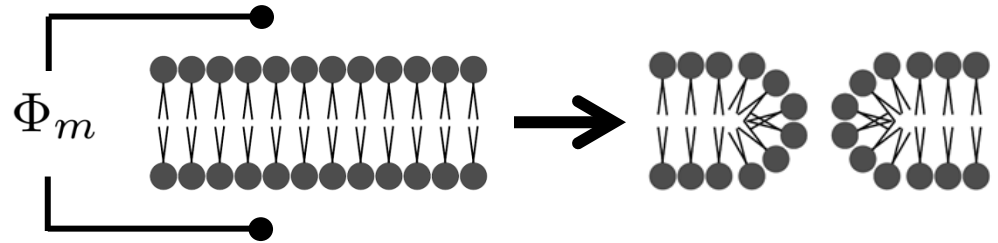


Theoretical Studies

Four components

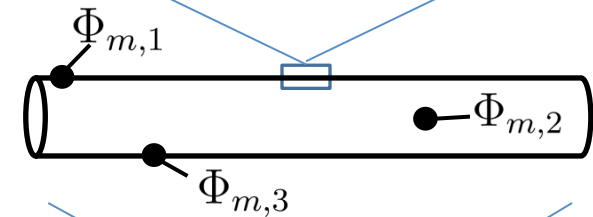
1. Membrane level

Creation of pores in response to Φ_m



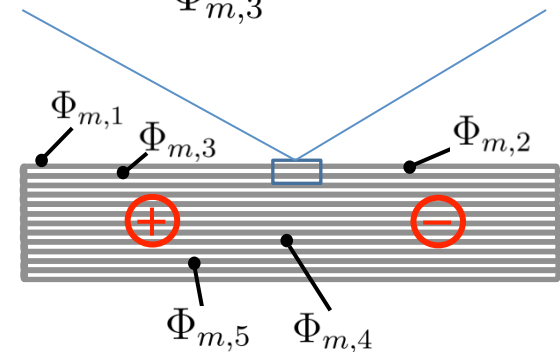
2. Cellular level

Distribution of Φ_m due to microscale geometry of individual cells



3. Tissue level

Connects macroscale tissue distribution of potential from electric stimulus to Φ_m at cellular level

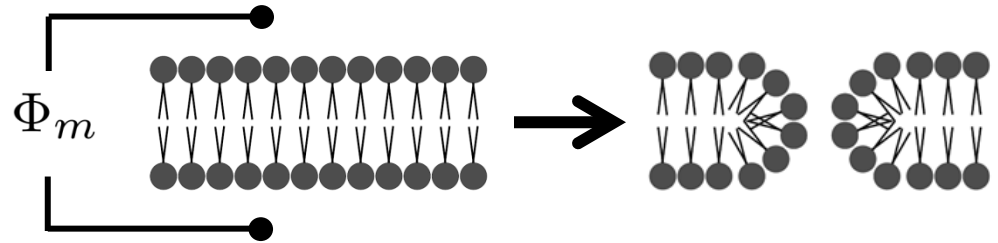


Theoretical Studies

Four components

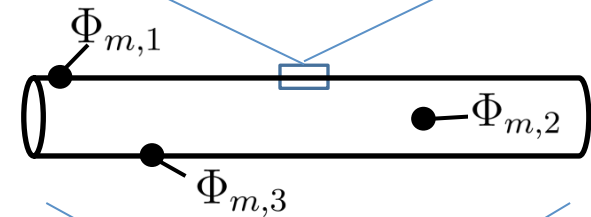
1. Membrane level

Creation of pores in response to Φ_m



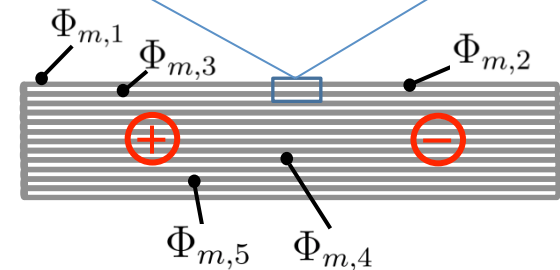
2. Cellular level

Distribution of Φ_m due to microscale geometry of individual cells



3. Tissue level

Connects macroscale tissue distribution of potential from electric stimulus to Φ_m at cellular level



4. Mass transport

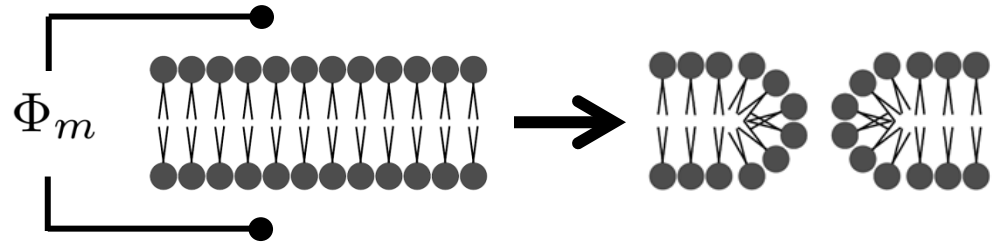


Theoretical Studies

Four components

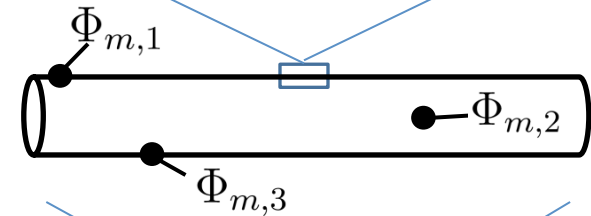
1. Membrane level

Creation of pores in response to Φ_m



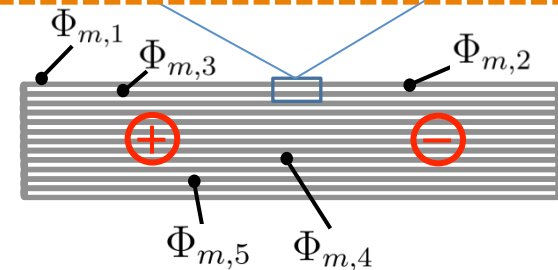
2. Cellular level

Distribution of Φ_m due to microscale geometry of individual cells



3. Tissue level

Connects macroscale tissue distribution of potential from electric stimulus to Φ_m at cellular level



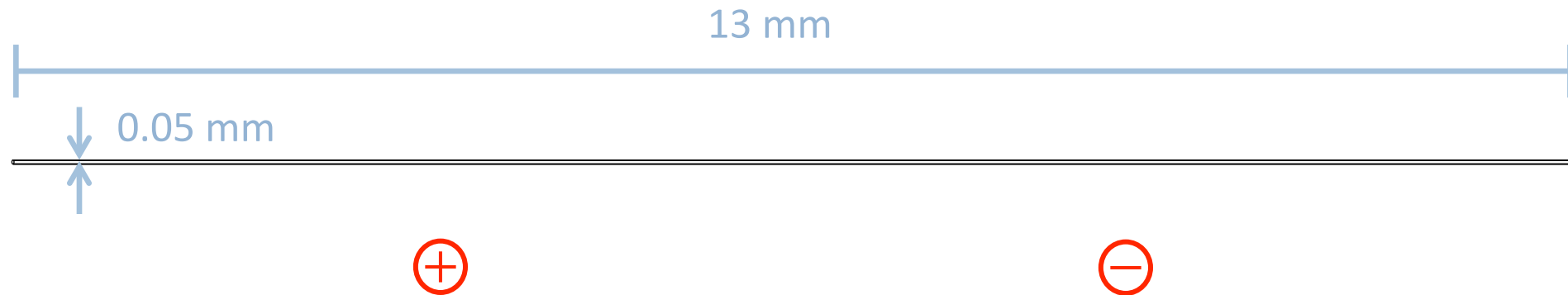
4. Mass transport



Theoretical Studies

Cellular level

Distribution of Φ_m due to microscale geometry of individual fibers



Transverse direction



Neuron modeling community:

Longitudinal component is
indispensable!

Electroporation modeling community:

Longitudinal component is neglected

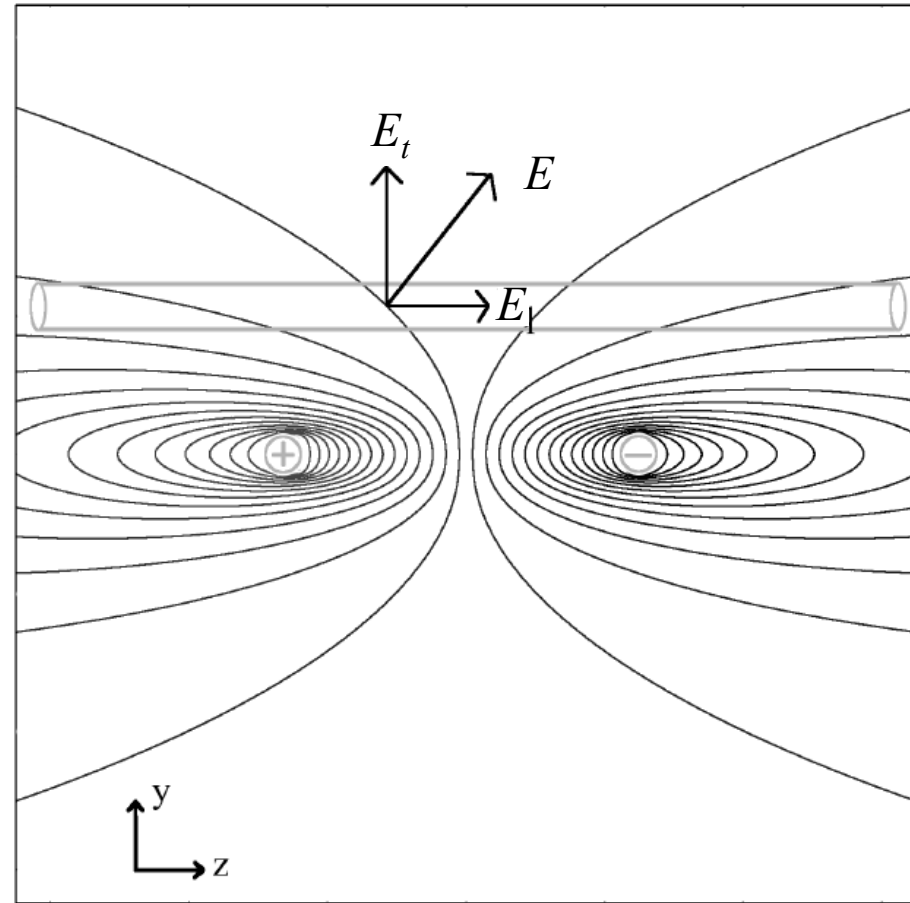
Theoretical Studies

Cellular level

Distribution of Φ_m due to microscale geometry of individual cells

EP community models

- $|\Phi_m| \propto |E|$



Theoretical Studies

Cellular level

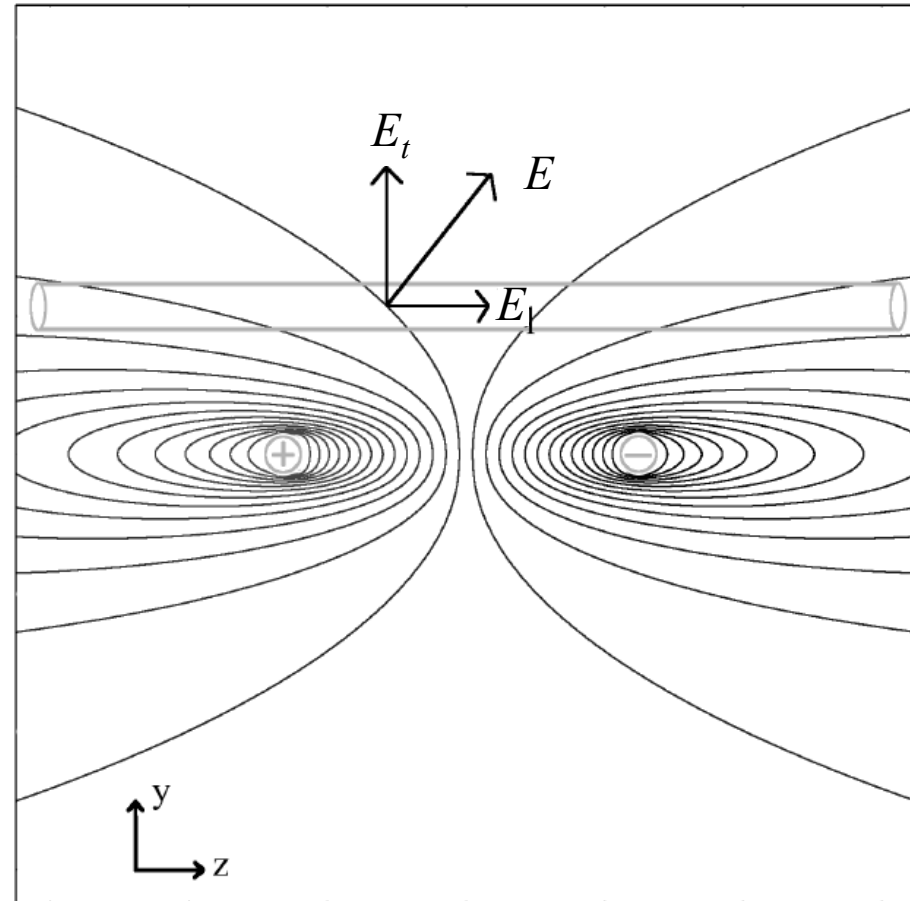
Distribution of Φ_m due to microscale geometry of individual cells

EP community models

- $|\Phi_m| \propto |E|$

Neuron community models

- “transverse charging”
 $|\Phi_m| \propto |E_t|$
- “longitudinal charging”
 $|\Phi_m| \propto \left| \frac{\partial E_l}{\partial z} \right|$
- $\left| \frac{\partial E_l}{\partial z} \right|$ varies by 7 orders of mag.



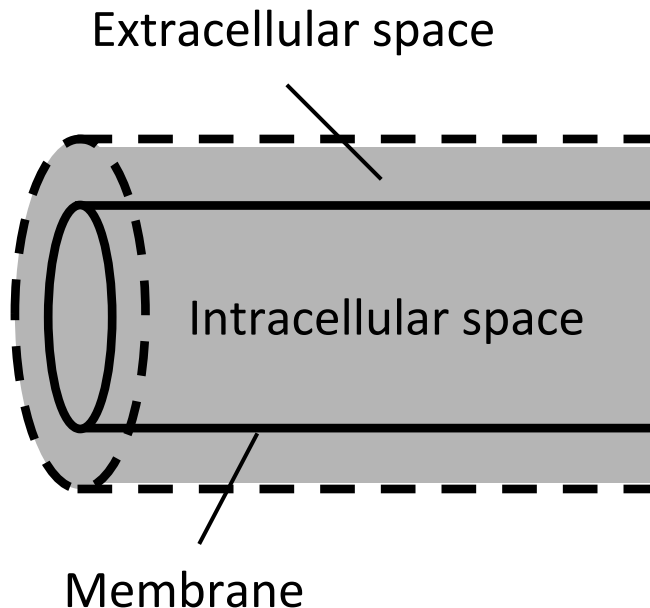
Effect of longitudinal charging on EP not characterized

Derivation of Asymptotic Fiber Model

Purpose: Compute distribution of electric potential in fiber

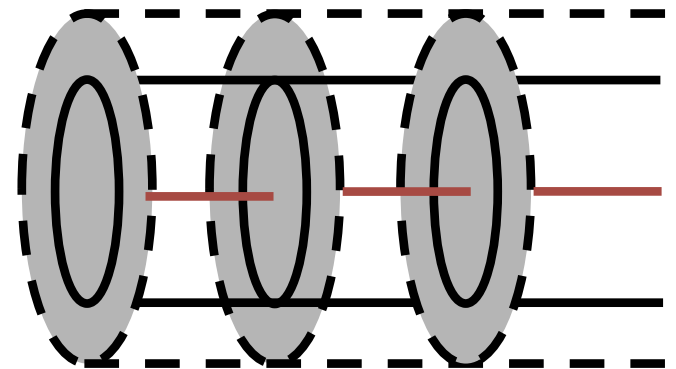
Full Fiber Model

(3D BVP)



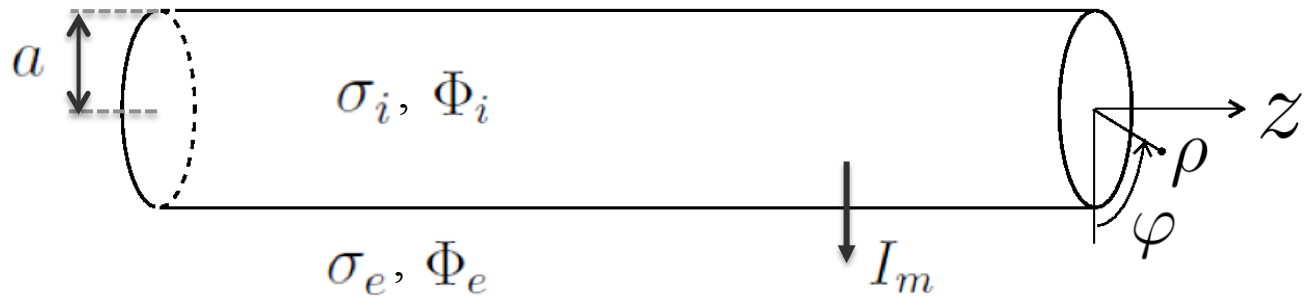
Asymptotic Fiber Model¹

(series of 2D transverse BVPs connected by 1D longitudinal problem)

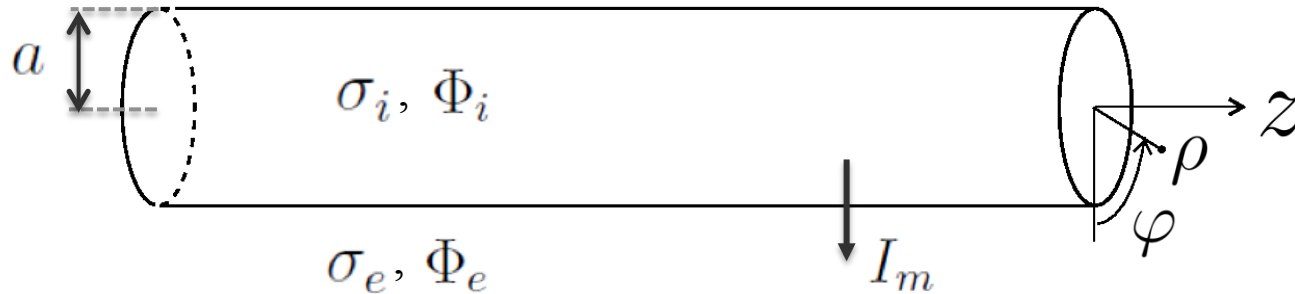


¹Cranford, J.P., Kim, B.J., Neu, W.K. (2012) MBEC, 50, 243-251

Derivation of Asymptotic Fiber Model



Derivation of Asymptotic Fiber Model



Scaled governing equations in dimensionless form

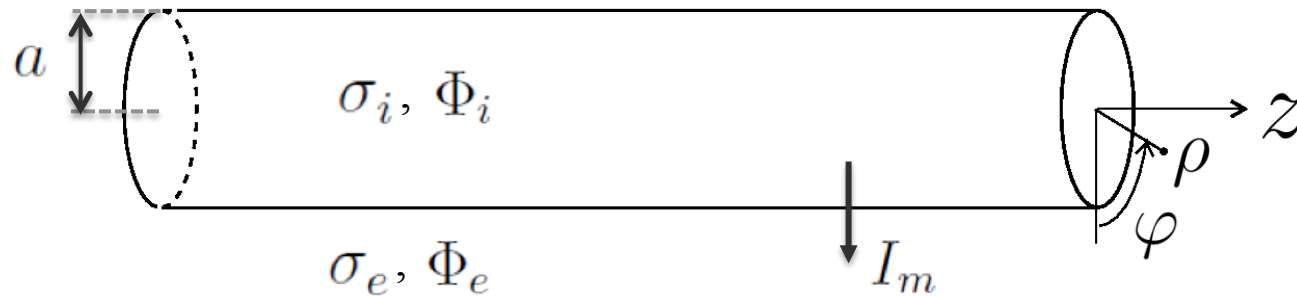
$$\frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial \Phi_i}{\partial \rho} \right) + \frac{1}{\rho^2} \frac{\partial^2 \Phi_i}{\partial \varphi^2} + \epsilon \frac{\partial^2 \Phi_i}{\partial z^2} = 0 \quad \text{in } \rho < a,$$

$$\frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial \Phi_e}{\partial \rho} \right) + \frac{1}{\rho^2} \frac{\partial^2 \Phi_e}{\partial \varphi^2} + \epsilon \frac{\partial^2 \Phi_e}{\partial z^2} = 0 \quad \text{in } \rho > a,$$

$$\frac{\partial \Phi_i}{\partial \rho} = -\frac{\partial \Phi_m}{\partial \tau} - \epsilon \frac{\partial \Phi_m}{\partial T_s} - I_m - \frac{\partial \Psi}{\partial \rho} \quad \text{on } \rho = a$$

$$\frac{\partial \Phi_e}{\partial \rho} = -\frac{\sigma_i}{\sigma_e} \left\{ \frac{\partial \Phi_m}{\partial \tau} + \epsilon \frac{\partial \Phi_m}{\partial T_s} + I_m \right\} - \frac{\partial \Psi}{\partial \rho} \quad \text{on } \rho = a$$

Derivation of Asymptotic Fiber Model



Scaled governing equations in dimensionless form

$$\frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial \Phi_i}{\partial \rho} \right) + \frac{1}{\rho^2} \frac{\partial^2 \Phi_i}{\partial \varphi^2} + \epsilon \frac{\partial^2 \Phi_i}{\partial z^2} = 0 \quad \text{in } \rho < a,$$

$$\frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial \Phi_e}{\partial \rho} \right) + \frac{1}{\rho^2} \frac{\partial^2 \Phi_e}{\partial \varphi^2} + \epsilon \frac{\partial^2 \Phi_e}{\partial z^2} = 0 \quad \text{in } \rho > a,$$

$$\frac{\partial \Phi_i}{\partial \rho} = -\frac{\partial \Phi_m}{\partial \tau} - \epsilon \frac{\partial \Phi_m}{\partial T_s} - I_m - \frac{\partial \Psi}{\partial \rho} \quad \text{on } \rho = a$$

$$\frac{\partial \Phi_e}{\partial \rho} = -\frac{\sigma_i}{\sigma_e} \left\{ \frac{\partial \Phi_m}{\partial \tau} + \epsilon \frac{\partial \Phi_m}{\partial T_s} + I_m \right\} - \frac{\partial \Psi}{\partial \rho} \quad \text{on } \rho = a$$

Derivation of Asymptotic Fiber Model

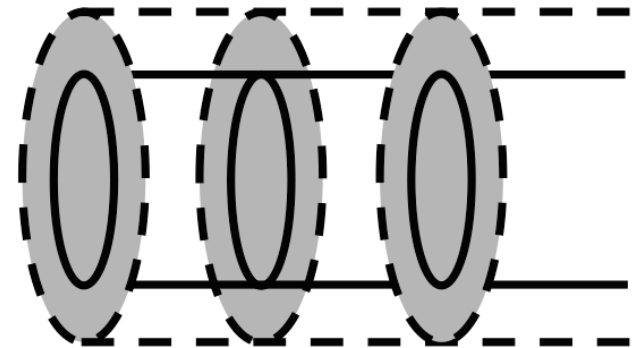
Mean-free equations of the transverse problem

$$\frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial \phi_i^0}{\partial \rho} \right) + \frac{1}{\rho^2} \frac{\partial^2 \phi_i^0}{\partial \varphi^2} = 0 \quad \text{in } \rho < a$$

$$\frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial \phi_e^0}{\partial \rho} \right) + \frac{1}{\rho^2} \frac{\partial^2 \phi_e^0}{\partial \varphi^2} = 0 \quad \text{in } \rho > a$$

$$\sigma_i \frac{\partial \phi_i^0}{\partial \rho} = -C_m \frac{\partial \phi_m^0}{\partial t} - \widetilde{I}_m - \sigma_i \frac{\partial \widetilde{\Psi}}{\partial \rho} \quad \text{on } \rho = a$$

$$\sigma_e \frac{\partial \phi_e^0}{\partial \rho} = -C_m \frac{\partial \phi_m^0}{\partial t} - \widetilde{I}_m - \sigma_e \frac{\partial \widetilde{\Psi}}{\partial \rho} \quad \text{on } \rho = a$$



Derivation of Asymptotic Fiber Model

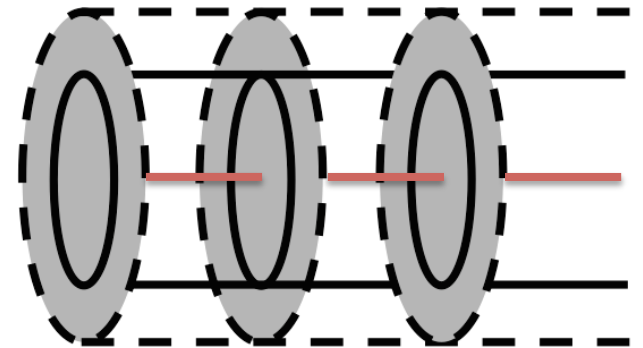
Mean-free equations of the transverse problem

$$\frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial \phi_i^0}{\partial \rho} \right) + \frac{1}{\rho^2} \frac{\partial^2 \phi_i^0}{\partial \varphi^2} = 0 \quad \text{in } \rho < a$$

$$\frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial \phi_e^0}{\partial \rho} \right) + \frac{1}{\rho^2} \frac{\partial^2 \phi_e^0}{\partial \varphi^2} = 0 \quad \text{in } \rho > a$$

$$\sigma_i \frac{\partial \phi_i^0}{\partial \rho} = -C_m \frac{\partial \phi_m^0}{\partial t} - \widetilde{I}_m - \sigma_i \frac{\partial \tilde{\Psi}}{\partial \rho} \quad \text{on } \rho = a$$

$$\sigma_e \frac{\partial \phi_e^0}{\partial \rho} = -C_m \frac{\partial \phi_m^0}{\partial t} - \widetilde{I}_m - \sigma_e \frac{\partial \tilde{\Psi}}{\partial \rho} \quad \text{on } \rho = a$$

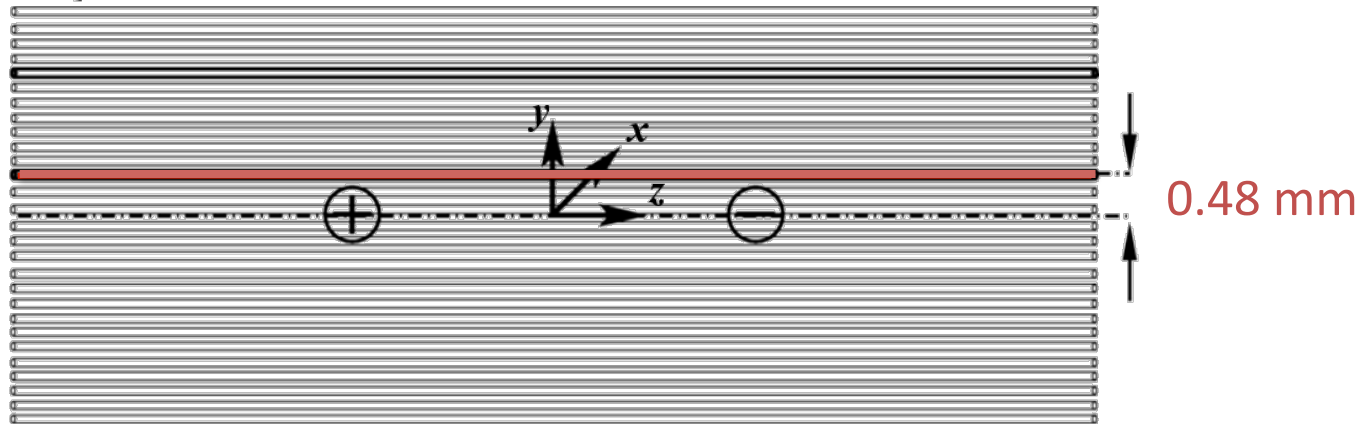


Mean equation of the longitudinal problem

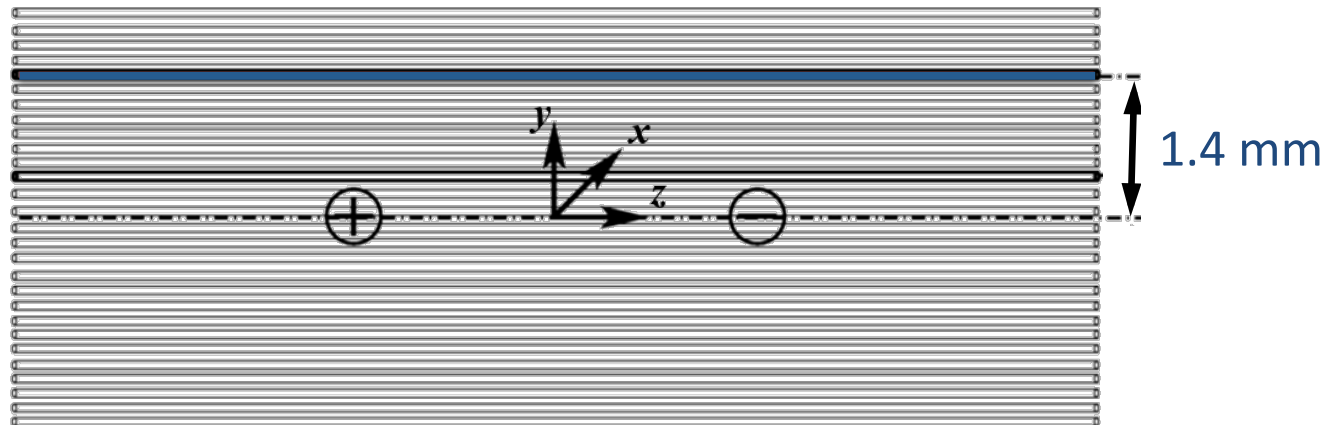
$$\frac{\sigma_i a}{2} \frac{\partial^2 f_m^0}{\partial z^2} = C_m \frac{\partial f_m^0}{\partial t} + \overline{I}_m - \frac{\sigma_i a}{2} \frac{\partial^2 \langle \Psi \rangle}{\partial z^2}$$

Uptake Results Graphical Aid

Close Fiber

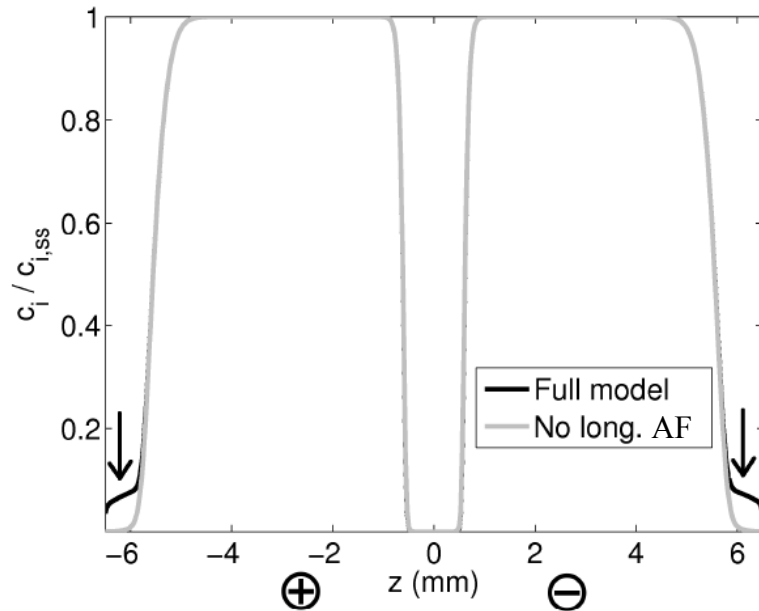


Far Fiber



Effect of Longitudinal Charging

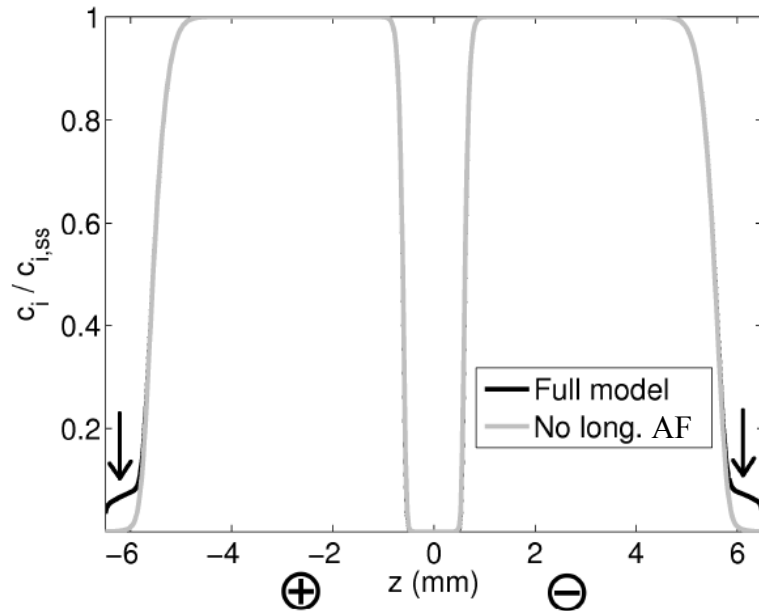
Close Fiber



Longitudinal charging: 0.91 %

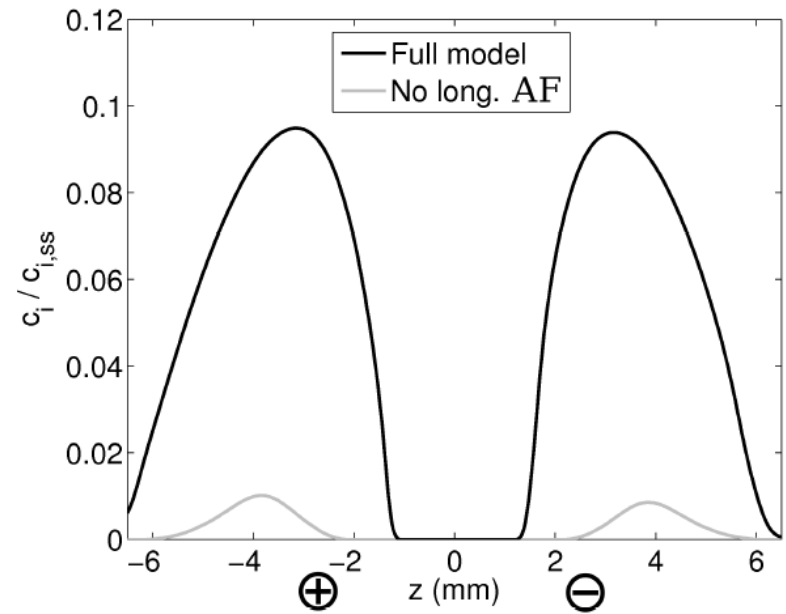
Effect of Longitudinal Charging

Close Fiber



Longitudinal charging: 0.91 %

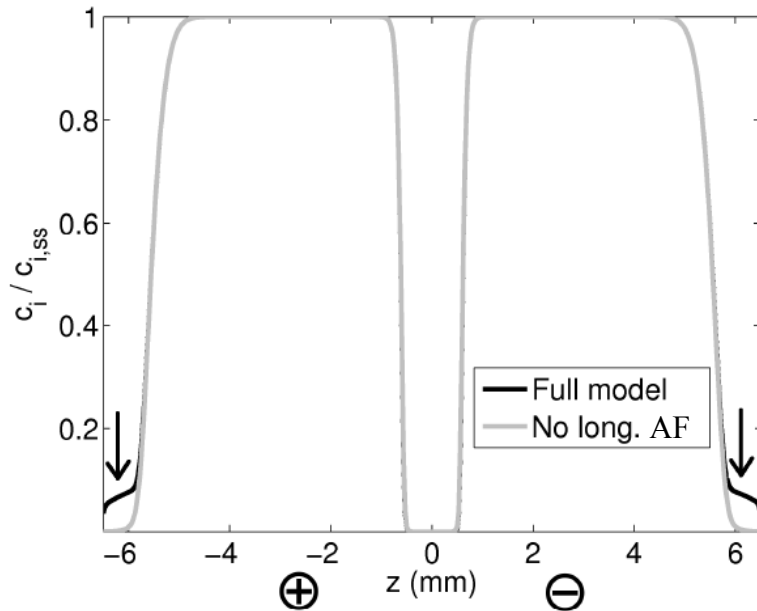
Far Fiber



Longitudinal charging: 2000 %

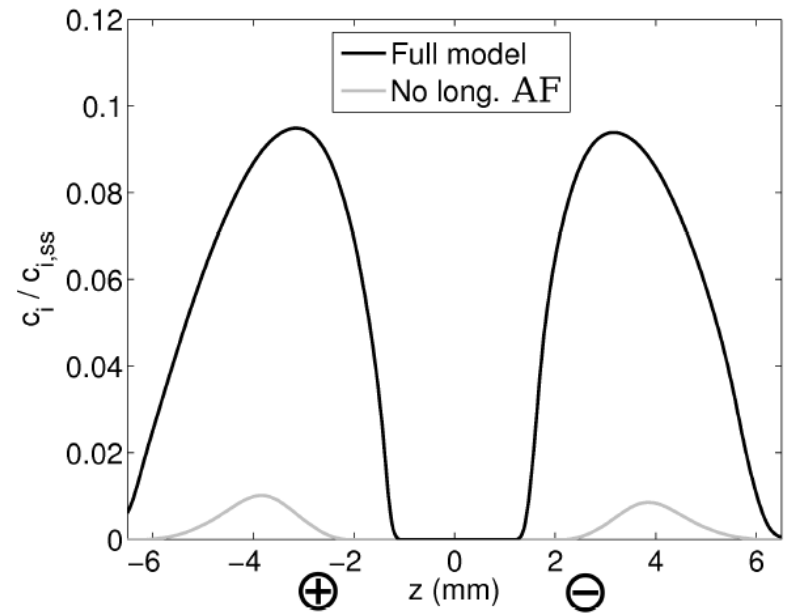
Effect of Longitudinal Charging

Close Fiber

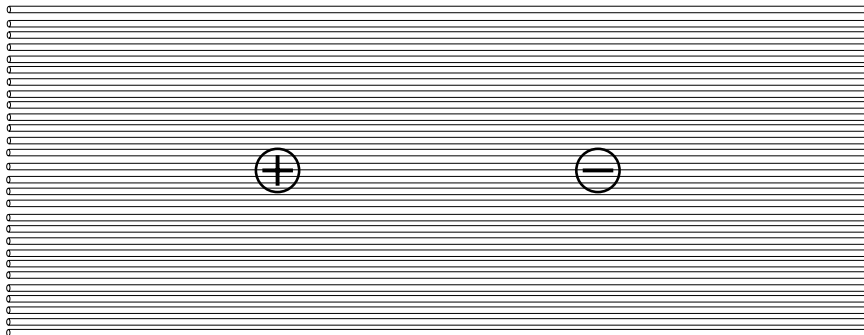


Longitudinal charging: 0.91 %

Far Fiber



Longitudinal charging: 2000 %

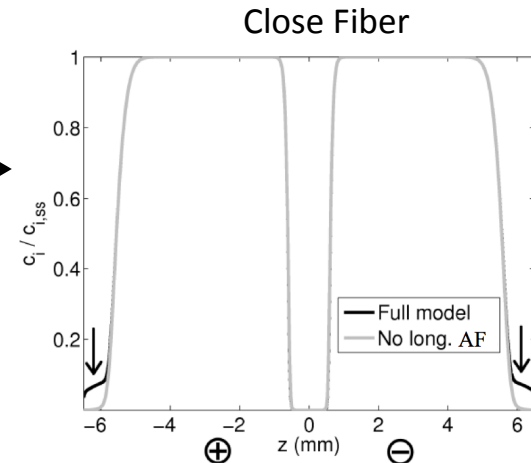
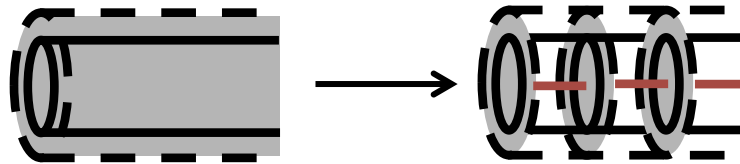


Total uptake over all cells in tissue is increased by only 3.95% by including longitudinal AF, which is less than experimental variability (6%)

Versatility of Asymptotic Fiber Model

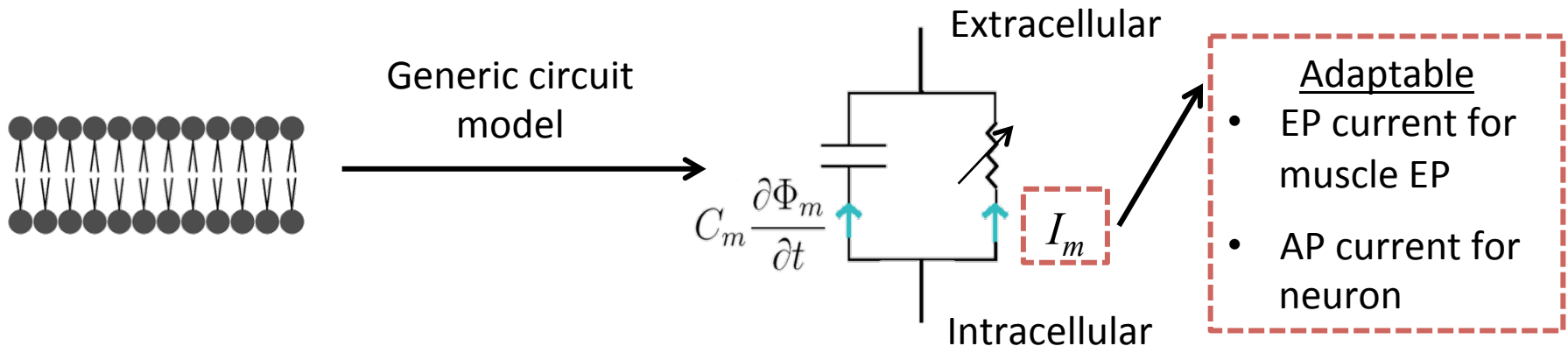
Neuron modeling community interested in stimulation of fibers

- Emphasis on longitudinal membrane charging (1D core conductor model), but ...
- Longitudinal charging not important for fibers close to electrodes →
- Computationally efficient model to study transverse, as well as longitudinal charging



Model is valid for neurons

- Axon of neuron has similar geometry and electrical properties
- Membrane level model adapts to many physiological phenomena



Conclusions

- Longitudinal charging dominates for skeletal muscle fibers far from electrodes, while transverse charging dominates for fibers close to the electrodes
- Majority of molecular uptake occurs for fibers close to the electrodes, thus, I corroborated traditional transverse-only models of electroporation in skeletal muscle
- Asymptotic model is an efficient and versatile tool to simulate bioelectric phenomena
- Presence of transverse charging in asymptotic model may have significant added value in simulating activation of neurons

Acknowledgements

Thank You

All Individuals in the Audience

Duke University Administrative and Technical Support

Kathy Barbour
Ned Danieleley
Susan Story
Dr. Adam P. Wax
Duke OIT

Colleagues and Mentors

Dr. Roger C. Barr	Dr. Ann Pitruzzello
Dr. Margot E. Bowen	Dr. Caroline Ring
Dr. James Esterline	Dr. John Wambaugh
Dr. Mazella B. Fuller	Dr. Thomas P. Witelski
Dr. Thomas J. McIntosh	Dr. Fan Yuan
Dr. John C. Neu	
Dr. Wanda K. Neu	

Extra Slides

Experimental Studies and Modeling

Goal in published studies is to maximize gene expression

- Optimize pulsing protocol
- Electrode geometry

Modeling DNA and expression: complex

- Relate DNA uptake to expression
- DNA is long polymer
 - Stretches/shrinks
 - Endocytotic mechanisms
 - Membrane adsorption
 - Molecular dynamics (expensive)

Modeling small molecules: less complex

- Radiolabelled molecules
- Radioactivity → molecular concentration
- Molecular concentration: convection-diffusion equation (less expensive)

Theoretical Studies

Membrane level

Creation of pores and current through pores in response to Φ_m

Molecular dynamics model

- Motion of each lipid molecule
- Not suitable for spatially extended system

Asymptotic model of EP¹

- Simple ODE
- Membrane of spherical cells

Longitudinal behavior of EP in fibers not characterized

¹Neu, J.C. & Krassowska, W. (1999) *PRE*, 59, 3471-3482

Theoretical Studies

Cellular level

Distribution of Φ_m due to microscale geometry of individual cells

EP community models

- $|\Phi_m|^\alpha \propto |E|$

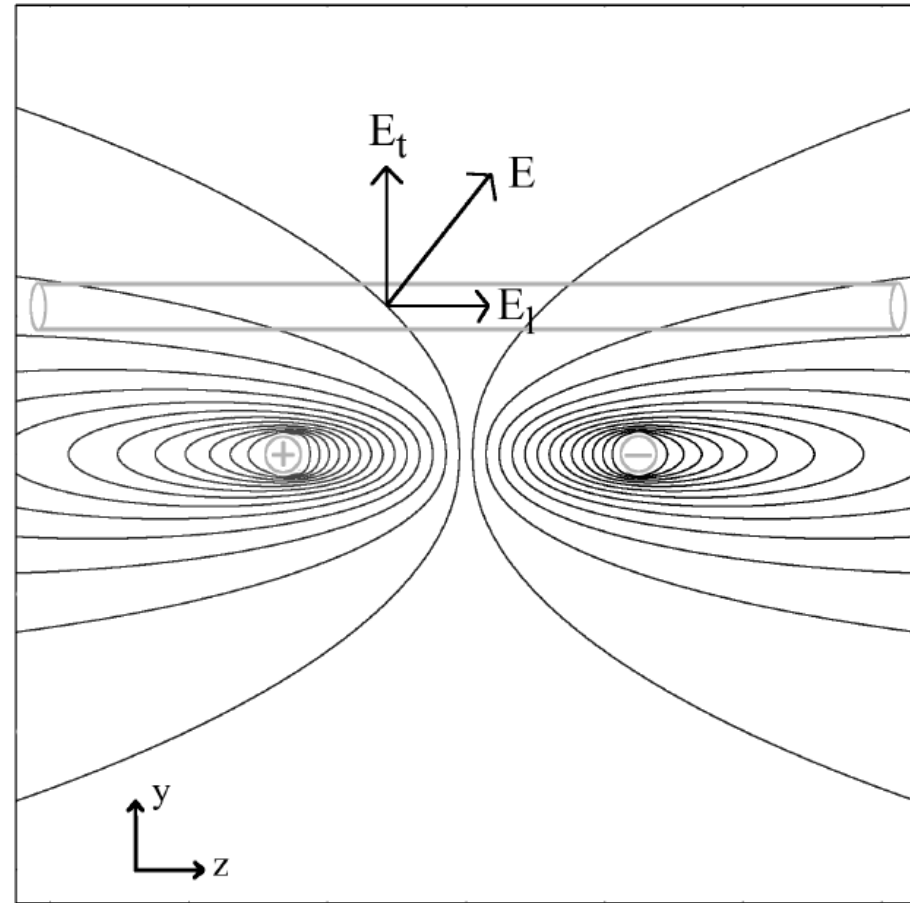
Theoretical Studies

Cellular level

Distribution of Φ_m due to microscale geometry of individual cells

EP community models

- $|\Phi_m| \propto |E|$



Theoretical Studies

Cellular level

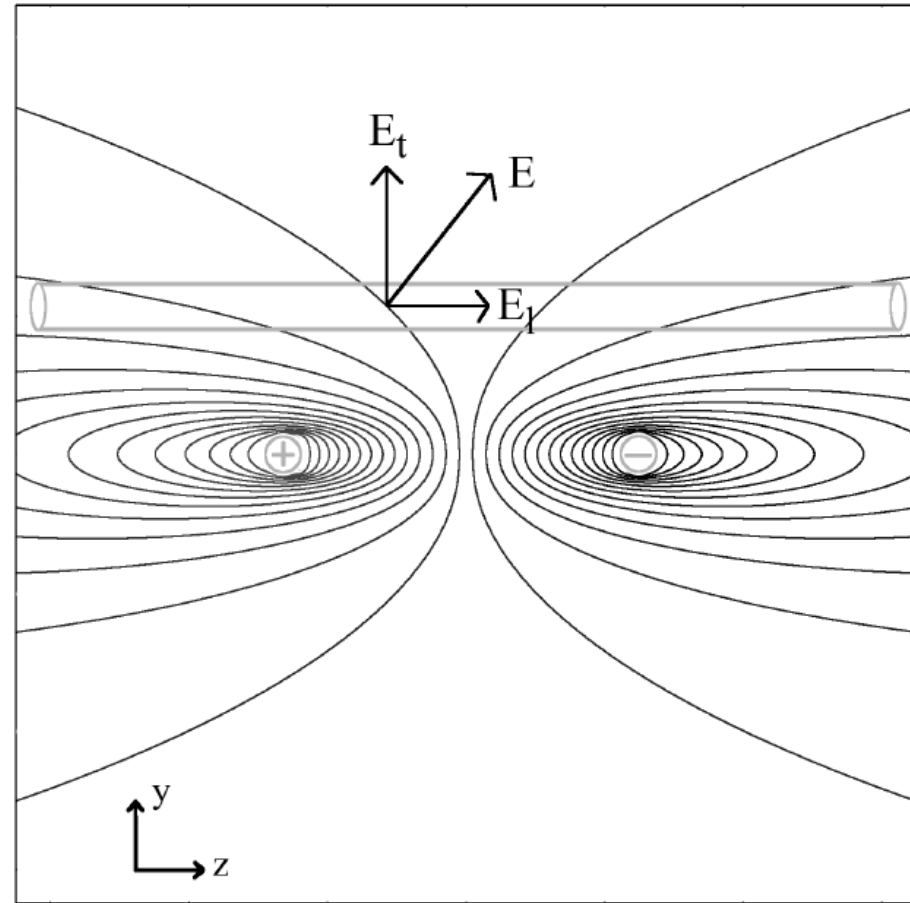
Distribution of Φ_m due to microscale geometry of individual cells

EP community models

- $|\Phi_m| \propto |E|$

Neuron community models

- “transverse charging”
 $|\Phi_m| \propto |E_t|$
- “longitudinal charging”
 $|\Phi_m| \propto \left| \frac{\partial E_l}{\partial z} \right|$
- $\left| \frac{\partial E_l}{\partial z} \right|$ varies by 7 orders of mag.



Effect of longitudinal charging on EP not characterized

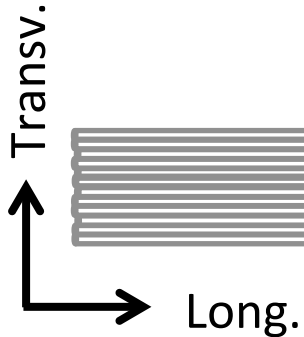
Theoretical Studies

Tissue level

Connects macroscale tissue distribution of potential from electric stimulus to Φ_m at **cellular level**

Ideal tissue model

Individual fibers



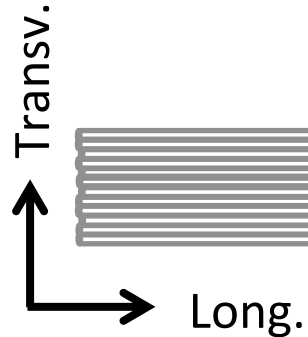
Theoretical Studies

Tissue level

Connects macroscale tissue distribution of potential from electric stimulus to Φ_m at **cellular level**

Ideal tissue model

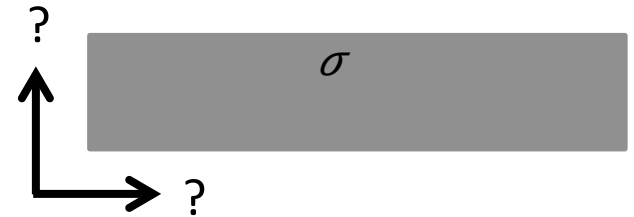
Individual fibers



Vs.

Existing models

No individual fiber geometry



Existing models cannot capture transverse and longitudinal membrane charging

Research Aims

1. Investigate mechanism of transverse and longitudinal membrane charging during EP
2. Determine contribution of longitudinal charging to molecular uptake during EP
3. Compare uptake to traditional models of EP in muscle that neglect individual fiber geometry, and thus, directional membrane charging

Outline

Introduction

- Background
- Published experiments and models

Derivation and Validation of Models

- **Asymptotic fiber model**
- Alterations for EP-mediated uptake
- Mass transport model

Results

- Single fiber close to electrodes
- Single fiber far from electrodes
- Model variations and tissue-wide effect

Conclusions and Connections to LLNL

Theoretical Studies

Tissue level

- Macroscale distribution of potential throughout entire tissue
- Connects electric stimulus to membrane charging at cellular level

Liver tissue models

- Compact, spherical cells
- Asymptotic model of EP

Muscle tissue models

- No individual cell geometry
- Cannot capture transverse and longitudinal membrane charging

Effect of individual fiber geometry on EP in muscle tissue not characterized

Theoretical Studies

Mass transport

Movement of small molecules throughout tissue and across membrane

Transport mechanisms

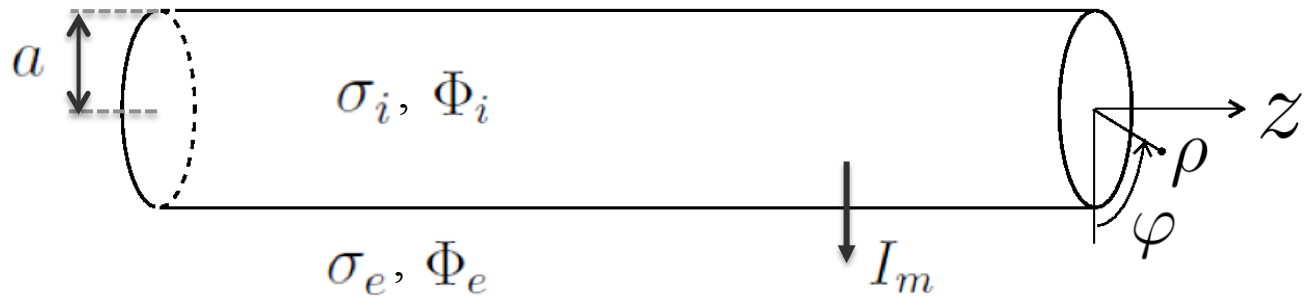
- Electric drift
- Diffusion
- Convection-diffusion equation

Applied to cancer tissue

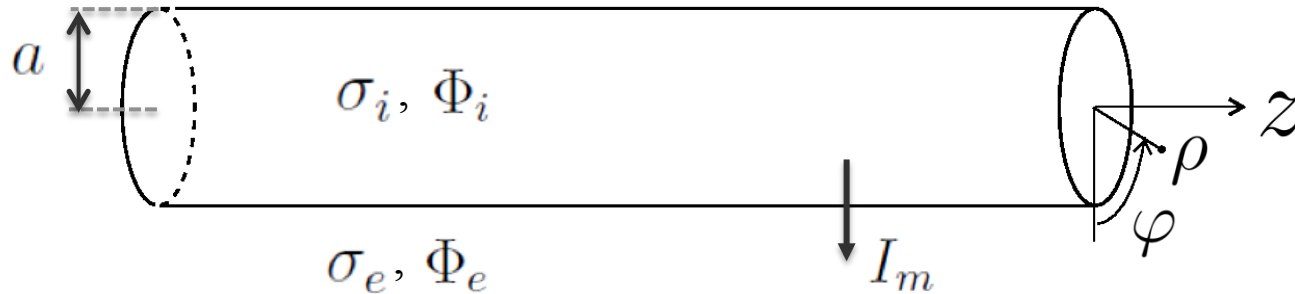
- Compact spherical cells

Uptake in muscle tissue with longitudinal and transverse membrane charging not characterized

Derivation of Asymptotic Fiber Model



Derivation of Asymptotic Fiber Model



Scaled governing equations in dimensionless form

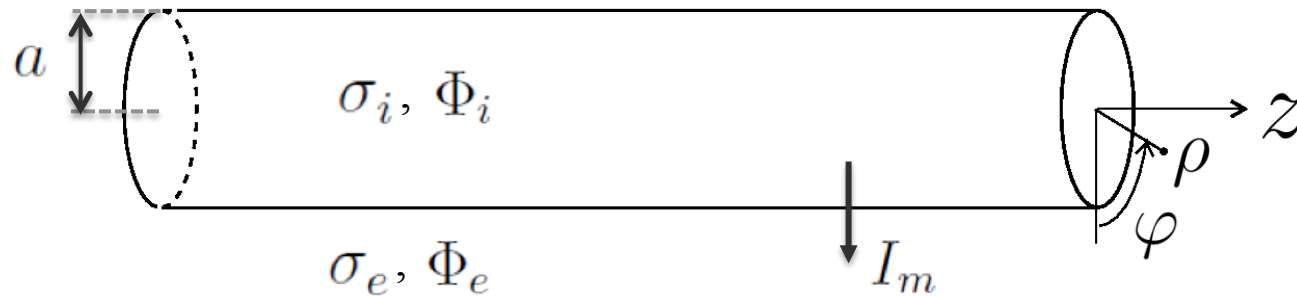
$$\frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial \Phi_i}{\partial \rho} \right) + \frac{1}{\rho^2} \frac{\partial^2 \Phi_i}{\partial \varphi^2} + \epsilon \frac{\partial^2 \Phi_i}{\partial z^2} = 0 \quad \text{in } \rho < a,$$

$$\frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial \Phi_e}{\partial \rho} \right) + \frac{1}{\rho^2} \frac{\partial^2 \Phi_e}{\partial \varphi^2} + \epsilon \frac{\partial^2 \Phi_e}{\partial z^2} = 0 \quad \text{in } \rho > a,$$

$$\frac{\partial \Phi_i}{\partial \rho} = -\frac{\partial \Phi_m}{\partial \tau} - \epsilon \frac{\partial \Phi_m}{\partial T_s} - I_m - \frac{\partial \Psi}{\partial \rho} \quad \text{on } \rho = a$$

$$\frac{\partial \Phi_e}{\partial \rho} = -\frac{\sigma_i}{\sigma_e} \left\{ \frac{\partial \Phi_m}{\partial \tau} + \epsilon \frac{\partial \Phi_m}{\partial T_s} + I_m \right\} - \frac{\partial \Psi}{\partial \rho} \quad \text{on } \rho = a$$

Derivation of Asymptotic Fiber Model



Scaled governing equations in dimensionless form

$$\frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial \Phi_i}{\partial \rho} \right) + \frac{1}{\rho^2} \frac{\partial^2 \Phi_i}{\partial \varphi^2} + \epsilon \frac{\partial^2 \Phi_i}{\partial z^2} = 0 \quad \text{in } \rho < a,$$

$$\frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial \Phi_e}{\partial \rho} \right) + \frac{1}{\rho^2} \frac{\partial^2 \Phi_e}{\partial \varphi^2} + \epsilon \frac{\partial^2 \Phi_e}{\partial z^2} = 0 \quad \text{in } \rho > a,$$

$$\frac{\partial \Phi_i}{\partial \rho} = -\frac{\partial \Phi_m}{\partial \tau} - \epsilon \frac{\partial \Phi_m}{\partial T_s} - I_m - \frac{\partial \Psi}{\partial \rho} \quad \text{on } \rho = a$$

$$\frac{\partial \Phi_e}{\partial \rho} = -\frac{\sigma_i}{\sigma_e} \left\{ \frac{\partial \Phi_m}{\partial \tau} + \epsilon \frac{\partial \Phi_m}{\partial T_s} + I_m \right\} - \frac{\partial \Psi}{\partial \rho} \quad \text{on } \rho = a$$

Derivation of Asymptotic Fiber Model

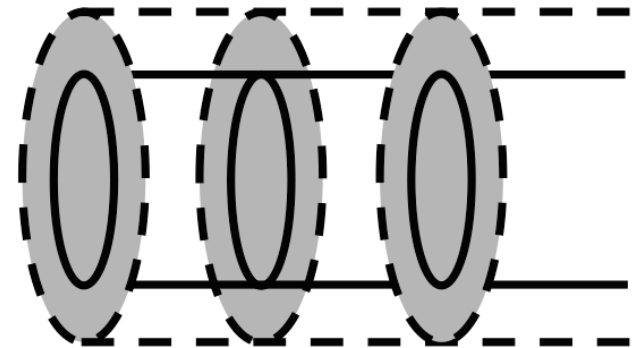
Mean-free equations of the transverse problem

$$\frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial \phi_i^0}{\partial \rho} \right) + \frac{1}{\rho^2} \frac{\partial^2 \phi_i^0}{\partial \varphi^2} = 0 \quad \text{in } \rho < a$$

$$\frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial \phi_e^0}{\partial \rho} \right) + \frac{1}{\rho^2} \frac{\partial^2 \phi_e^0}{\partial \varphi^2} = 0 \quad \text{in } \rho > a$$

$$\sigma_i \frac{\partial \phi_i^0}{\partial \rho} = -C_m \frac{\partial \phi_m^0}{\partial t} - \widetilde{I}_m - \sigma_i \frac{\partial \widetilde{\Psi}}{\partial \rho} \quad \text{on } \rho = a$$

$$\sigma_e \frac{\partial \phi_e^0}{\partial \rho} = -C_m \frac{\partial \phi_m^0}{\partial t} - \widetilde{I}_m - \sigma_e \frac{\partial \widetilde{\Psi}}{\partial \rho} \quad \text{on } \rho = a$$



Derivation of Asymptotic Fiber Model

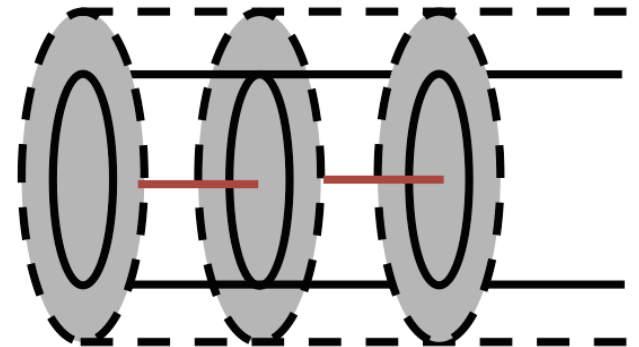
Mean-free equations of the transverse problem

$$\frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial \phi_i^0}{\partial \rho} \right) + \frac{1}{\rho^2} \frac{\partial^2 \phi_i^0}{\partial \varphi^2} = 0 \quad \text{in } \rho < a$$

$$\frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial \phi_e^0}{\partial \rho} \right) + \frac{1}{\rho^2} \frac{\partial^2 \phi_e^0}{\partial \varphi^2} = 0 \quad \text{in } \rho > a$$

$$\sigma_i \frac{\partial \phi_i^0}{\partial \rho} = -C_m \frac{\partial \phi_m^0}{\partial t} - \widetilde{I}_m - \sigma_i \frac{\partial \widetilde{\Psi}}{\partial \rho} \quad \text{on } \rho = a$$

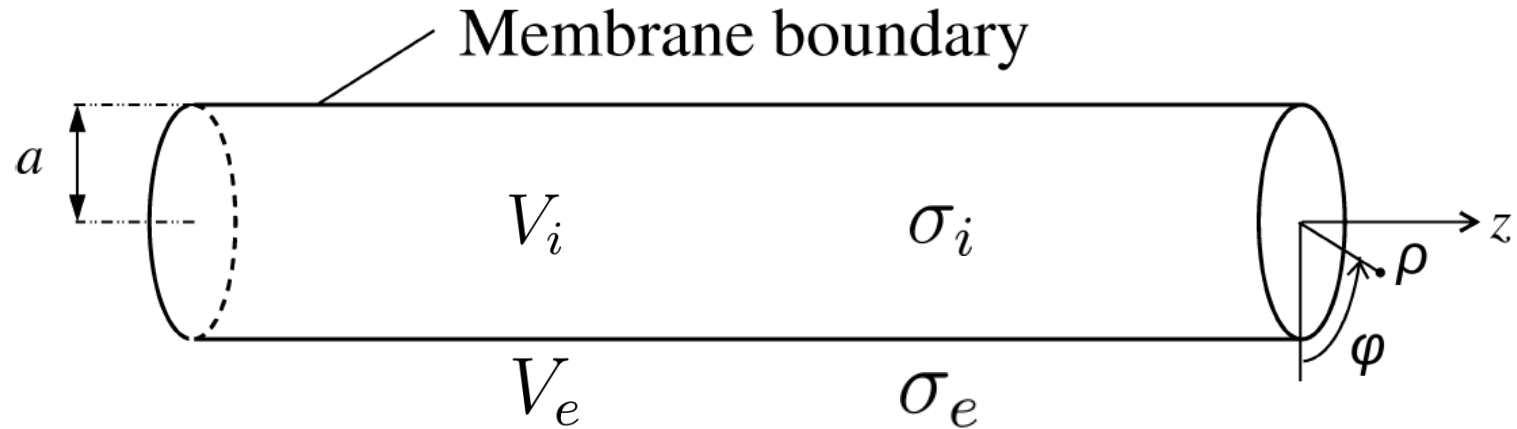
$$\sigma_e \frac{\partial \phi_e^0}{\partial \rho} = -C_m \frac{\partial \phi_m^0}{\partial t} - \widetilde{I}_m - \sigma_e \frac{\partial \widetilde{\Psi}}{\partial \rho} \quad \text{on } \rho = a$$



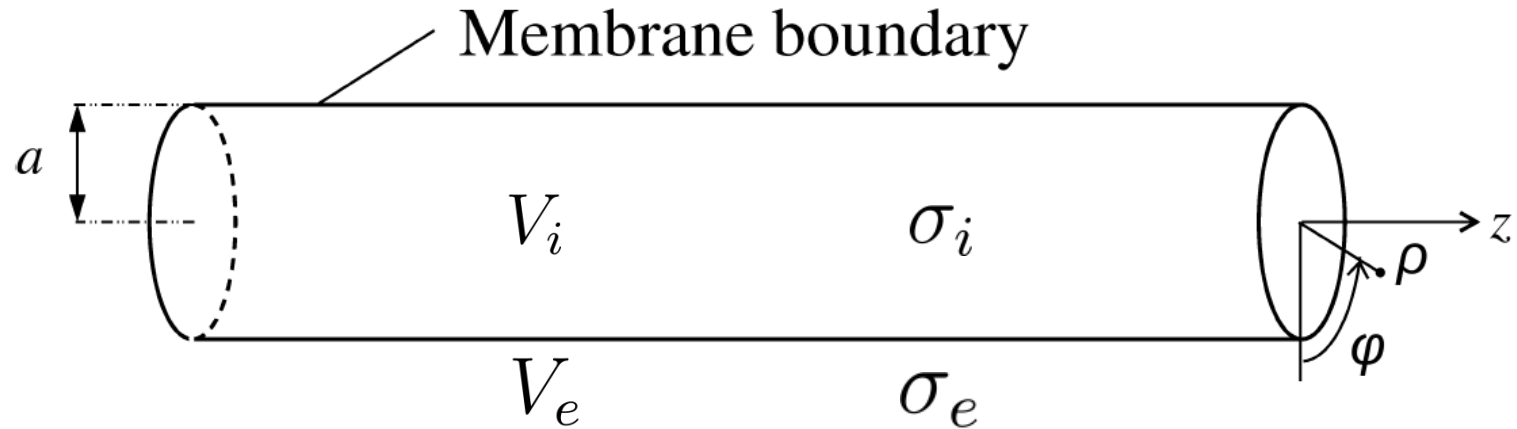
Mean equation of the longitudinal problem

$$\frac{\sigma_i a}{2} \frac{\partial^2 f_m^0}{\partial z^2} = C_m \frac{\partial f_m^0}{\partial t} + \overline{I}_m - \frac{\sigma_i a}{2} \frac{\partial^2 \langle \Psi \rangle}{\partial z^2}$$

Derivation of Asymptotic Fiber Model



Derivation of Asymptotic Fiber Model



Split potential

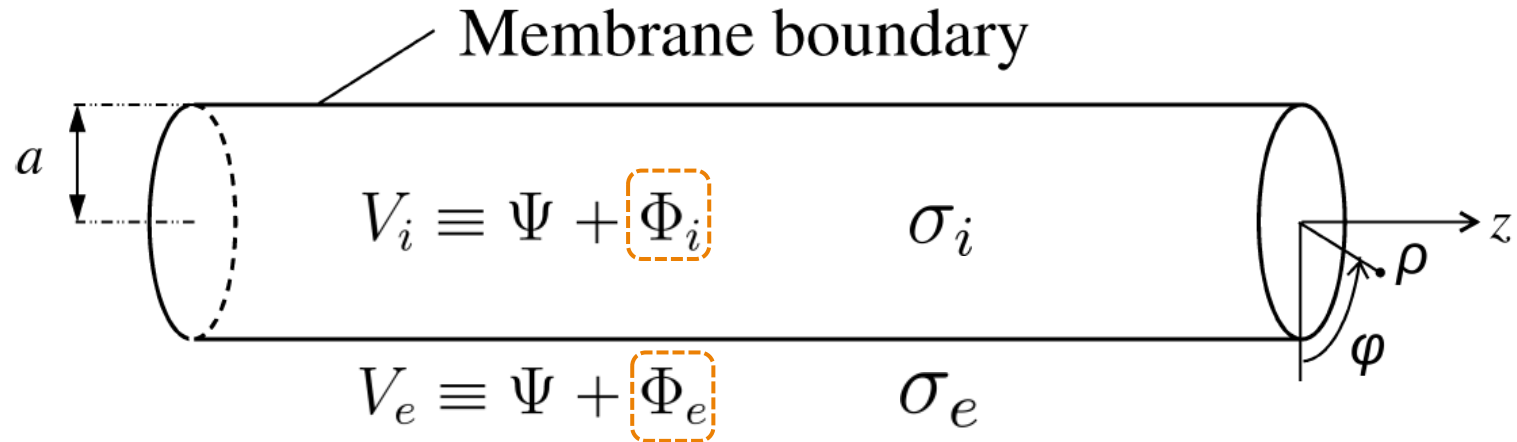
$$V_i \equiv \boxed{\Psi} + \Phi_i$$

$$V_e \equiv \boxed{\Psi} + \Phi_e$$

Primary potential $\boxed{\Psi}$

- Excludes microscopic presence of fibers
- Analytical solution

Derivation of Asymptotic Fiber Model



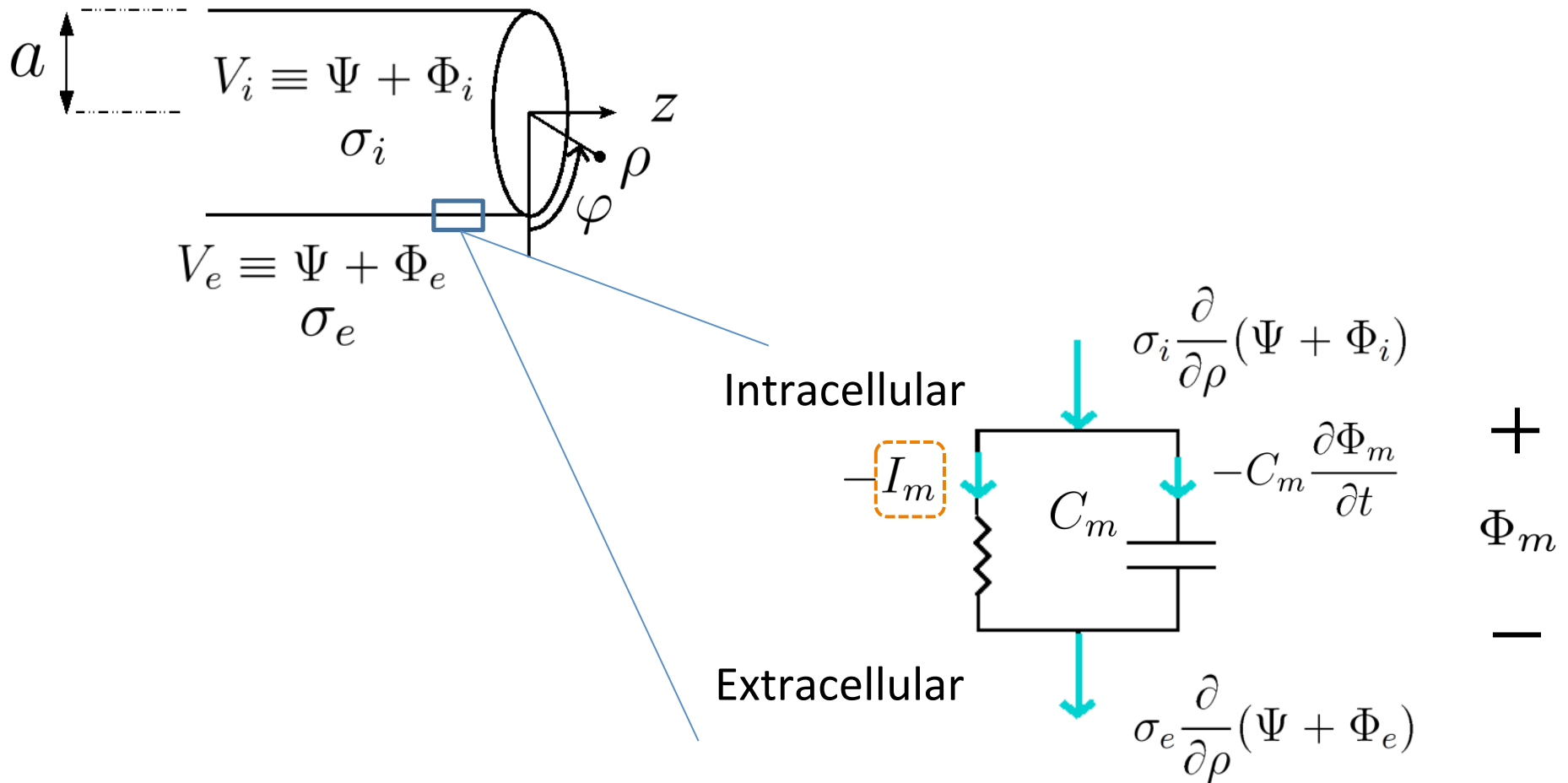
Secondary potentials Φ_i and Φ_e

- Includes microscopic presence of fiber
- Laplace's equation

$$\frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial \Phi_i}{\partial \rho} \right) + \frac{1}{\rho^2} \frac{\partial^2 \Phi_i}{\partial \varphi^2} + \frac{\partial^2 \Phi_i}{\partial z^2} = 0 \quad \text{in } \rho < a$$

$$\frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial \Phi_e}{\partial \rho} \right) + \frac{1}{\rho^2} \frac{\partial^2 \Phi_e}{\partial \varphi^2} + \frac{\partial^2 \Phi_e}{\partial z^2} = 0 \quad \text{in } \rho > a$$

Derivation of Asymptotic Fiber Model



- Matching conditions

$$\sigma_i \frac{\partial}{\partial \rho} (\Psi + \Phi_i) = \sigma_e \frac{\partial}{\partial \rho} (\Psi + \Phi_e) = -C_m \frac{\partial \Phi_m}{\partial t} - \boxed{I_m} \quad \text{on } \rho = a$$

Derivation of Asymptotic Fiber Model

Spatial scales

Short: microns (fiber radius)

Long: millimeters (length scale of fiber)

Temporal scales

Fast: microseconds (charging by transverse currents, aC_m/σ_i)

Slow: milliseconds (charging by longitudinal currents, $R_m C_m$)

Separate *fast, short*-distance response from *slow, long*-distance response

Method of multiple scales $\frac{\partial}{\partial t} = \frac{\partial}{\partial \tau} + \epsilon \frac{\partial}{\partial T_s}$

Small parameter: $\epsilon = \frac{aC_m/\sigma_i}{R_m C_m} = \frac{a}{\sigma_i R_m} \rightarrow O(10^{-5})$

Derivation of Asymptotic Fiber Model

Scale and convert governing equations to dimensionless form

Expand potentials in powers of ϵ , e.g., $\Phi_i \sim \Phi_i^0 + \epsilon \Phi_i^1 + \dots$

Limit $\epsilon \rightarrow 0$, LO problem suggests separation of potentials

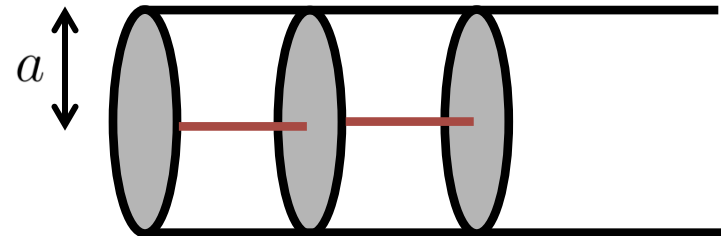
$$\Phi_i^0 = \phi_i^0(\rho, \varphi, \tau) + f_i^0(z, T_s)$$

$$\Phi_e^0 = \phi_e^0(\rho, \varphi, \tau) + f_e^0(z, T_s)$$

$$\Phi_m^0 = \underbrace{\phi_m^0(\varphi, \tau)}_{\text{mean-free}} + \underbrace{f_m^0(z, T_s)}_{\text{mean}}$$

Definition of mean and mean-free components

$$\underbrace{\phi_i^0}_{\text{mean-free}} = \underbrace{\Phi_i^0}_{\text{total}} - \underbrace{\frac{1}{\pi a^2} \int_0^{2\pi} \int_0^a \Phi_i^0 \rho d\rho d\varphi}_{\text{mean}}$$



Derivation of Asymptotic Fiber Model

Scale and convert governing equations to dimensionless form

$$\frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial \Phi_i}{\partial \rho} \right) + \frac{1}{\rho^2} \frac{\partial^2 \Phi_i}{\partial \varphi^2} + \epsilon \frac{\partial^2 \Phi_i}{\partial z^2} = 0 \quad \text{in } \rho < a$$

$$\frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial \Phi_e}{\partial \rho} \right) + \frac{1}{\rho^2} \frac{\partial^2 \Phi_e}{\partial \varphi^2} + \epsilon \frac{\partial^2 \Phi_e}{\partial z^2} = 0 \quad \text{in } \rho > a$$

$$\frac{\partial \Phi_i}{\partial \rho} = -\frac{\partial \Phi_m}{\partial \tau} - \epsilon \frac{\partial \Phi_m}{\partial T_s} - I_m - \frac{\partial \Psi}{\partial \rho} \quad \text{on } \rho = a$$

$$\frac{\partial \Phi_e}{\partial \rho} = -\frac{\sigma_i}{\sigma_e} \left\{ \frac{\partial \Phi_m}{\partial \tau} + \epsilon \frac{\partial \Phi_m}{\partial T_s} + I_m \right\} - \frac{\partial \Psi}{\partial \rho} \quad \text{on } \rho = a$$



Derivation of Asymptotic Fiber Model

Limit $\epsilon \rightarrow 0$, LO problem suggests separation of potentials

$$\Phi_i^0 = \phi_i^0(\rho, \varphi, \tau) + f_i^0(z, T_s)$$

$$\Phi_e^0 = \phi_e^0(\rho, \varphi, \tau) + f_e^0(z, T_s)$$

$$\Phi_m^0 = \underbrace{\phi_m^0(\varphi, \tau)}_{mean-free} + \underbrace{f_m^0(z, T_s)}_{mean}$$

Derivation of Asymptotic Fiber Model

Limit $\epsilon \rightarrow 0$, LO problem suggests separation of potentials

$$\Phi_i^0 = \phi_i^0(\rho, \varphi, \tau) + f_i^0(z, T_s)$$

$$\Phi_e^0 = \phi_e^0(\rho, \varphi, \tau) + f_e^0(z, T_s)$$

$$\Phi_m^0 = \underbrace{\phi_m^0(\varphi, \tau)}_{\text{mean-free}} + \underbrace{f_m^0(z, T_s)}_{\text{mean}}$$

Mean-free equations of the transverse problem

$$\frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial \phi_i^0}{\partial \rho} \right) + \frac{1}{\rho^2} \frac{\partial^2 \phi_i^0}{\partial \varphi^2} = 0 \quad \text{in } \rho < a,$$

$$\frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial \phi_e^0}{\partial \rho} \right) + \frac{1}{\rho^2} \frac{\partial^2 \phi_e^0}{\partial \varphi^2} = 0 \quad \text{in } \rho > a,$$

$$\sigma_i \frac{\partial \phi_i^0}{\partial \rho} = -C_m \frac{\partial \phi_m^0}{\partial t} - \widetilde{I}_m - \sigma_i \frac{\partial \widetilde{\Psi}}{\partial \rho} \quad \text{on } \rho = a$$

$$\sigma_e \frac{\partial \phi_e^0}{\partial \rho} = -C_m \frac{\partial \phi_m^0}{\partial t} - \widetilde{I}_m - \sigma_e \frac{\partial \widetilde{\Psi}}{\partial \rho} \quad \text{on } \rho = a$$

Derivation of Asymptotic Fiber Model

Mean-free equations of the transverse problem

$$\frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial \phi_i^0}{\partial \rho} \right) + \frac{1}{\rho^2} \frac{\partial^2 \phi_i^0}{\partial \varphi^2} = 0 \quad \text{in } \rho < a,$$

$$\frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial \phi_e^0}{\partial \rho} \right) + \frac{1}{\rho^2} \frac{\partial^2 \phi_e^0}{\partial \varphi^2} = 0 \quad \text{in } \rho > a,$$

$$\sigma_i \frac{\partial \phi_i^0}{\partial \rho} = -C_m \frac{\partial \phi_m^0}{\partial t} - \widetilde{I}_m - \sigma_i \frac{\partial \tilde{\Psi}}{\partial \rho} \quad \text{on } \rho = a$$

$$\sigma_e \frac{\partial \phi_e^0}{\partial \rho} = -C_m \frac{\partial \phi_m^0}{\partial t} - \widetilde{I}_m - \sigma_e \frac{\partial \tilde{\Psi}}{\partial \rho} \quad \text{on } \rho = a$$

$\partial \tilde{\Psi} / \partial \rho$: “transverse AF”

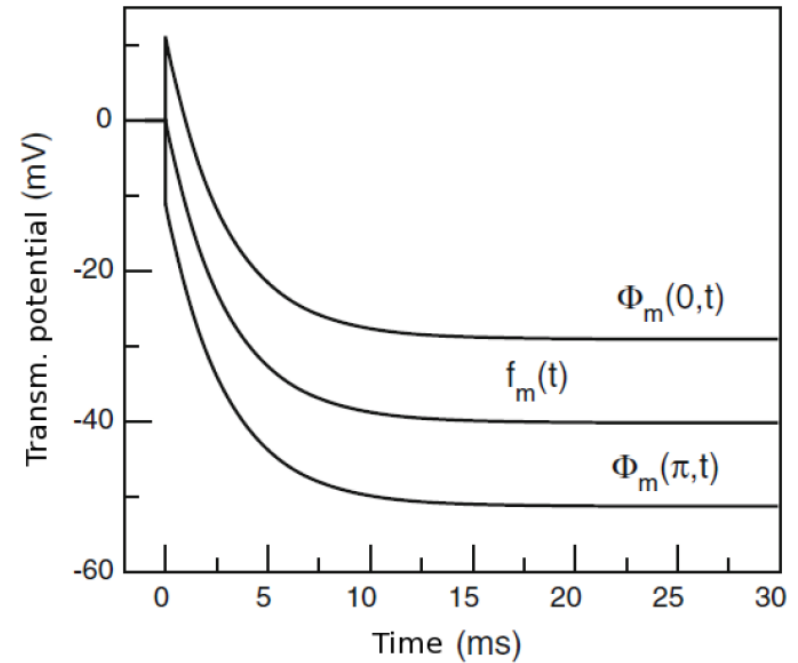
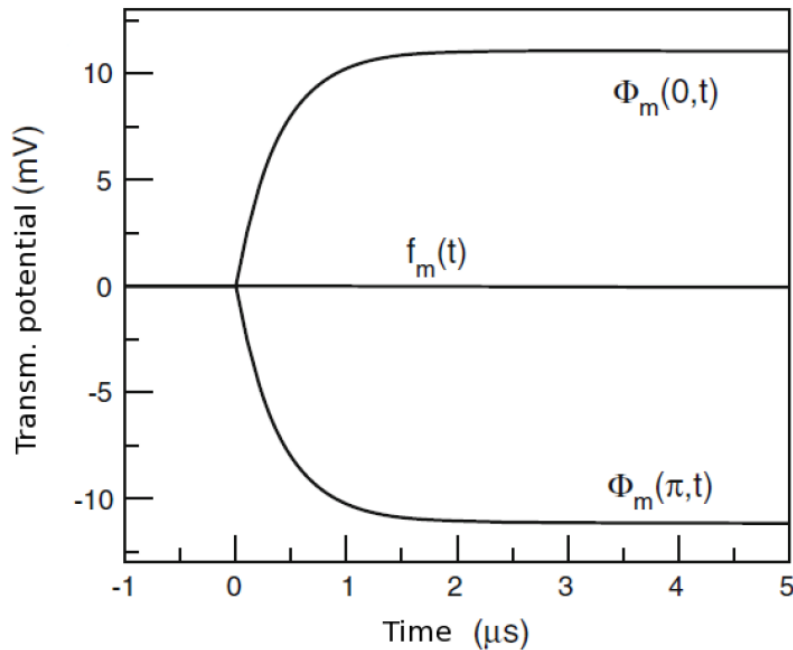
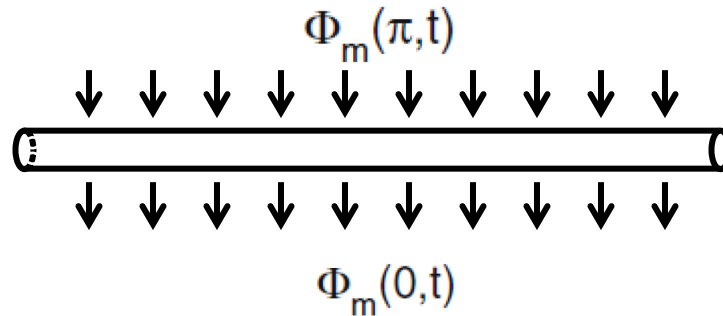
Mean equation of the longitudinal problem

$$\frac{\sigma_i a}{2} \frac{\partial^2 f_m^0}{\partial z^2} = C_m \frac{\partial f_m^0}{\partial t} + \overline{I}_m - \frac{\sigma_i a}{2} \frac{\partial^2 \langle \Psi \rangle}{\partial z^2}$$

$\frac{\partial^2 \langle \Psi \rangle}{\partial z^2}$: “longitudinal AF”

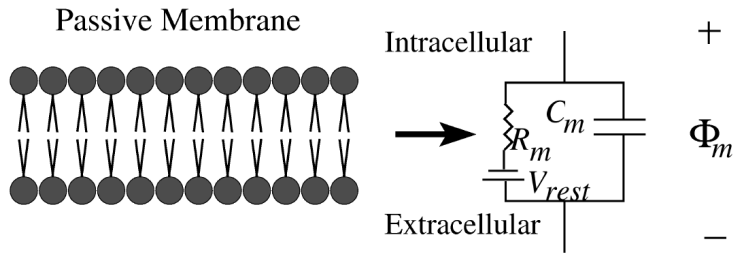
Validation of Asymptotic Fiber Model

Validation of temporal scales separation

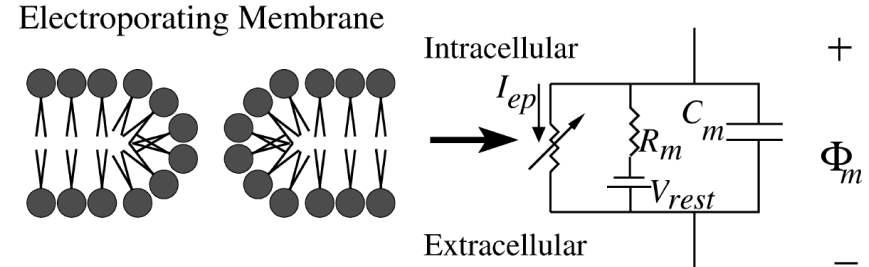


Alterations for EP-Mediated Uptake

Addition of current due to EP



$$I_m(\Phi_m) = \frac{\Phi_m - V_{rest}}{R_m}$$



$$I_m(\Phi_m) = \frac{\Phi_m - V_{rest}}{R_m} + I_{ep}$$

$$I_{ep} = N_{ep} i_{ep}$$

Pore density

Model: Asymptotic model of EP¹

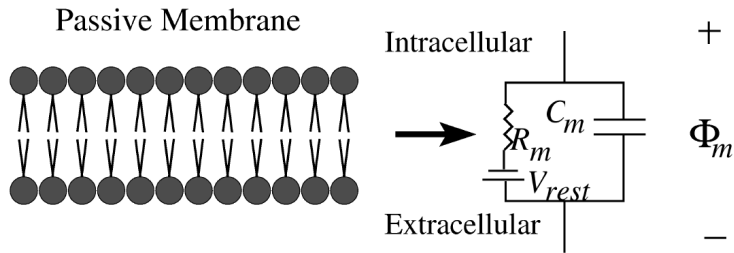
Assumptions

- Non-growing, constant pore radius
- Parameters fit to artificial membrane

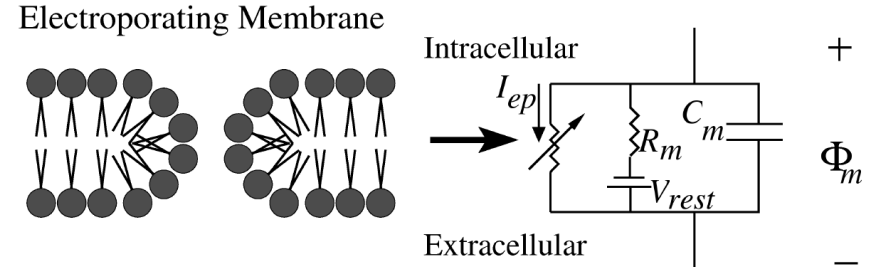
¹Neu, J.C. & Krassowska, W. (1999) *PRE*, 59, 3471-3482

Alterations for EP-Mediated Uptake

Addition of current due to EP



$$I_m(\Phi_m) = \frac{\Phi_m - V_{rest}}{R_m}$$



$$I_m(\Phi_m) = \frac{\Phi_m - V_{rest}}{R_m} + I_{ep}$$

$$I_{ep} = N_{ep} i_{ep}$$

Pore density

Model: Asymptotic model of EP¹

Assumptions

- Non-growing, constant pore radius
- Parameters fit to artificial membrane

Current through an electropore

Model: Aqueous pore²

Assumptions

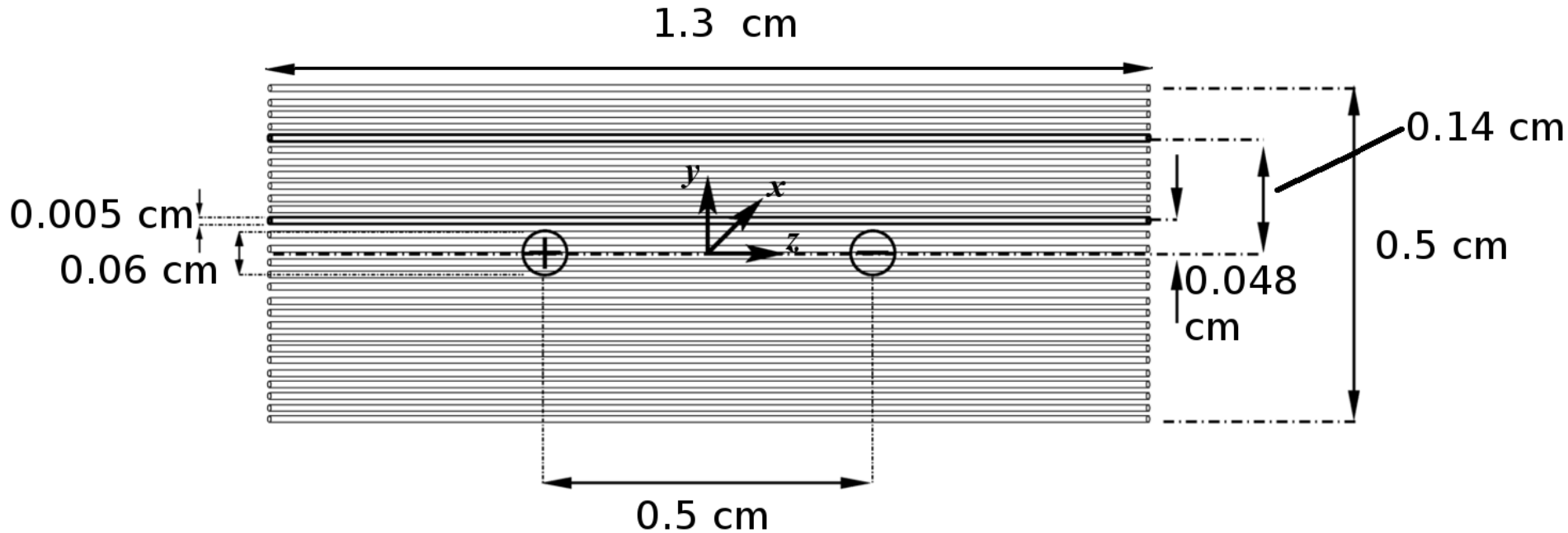
- Born energy due to low permittivity
- Two singularly charged ions (+ and -)
- Parameters fit to artificial membrane

¹Neu, J.C. & Krassowska, W. (1999) *PRE*, 59, 3471-3482

²Barnett, A. (1999) *Biochim. Biophys. Acta*, 1025, 10-14

Alterations for EP-Mediated Uptake

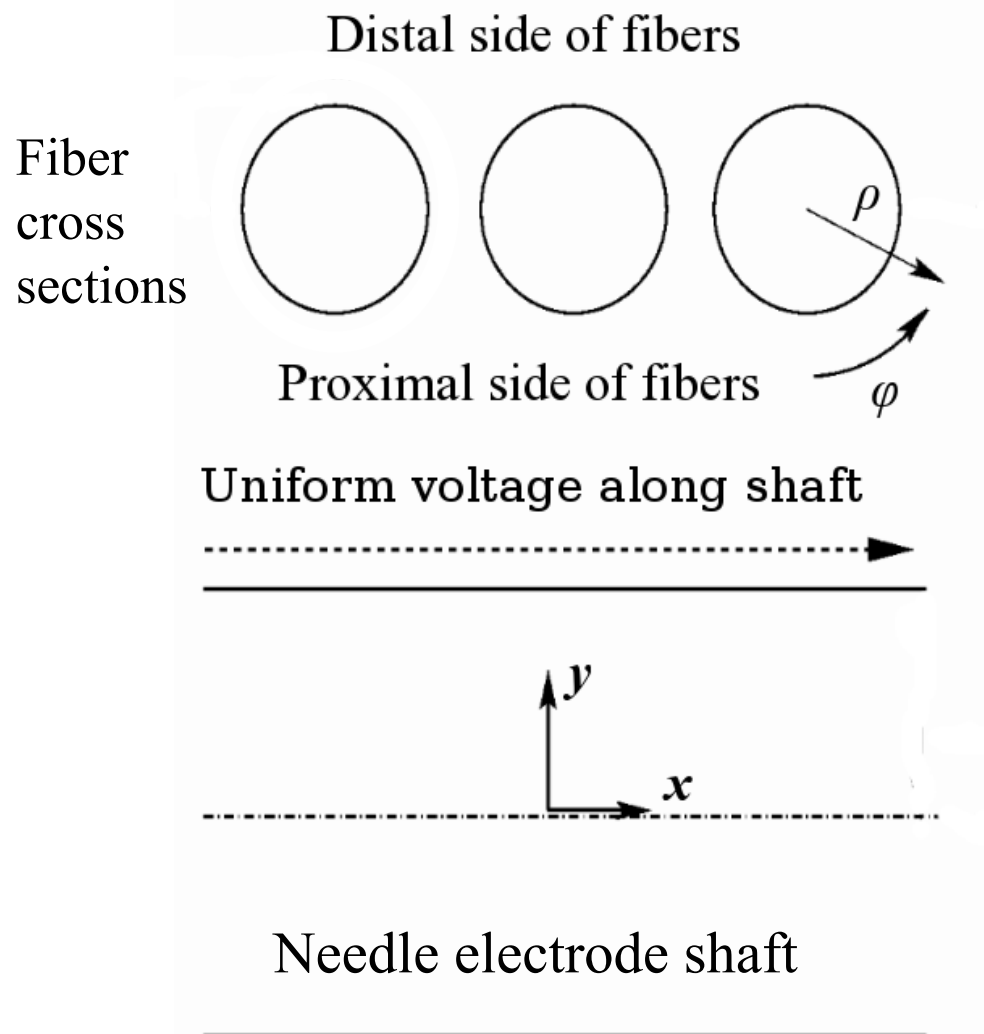
Conversion to muscle tissue



- Geometry from experiments
- Interfiber distance 1/50 fiber diameter
- Needle diameter is 12x fiber diameter
- Fiber length is 260x fiber diameter
- Far fiber is 3x distance of close fiber

Alterations for EP-Mediated Uptake

Conversion to muscle tissue



Fibers identical in x -dimension

Simulate fibers at different distances in y -dimension

- Each fiber has different AFs $\partial \tilde{\Psi} / \partial \rho$ and $\partial^2 \langle \Psi \rangle / \partial z^2$
- Analytical solution to primary potential Ψ in anisotropic medium¹

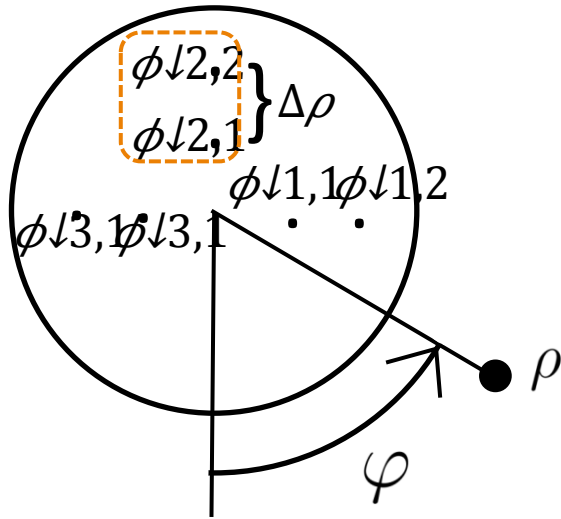
$$\sigma_y \frac{\partial^2 \Psi}{\partial y^2} + \sigma_z \frac{\partial^2 \Psi}{\partial z^2} = 0$$

¹Guo, L., Cranford, J.P., Neu, J.C., and Neu, W.K. (2009) *Med. Biol. Eng. Comput.*, 47, 1001-1010

Numerical Implementation and Simulation

Transverse 2D problem

Spatial: Finite difference method (FDM)



Discretization: $\frac{\partial \phi}{\partial \rho} \rightarrow \frac{\phi_{2,2} - \phi_{2,1}}{\Delta \rho}$

Temporal: Backward difference

Discretization: $\frac{\partial \phi}{\partial t} = \mathbb{F} \rightarrow \frac{\phi^{t+\Delta t} - \phi^t}{\Delta t} = \mathbb{F}^{t+\Delta t}$

Longitudinal 1D problem: Crank-Nicolson method

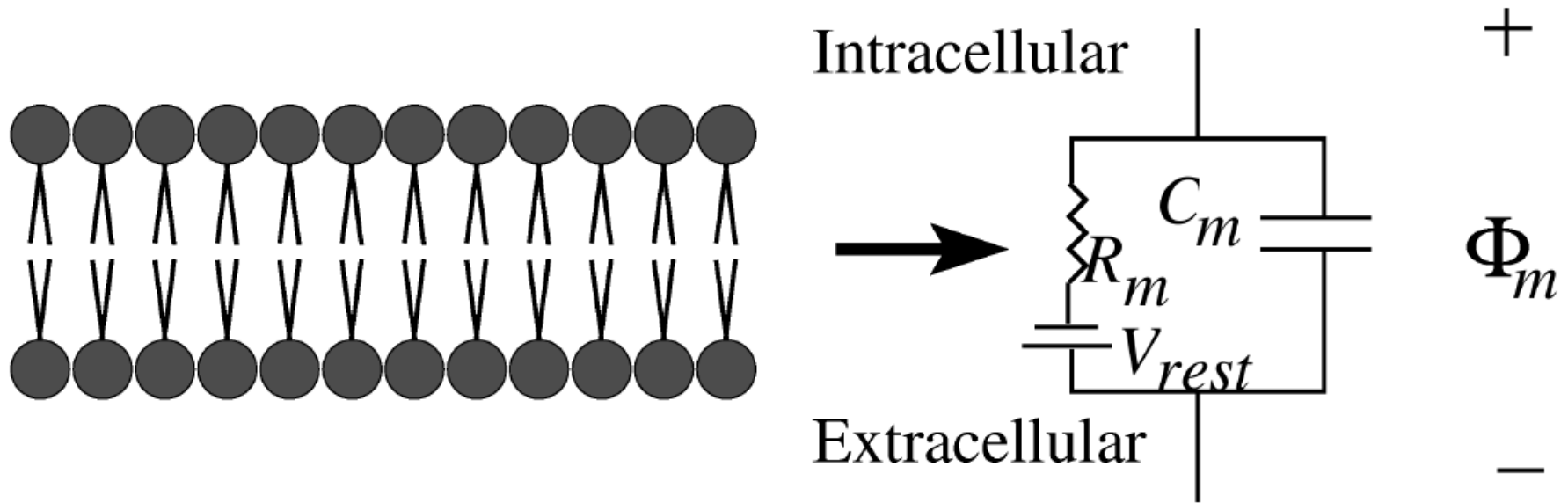
$f_{m-1}, f_m, f_{m+1} \rightarrow$ Discretization: $\frac{\partial f_m}{\partial t} = \mathbb{F} \rightarrow \frac{f_m^{t+\Delta t} - f_m^t}{\Delta t} = \frac{1}{2} \{ \mathbb{F}^{t+\Delta t} + \mathbb{F}^t \}$

Parallel computing?

System in “anti-Goldilocks zone”

Validation of Asymptotic Fiber Model

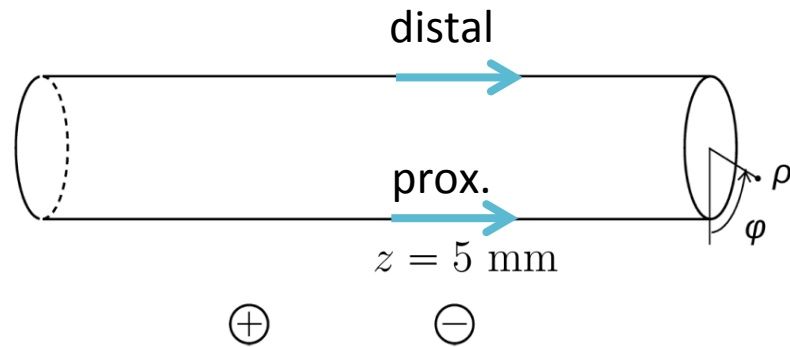
Passive membrane dynamics



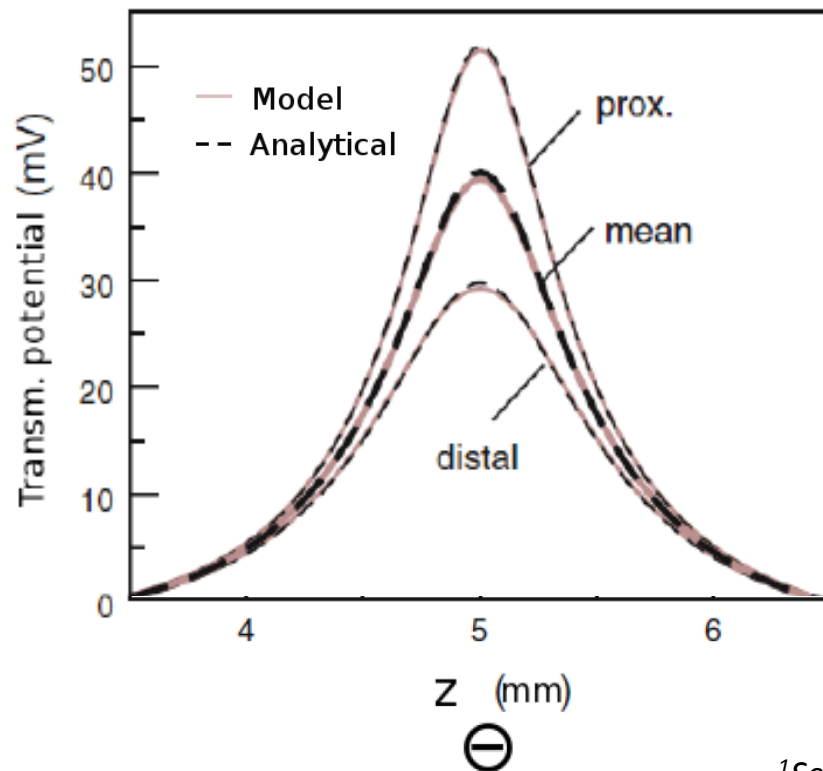
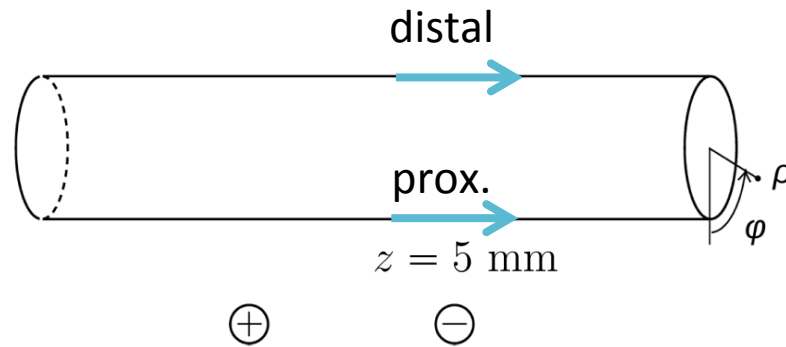
Q: Why passive (no pores/action potentials in membrane)?

A: So there is an analytical solution!

Validation of Asymptotic Fiber Model



Validation of Asymptotic Fiber Model



Compare distribution of steady state Φ_m in model to analytical solution¹

Root mean square error: 0.153 mV (0.295% of max $|\Phi_m|$)

¹Schnabel, V. and Struijk, J.J. (2001) IEEE T. Bio-Med. Eng., 48, 1027-1033

Outline

Introduction

- Background
- Published experiments and models

Derivation and Validation of Models

- Asymptotic fiber model
- **Alterations for EP-mediated uptake**
- Mass transport model

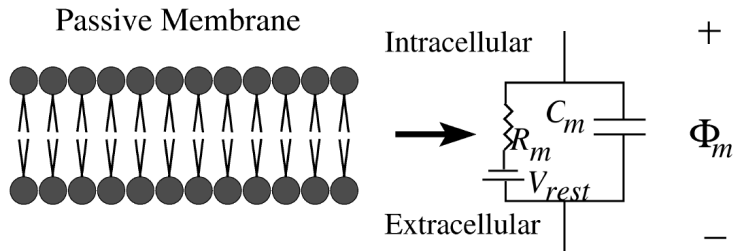
Results

- Single fiber close to electrodes
- Single fiber far from electrodes
- Model variations and tissue-wide effect

Conclusions and Connections to LLNL

Alterations for EP-Mediated Uptake

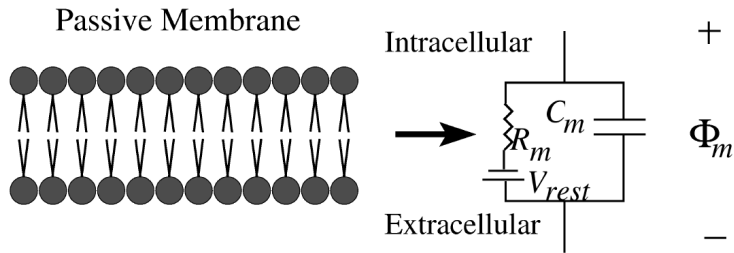
Addition of current due to EP



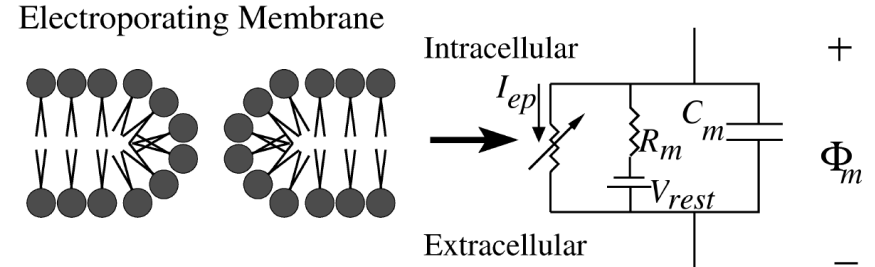
$$I_m(\Phi_m) = \frac{\Phi_m - V_{rest}}{R_m}$$

Alterations for EP-Mediated Uptake

Addition of current due to EP



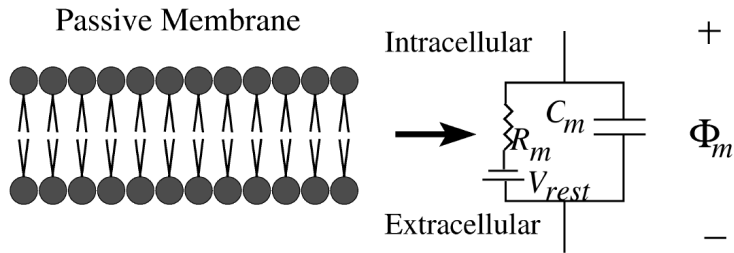
$$I_m(\Phi_m) = \frac{\Phi_m - V_{rest}}{R_m}$$



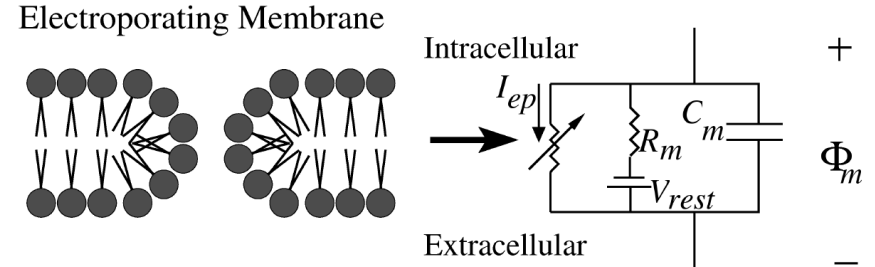
$$I_m(\Phi_m) = \frac{\Phi_m - V_{rest}}{R_m} + I_{ep}$$

Alterations for EP-Mediated Uptake

Addition of current due to EP



$$I_m(\Phi_m) = \frac{\Phi_m - V_{rest}}{R_m}$$

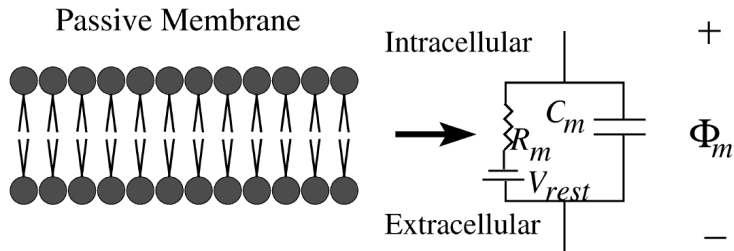


$$I_m(\Phi_m) = \frac{\Phi_m - V_{rest}}{R_m} + I_{ep}$$

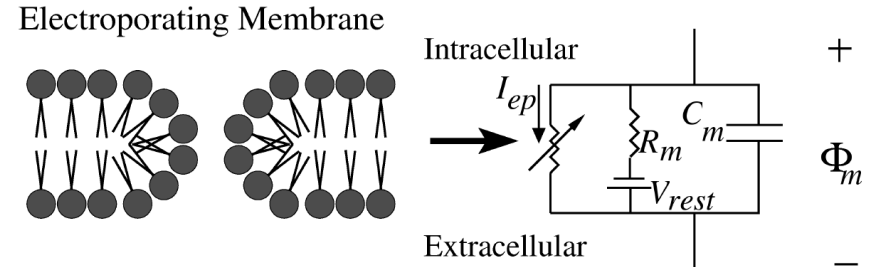
$$I_{ep} = N_{ep} i_{ep}$$

Alterations for EP-Mediated Uptake

Addition of current due to EP



$$I_m(\Phi_m) = \frac{\Phi_m - V_{rest}}{R_m}$$



$$I_m(\Phi_m) = \frac{\Phi_m - V_{rest}}{R_m} + I_{ep}$$

$$I_{ep} = N_{ep} i_{ep}$$

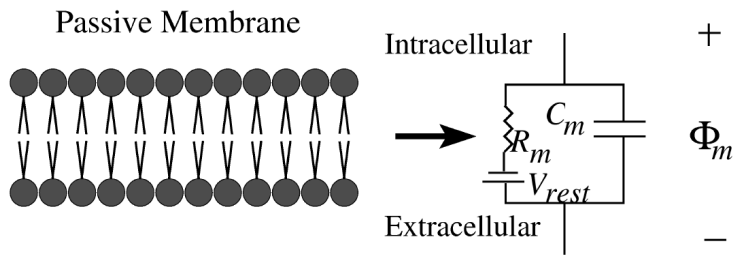
Pore density

Published ODE model ¹
(highly nonlinear)

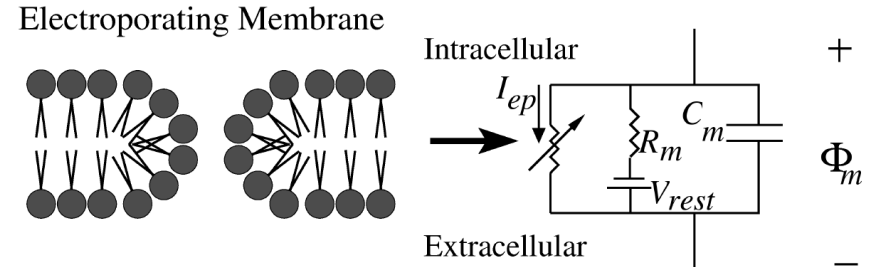
¹Neu, J.C. & Krassowska, W. (1999) *PRE*, 59, 3471-3482

Alterations for EP-Mediated Uptake

Addition of current due to EP



$$I_m(\Phi_m) = \frac{\Phi_m - V_{rest}}{R_m}$$



$$I_m(\Phi_m) = \frac{\Phi_m - V_{rest}}{R_m} + I_{ep}$$

$$I_{ep} = N_{ep} i_{ep}$$

Pore density

Published ODE model ¹
(highly nonlinear)

Current through an electropore

Published aqueous pore model ²

¹Neu, J.C. & Krassowska, W. (1999) *PRE*, 59, 3471-3482

²Barnett, A. (1999) *Biochim. Biophys. Acta*, 1025, 10-14

Outline

Introduction

- Background
- Published experiments and models

Derivation and Validation of Models

- Asymptotic fiber model
- Alterations for EP-mediated uptake
- **Mass transport model**

Results

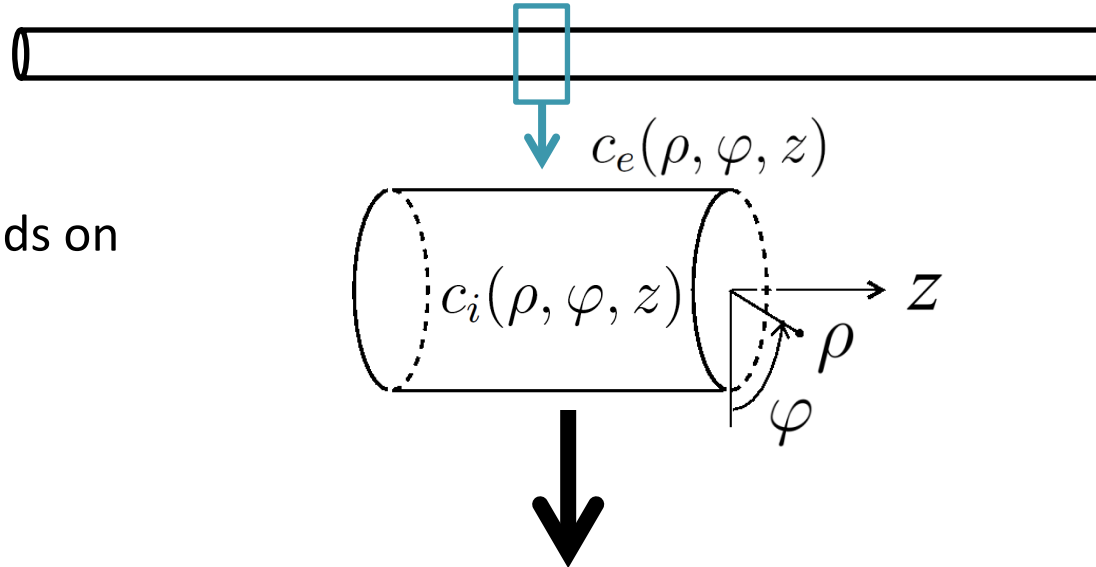
- Single fiber close to electrodes
- Single fiber far from electrodes
- Model variations and tissue-wide effect

Conclusions and Connections to LLNL

Derivation of Mass Transport Model

Model transport of molecules

Full 3D convection-diffusion equations



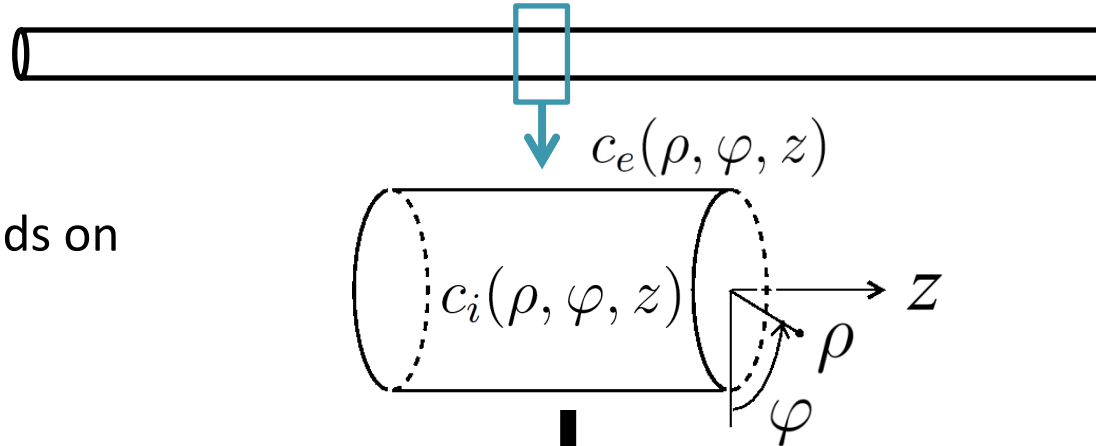
Transport depends on

- Electric drift
- Diffusion

Derivation of Mass Transport Model

Model transport of molecules

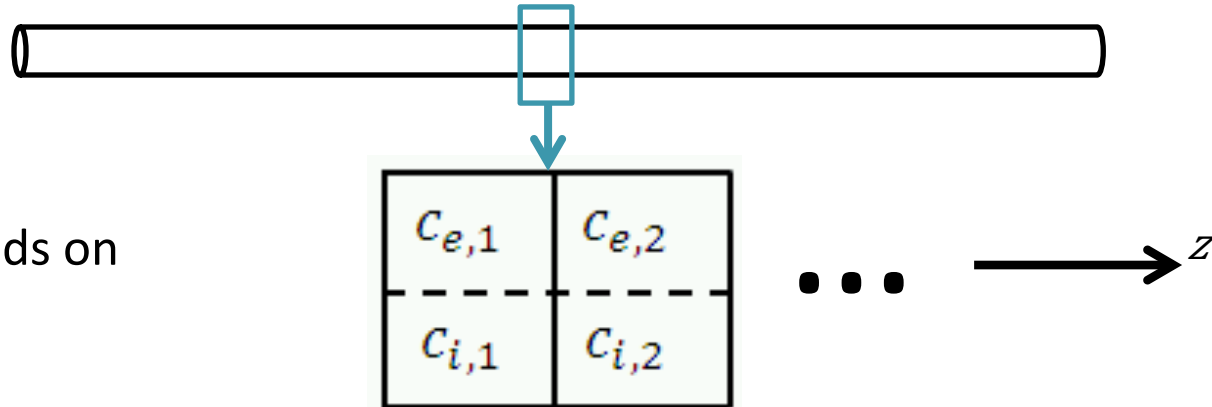
Full 3D convection-diffusion equations



Transport depends on

- Electric drift
- Diffusion

Series of longitudinally independent, two-compartment equations
with flux at membrane



Transport depends on

- Diffusion

Beginning of Mass Transport Model Deriv.

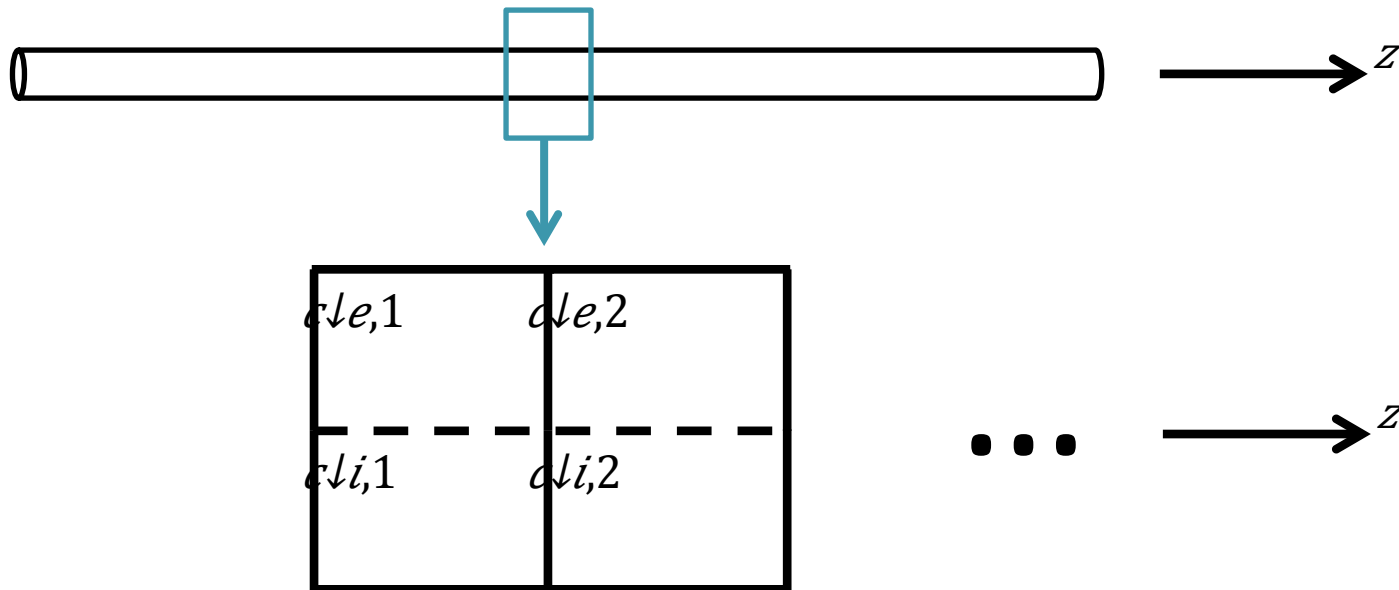
Derivation of Mass Transport Model

Model transport of molecules

Full 3D convection-diffusion equations



Series of longitudinally independent, two-compartment equations
with flux at membrane

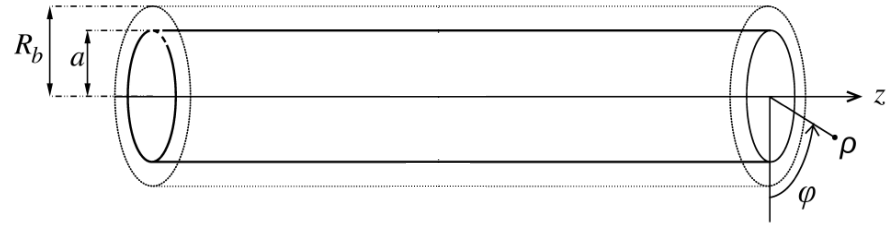


Derivation of Mass Transport Model

Convection-diffusion equation

$$\frac{\partial c_i}{\partial t} = \nabla \cdot \{D \nabla c_i - \mathbf{v}_i c_i\} \quad \text{in } \rho < a,$$

$$\frac{\partial c_e}{\partial t} = \nabla \cdot \{D \nabla c_e - \mathbf{v}_e c_e\} \quad \text{in } a < \rho < R_b$$



Electric drift velocity

$$\mathbf{v}_i = -D \frac{z_m e}{k_B T} \nabla V_i \quad \text{in } \rho < a,$$

$$\mathbf{v}_e = -D \frac{z_m e}{k_B T} \nabla V_e \quad \text{in } a < \rho < R_b$$

Continuity of molecular flux across membrane boundary

$$D \left\{ \frac{\partial c_i}{\partial \rho} + \frac{c_i z_m e}{k_B T} \frac{\partial V_i}{\partial \rho} \right\} = D \left\{ \frac{\partial c_e}{\partial \rho} + \frac{c_e z_m e}{k_B T} \frac{\partial V_e}{\partial \rho} \right\} = -j_m \quad \text{on } \rho = a$$

Derivation of Mass Transport Model

Molecule and pulsing parameters

Parameter	Definition	Units	Value
$ z_m $	Valence magnitude	Unitless	1
D	Diffusion coefficient	$\text{m}^2 \text{s}^{-1}$	2.01×10^{-10}
T	Temperature	K	310
τ_p	Pulse duration	μs	100
$\ E\ _\infty$	Max. magnitude of field	kV cm^{-1}	10
N	Number of pulses	Unitless	8
T_p	Period of pulse train	s	1
a	Fiber radius	μm	25

Derivation of Mass Transport Model

Molecule and pulsing parameters

Parameter	Definition	Units	Value
$ z_m $	Valence magnitude	Unitless	1
D	Diffusion coefficient	$\text{m}^2 \text{s}^{-1}$	2.01×10^{-10}
T	Temperature	K	310
τ_p	Pulse duration	μs	100
$\ E\ _\infty$	Max. magnitude of field	kV cm^{-1}	10
N	Number of pulses	Unitless	8
T_p	Period of pulse train	s	1
a	Fiber radius	μm	25

Two time scales

1. Electric drift: $N\tau_p = 800 \mu\text{s}$
2. Diffusion: 8.5 s

Derivation of Mass Transport Model

Molecule and pulsing parameters

Parameter	Definition	Units	Value
$ z_m $	Valence magnitude	Unitless	1
D	Diffusion coefficient	$\text{m}^2 \text{s}^{-1}$	2.01×10^{-10}
T	Temperature	K	310
τ_p	Pulse duration	μs	100
$\ E\ _\infty$	Max. magnitude of field	kV cm^{-1}	10
N	Number of pulses	Unitless	8
T_p	Period of pulse train	s	1
a	Fiber radius	μm	25

Two time scales

1. Electric drift: $N\tau_p = 800 \mu\text{s}$
2. Diffusion: 8.5 s

Two length scales

1. Electric drift: $L_e = \underbrace{\left(\frac{D|z_m|e}{k_B T} \|E\|_\infty \right)}_{\text{max. velocity}} 800 \mu\text{s} = 6.01 \mu\text{m}$
2. Diffusion: $L_d = \sqrt{6D8.5 \text{ s}} = 101 \mu\text{m}$

Derivation of Mass Transport Model

Molecule and pulsing parameters

Parameter	Definition	Units	Value
$ z_m $	Valence magnitude	Unitless	1
D	Diffusion coefficient	$\text{m}^2 \text{s}^{-1}$	2.01×10^{-10}
T	Temperature	K	310
τ_p	Pulse duration	μs	100
$\ E\ _\infty$	Max. magnitude of field	kV cm^{-1}	10
N	Number of pulses	Unitless	8
T_p	Period of pulse train	s	1
a	Fiber radius	μm	25

Two time scales

1. Electric drift: $N\tau_p = 800 \mu\text{s}$
2. Diffusion: 8.5 s

Two length scales

1. Electric drift: $L_e = \underbrace{\left(\frac{D|z_m|e}{k_B T} \|E\|_\infty \right)}_{\text{max. velocity}} 800 \mu\text{s} = 6.01 \mu\text{m}$
2. Diffusion: $L_d = \sqrt{6D8.5 \text{ s}} = 101 \mu\text{m}$

Two assumptions

1. Electric drift ignored
(L_e is 6% of L_d)
2. “Well-mixed” in transverse direction
(L_d is four-fold larger than fiber radius a)

Derivation of Mass Transport Model

Reduction to two-compartment, 1D longitudinal diffusion equations

Divergence theorem

- Neglect electric drift
- Uniform transverse concentration

$$\frac{\partial c_i}{\partial t} = D \frac{\partial^2 c_i}{\partial z^2} - \frac{2}{a} \overline{j_m},$$

$$\frac{\partial c_e}{\partial t} = D \frac{\partial^2 c_e}{\partial z^2} + \frac{2a}{R_b^2 - a^2} \overline{j_m}$$

$\overline{j_m}$ is instantaneous function of c_i , c_e , and N_{ep}

Derivation of Mass Transport Model

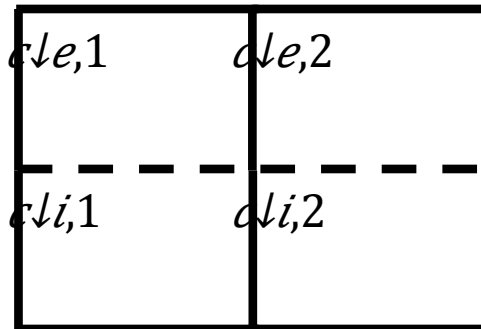
Remove longitudinal diffusion

$$\frac{\partial c_i}{\partial t} = \cancel{D \frac{\partial^2 c_i}{\partial z^2}} - \frac{2}{a} \dot{j}_m,$$

$$\frac{\partial c_e}{\partial t} = \cancel{D \frac{\partial^2 c_e}{\partial z^2}} + \frac{2a}{R_b^2 - a^2} \dot{j}_m$$

Maximum error in total uptake in entire fiber from neglecting longitudinal diffusion is 0.1%

$$\begin{aligned} \frac{\partial c_i}{\partial t} &= -\frac{2}{a} \dot{j}_m, \\ \frac{\partial c_e}{\partial t} &= \frac{2a}{R_b^2 - a^2} \dot{j}_m \end{aligned}$$



End of Mass Transport Model Deriv.

Outline

Introduction

- Background
- Published experiments and models

Derivation and Validation of Models

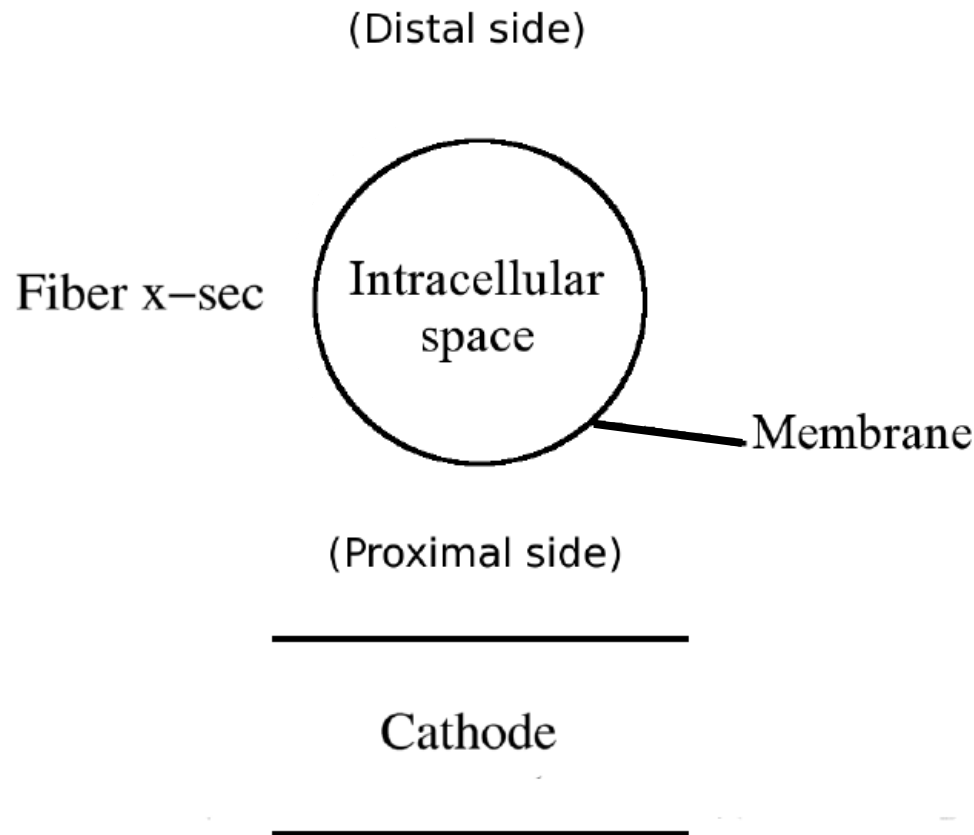
- Asymptotic fiber model
- Alterations for EP-mediated uptake
- Mass transport model

Results

- **Single fiber close to electrodes**
- Single fiber far from electrodes
- Model variations and tissue-wide effect

Conclusions and Connections to LLNL

Passive Versus Electroporating

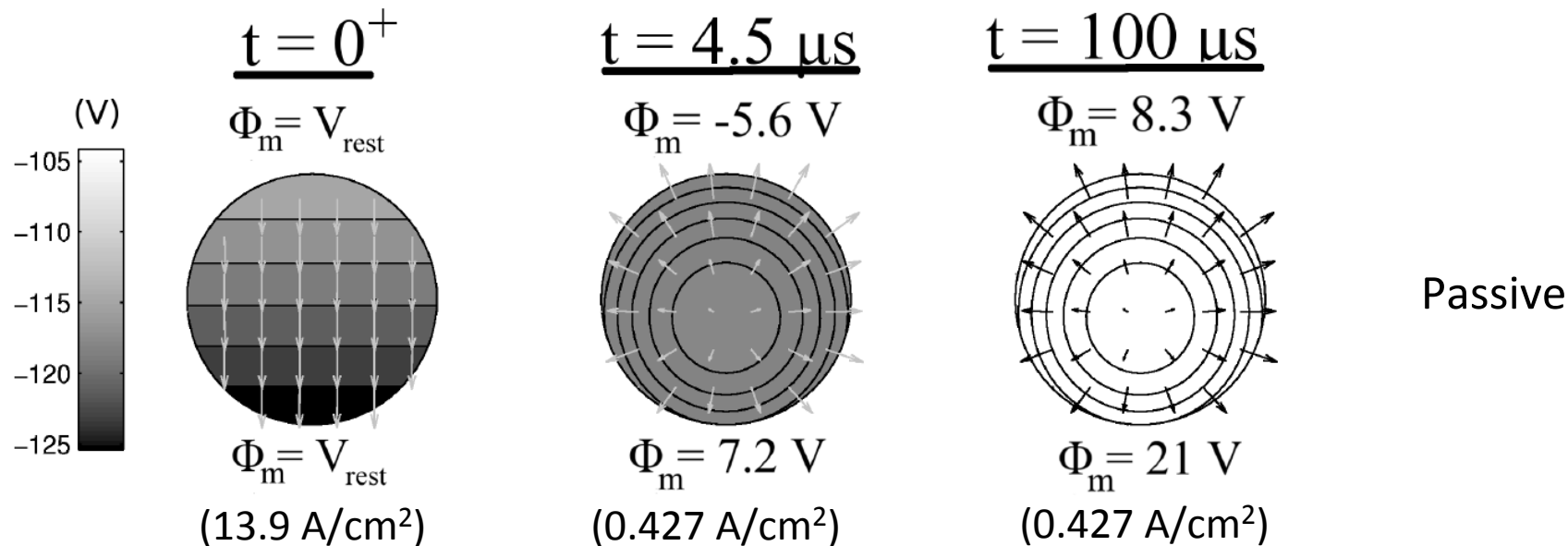


Contour plots viewed at 3 time instants

1. $t=0^+$, when pulse turned on
2. $t=4.5 \mu\text{s}$, ten-fold larger than time scale of membrane charging via transverse currents
3. $t=100 \mu\text{s}$, pulse duration

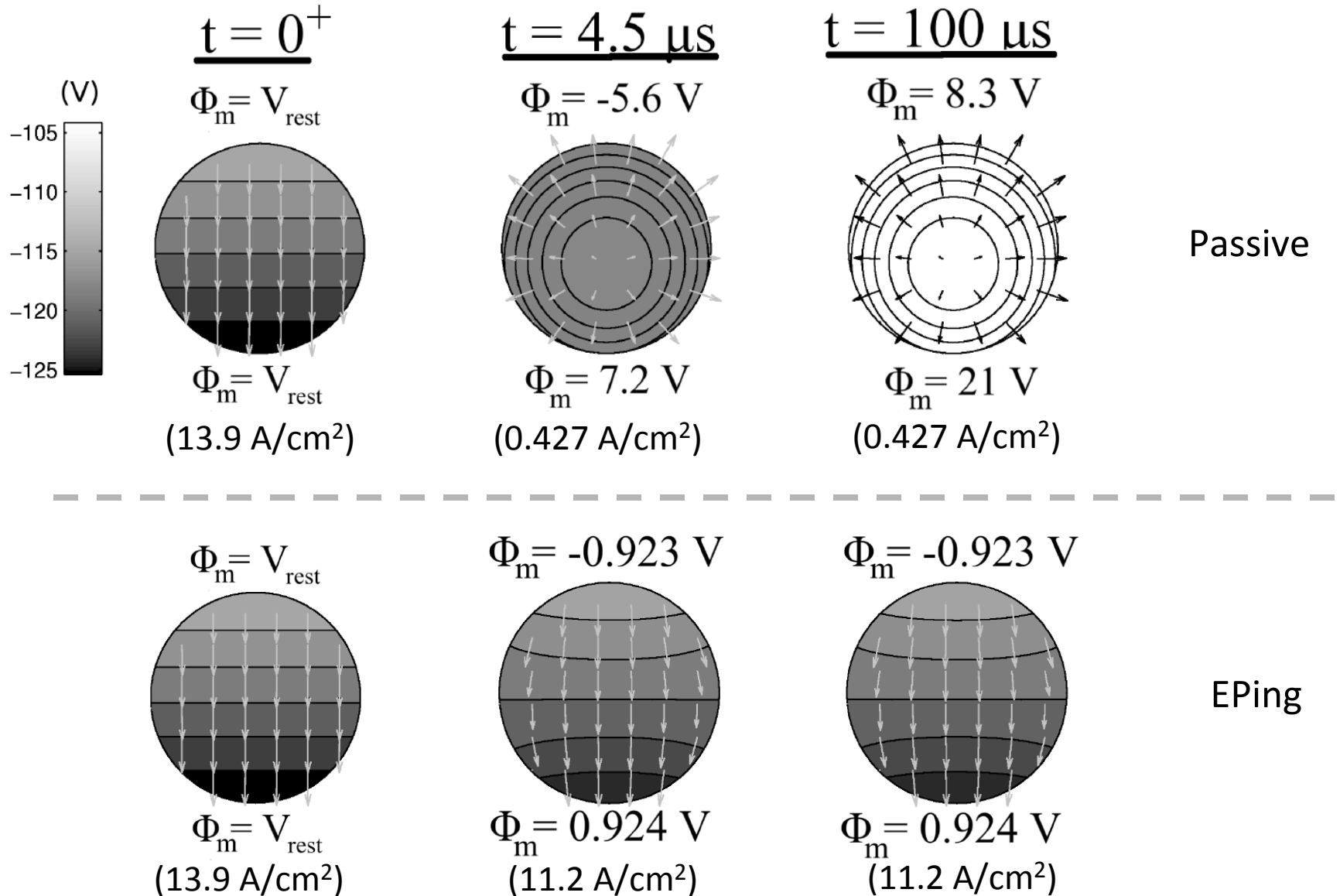
Passive Versus Electroporating

Distribution of potential and current density



Passive Versus Electroporating

Distribution of potential and current density



Passive Versus Electroporating

1D longitudinal problem for mean potential

$$\left(1 + \frac{1}{\gamma\mu}\right) C_m \frac{\partial f_m^0}{\partial t} = \left(1 + \frac{1}{\gamma}\right) \frac{\sigma_i a}{2} \frac{\partial^2 \langle \Psi \rangle}{\partial z^2} - \left(1 + \frac{1}{\gamma\mu}\right) \overline{I_m}$$

Passive Versus Electroporating

1D longitudinal problem for mean potential

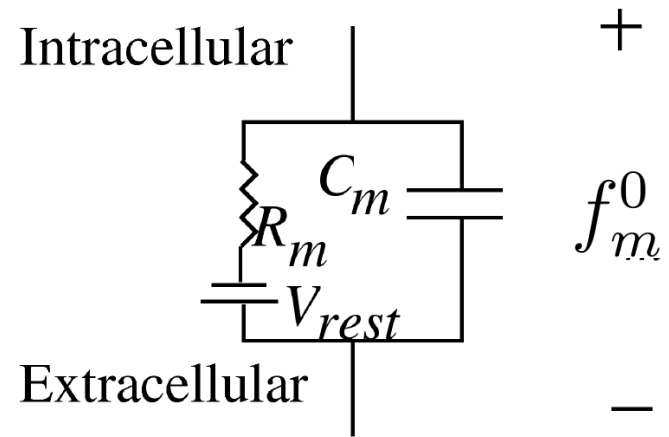
$$\left(1 + \frac{1}{\gamma\mu}\right) C_m \frac{\partial f_m^0}{\partial t} = \left(1 + \frac{1}{\gamma}\right) \frac{\sigma_i a}{2} \frac{\partial^2 \langle \Psi \rangle}{\partial z^2} - \left(1 + \frac{1}{\gamma\mu}\right) \overline{I}_m$$

Passive fiber

$$\overline{I}_m = \left(\frac{1}{R_m}\right) f_m^0 - \frac{V_{rest}}{R_m}$$



Time scale charging: 10 milliseconds



Passive Versus Electroporating

1D longitudinal problem for mean potential

$$\left(1 + \frac{1}{\gamma\mu}\right) C_m \frac{\partial f_m^0}{\partial t} = \left(1 + \frac{1}{\gamma}\right) \frac{\sigma_i a}{2} \frac{\partial^2 \langle \Psi \rangle}{\partial z^2} - \left(1 + \frac{1}{\gamma\mu}\right) \overline{I}_m$$

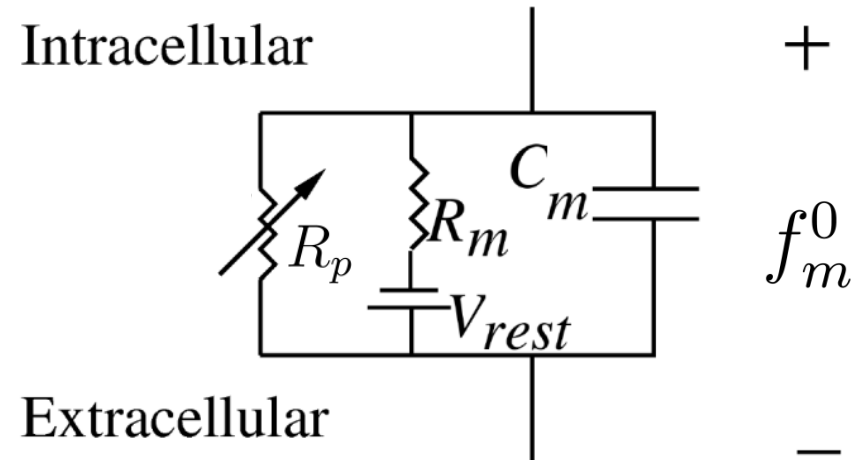
Electroporating fiber

$$\overline{I}_m = \left(\frac{R_m + R_p}{R_m R_p} \right) f_m^0 - \frac{V_{rest}}{R_m}$$

- Choose typical value of R_p that develops on short time scale due to transverse charging



Time scale of charging: 1 microsecond



Passive Versus Electroporating

1D longitudinal problem for mean potential

$$\left(1 + \frac{1}{\gamma\mu}\right) C_m \frac{\partial f_m^0}{\partial t} = \left(1 + \frac{1}{\gamma}\right) \frac{\sigma_i a}{2} \frac{\partial^2 \langle \Psi \rangle}{\partial z^2} - \left(1 + \frac{1}{\gamma\mu}\right) \overline{I}_m$$

Passive fiber

10 milliseconds



Membrane continues
charging after 4.5 μs

Electroporating fiber

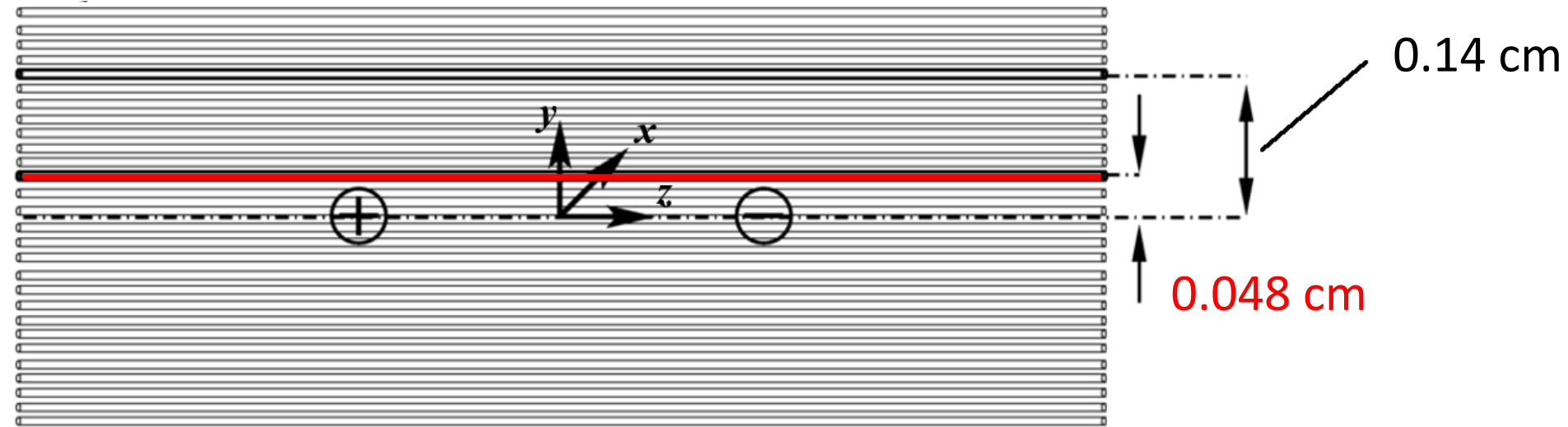
1 microseconds



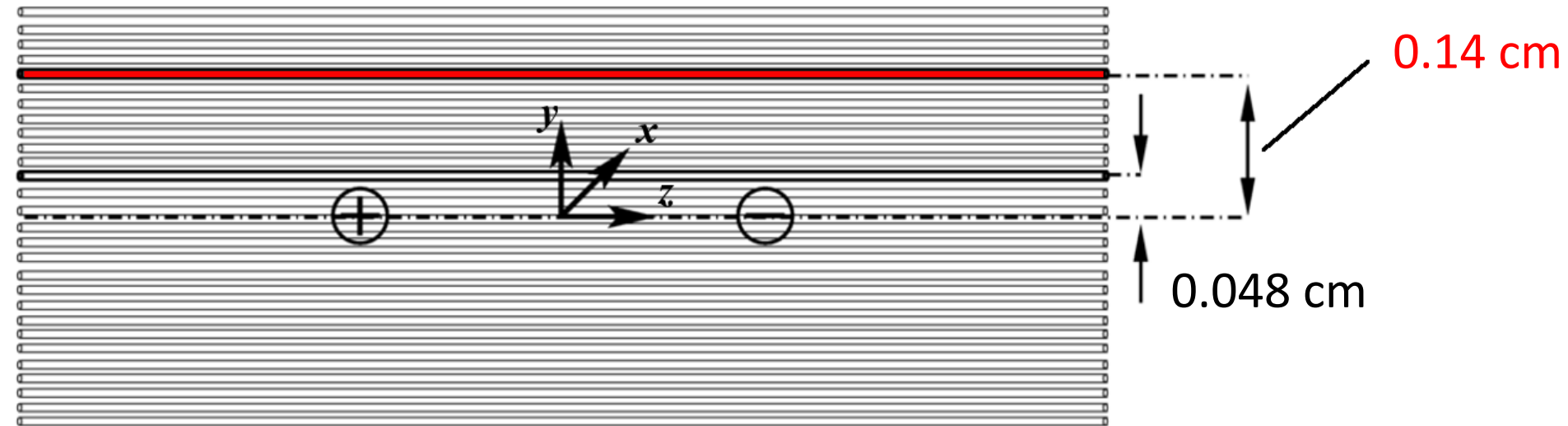
Membrane stops
charging and EPing
after 4.5 μs

**If EP occurs on time scale transverse charging (4.5 ,
it precludes further charging and pore creation in longitudinal problem**

Fiber Close to Electrodes

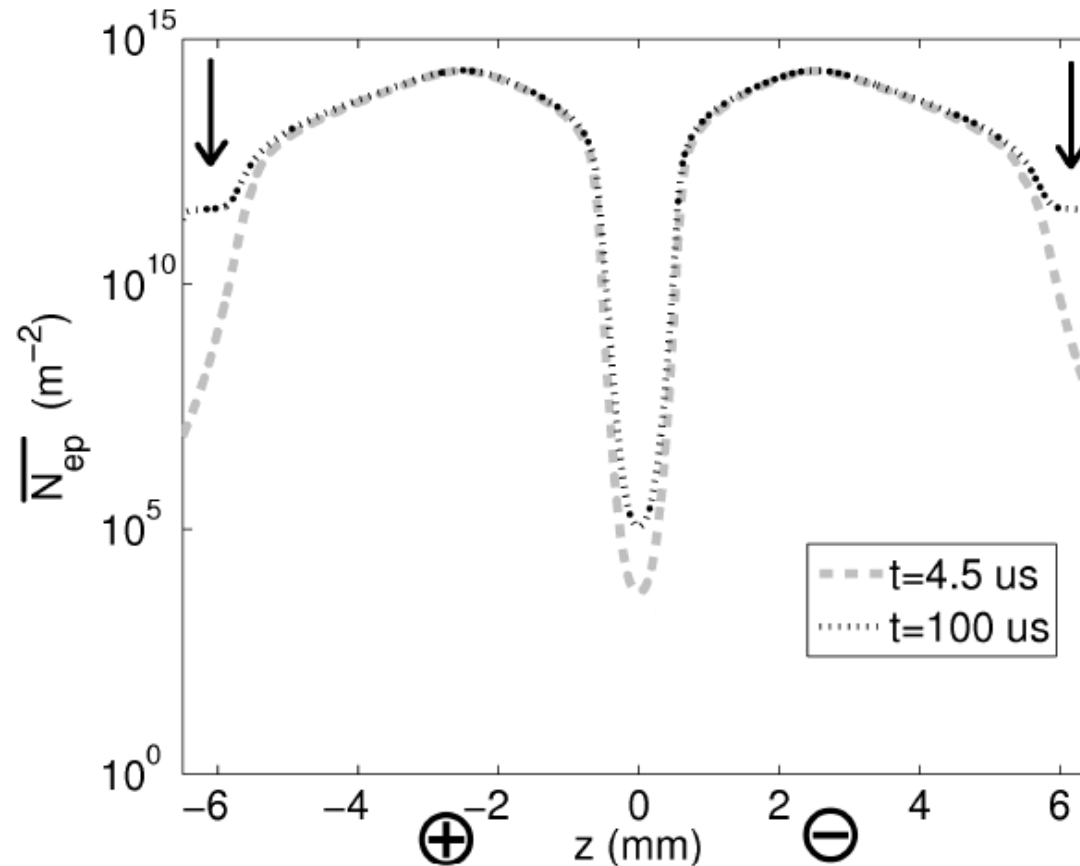


Fiber Far From Electrodes



Electroporating Fiber

Does charging and EPing cease after $4.5\mu s$ over the entire fiber?

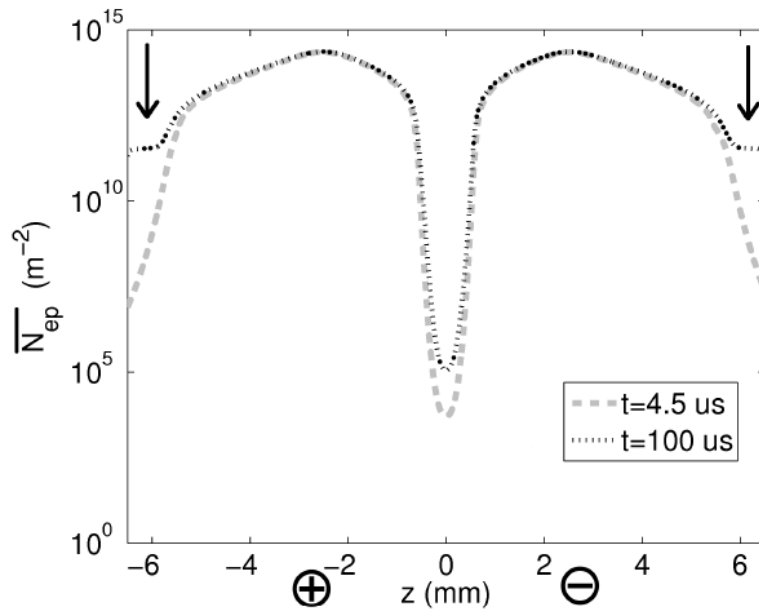


Membrane charging and additional pore creation after $4.5 \mu s$ for $|z| > 6$ mm

Effect of Longitudinal Charging

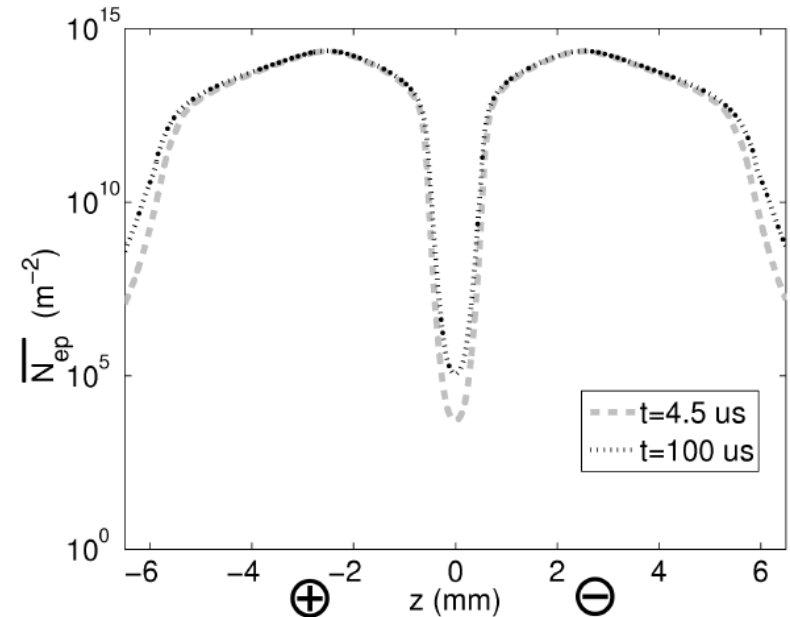
Creation of pores past 4.5 μs from transverse or longitudinal charging?

Pore Density Distribution,
Full Model



Charging and addition pore creation
after 4.5 μs for $|z| > 6$ mm

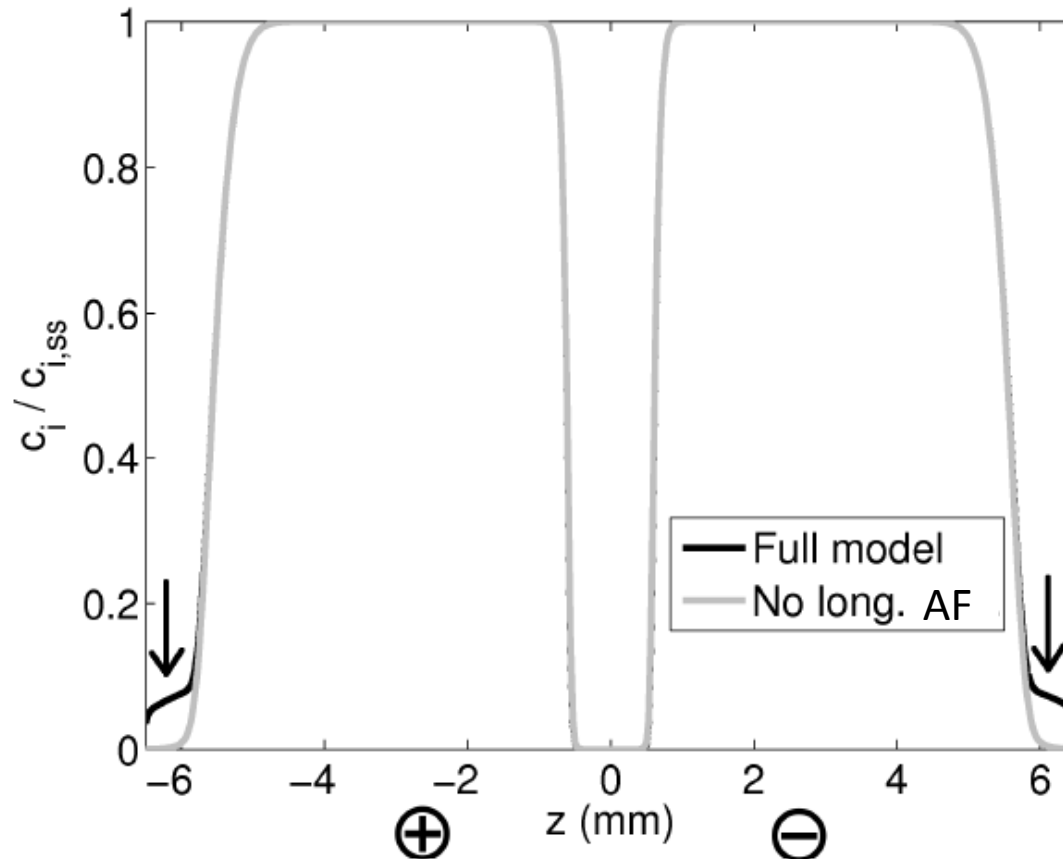
Pore Density Distribution
No Long. AF



No charging and addition pore
creation after 4.5 μs for $|z| > 6$ mm

Effect of Longitudinal Charging

Effect of longitudinal charging on uptake



Longitudinal AF effects uptake only for $|z| > 6$ mm

Difference in total uptake over entire fiber is 0.91 %

Outline

Introduction

- Background
- Published experiments and models

Derivation and Validation of Models

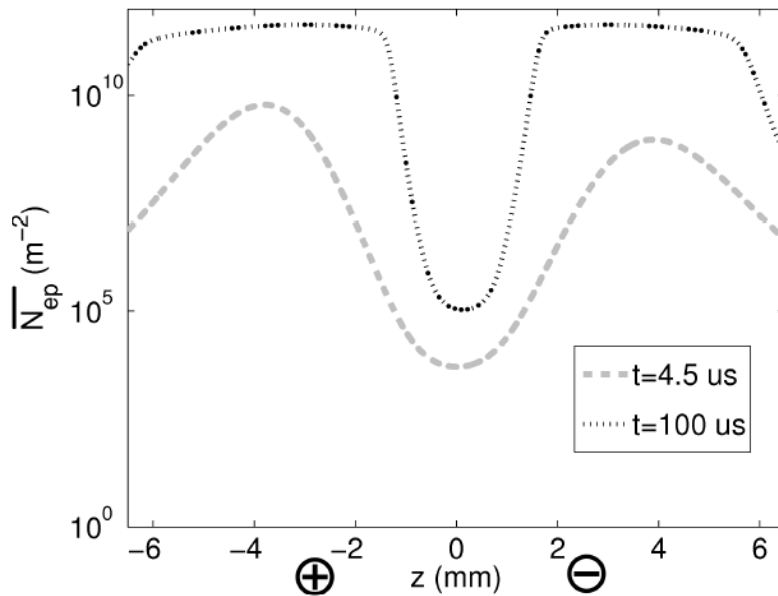
- Asymptotic fiber model
- Alterations for EP-mediated uptake
- Mass transport model

Results

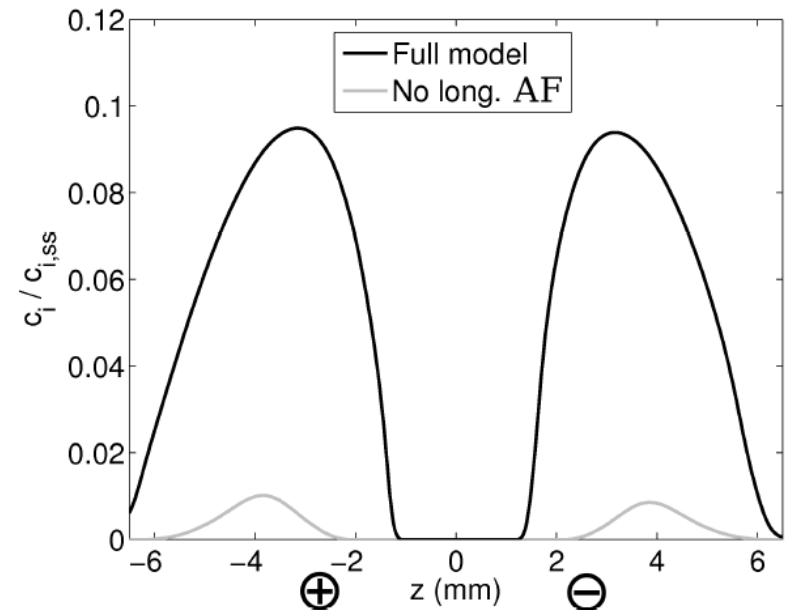
- Single fiber close to electrodes
- **Single fiber far from electrodes**
- Model variations and tissue-wide effect

Conclusions and Connections to LLNL

Effect of Longitudinal Charging



Membrane charging
and additional pore
creation after $4.5 \mu s$ for
all locations



Longitudinal AF effects
uptake for all positions

**Difference in total uptake
over entire fiber is 2000%**

Passive Versus Electroporating

1D longitudinal problem for mean potential

$$\left(1 + \frac{1}{\gamma\mu}\right) C_m \frac{\partial f_m^0}{\partial t} = \frac{\sigma_i a}{2} \frac{\partial^2 f_m^0}{\partial z^2} + \left(1 + \frac{1}{\gamma}\right) \frac{\sigma_i a}{2} \frac{\partial^2 \langle \Psi \rangle}{\partial z^2} - \left(1 + \frac{1}{\gamma\mu}\right) \overline{I}_m$$

Effect of Longitudinal Charging

Creation of pores past 4.5 μ s from transverse or longitudinal charging?

Simulate:

- Full transverse problem
- Full longitudinal problem

$$\left(1 + \frac{1}{\gamma\mu}\right) C_m \frac{\partial f_m^0}{\partial t} = \frac{\sigma_i a}{2} \frac{\partial^2 f_m^0}{\partial z^2} + \left(1 + \frac{1}{\gamma}\right) \frac{\sigma_i a}{2} \frac{\partial^2 \langle \Psi \rangle}{\partial z^2} - \left(1 + \frac{1}{\gamma\mu}\right) \overline{I_m}$$

Simulate:

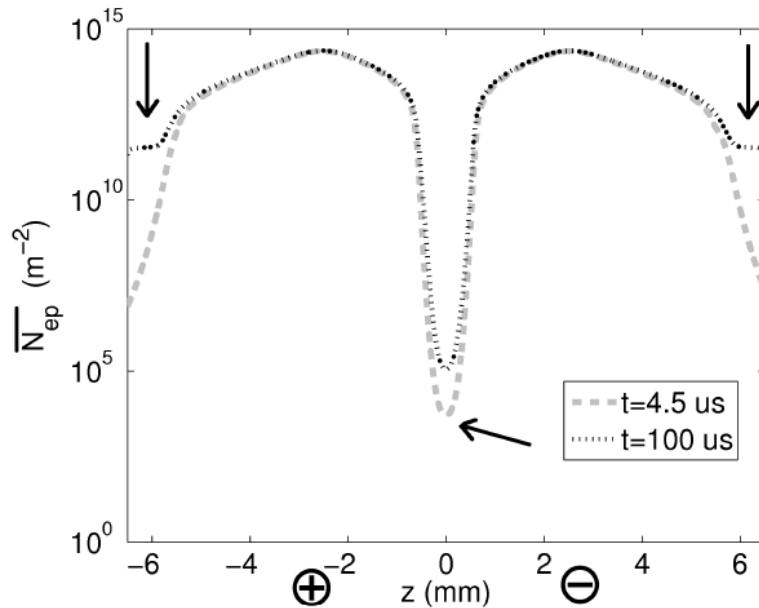
- Full transverse problem
- Partial longitudinal problem – no longitudinal AF

$$C_m \frac{\partial f_m^0}{\partial t} = -\overline{I_m}$$

Effect of Longitudinal Charging

Creation of pores past 4.5 μs from transverse or longitudinal charging?

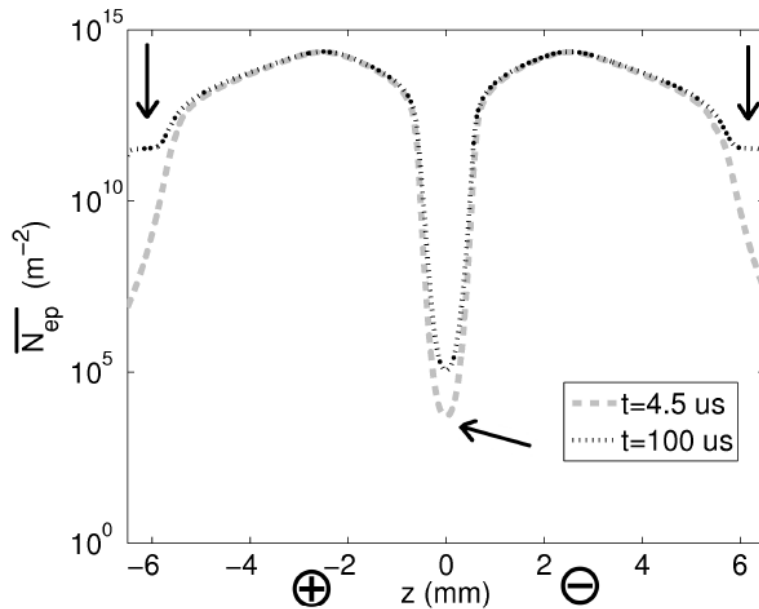
Pore Density Distribution,
Full Model



Effect of Longitudinal Charging

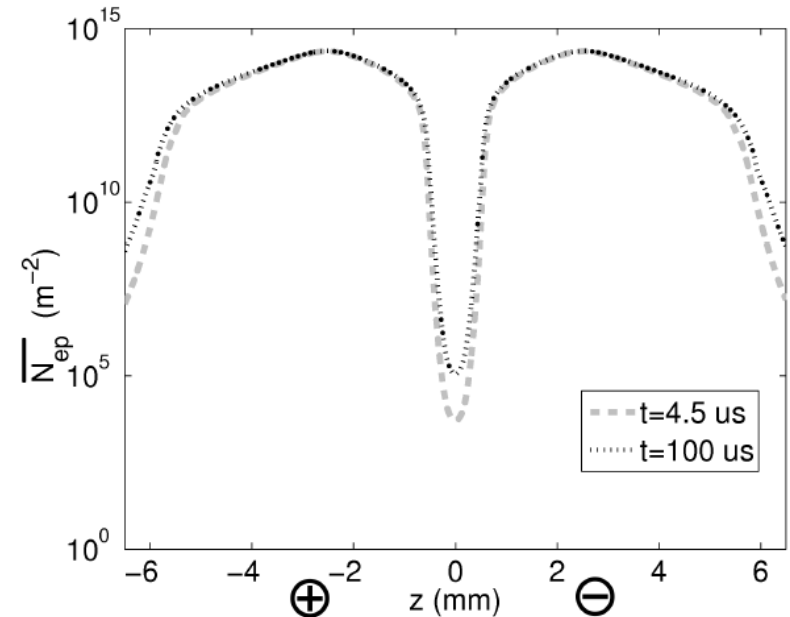
Creation of pores past 4.5 μs from transverse or longitudinal charging?

Pore Density Distribution,
Full Model



Charging and addition pore creation
after 4.5 μs for $|z| > 6$ mm

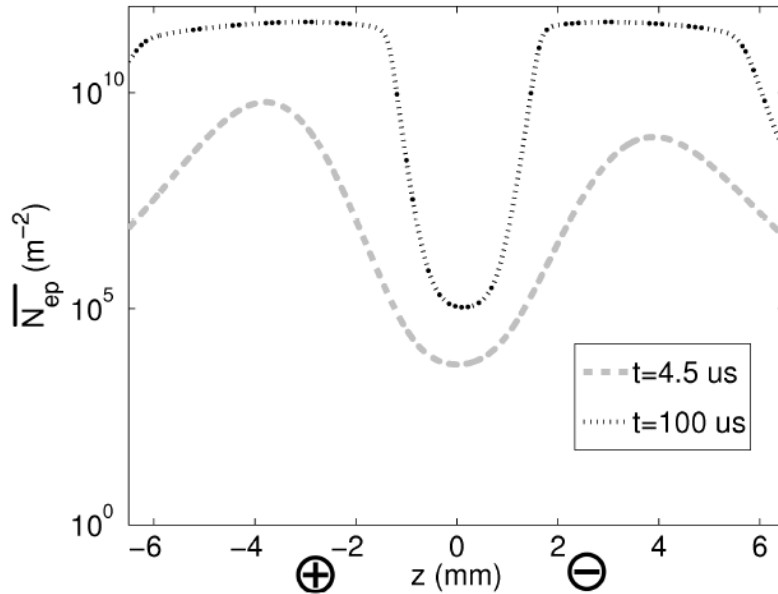
Pore Density Distribution
No Long. AF



No charging and addition pore
creation after 4.5 μs for $|z| > 6$ mm

Effect of Longitudinal Charging

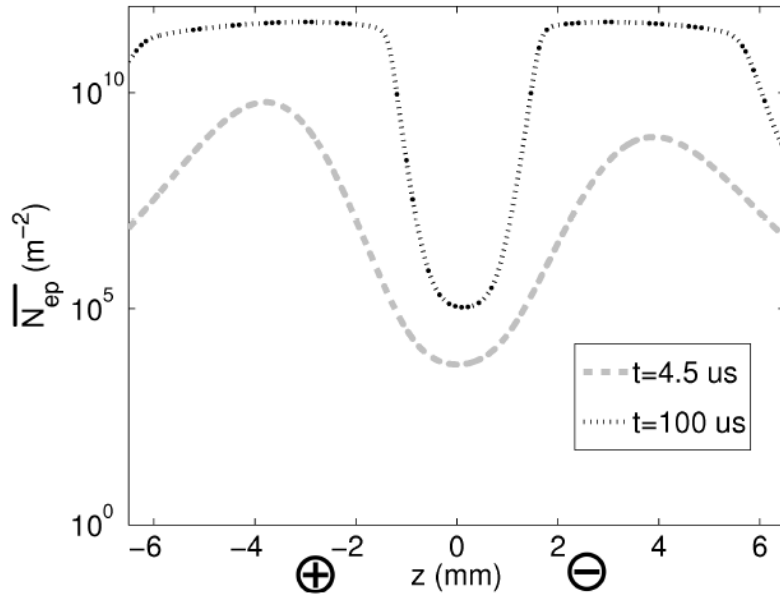
Pore Density Distribution, Full Model



Membrane charging
and additional pore
creation after $4.5 \mu\text{s}$
for all locations

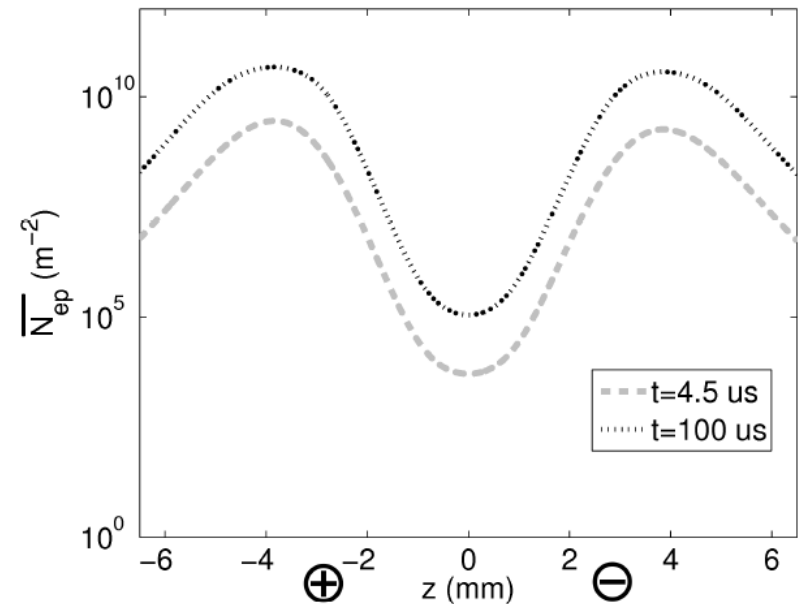
Effect of Longitudinal Charging

Pore Density Distribution, Full Model



Membrane charging
and additional pore
creation after $4.5 \mu\text{s}$
for all locations

Pore Density Distribution No Long. AF



Membrane charging
and additional pore
creation after $4.5 \mu\text{s}$
for all locations, but \overline{N}_{ep}
at peaks is one order
mag. smaller

Outline

Introduction

- Background
- Published experiments and models

Derivation and Validation of Models

- Asymptotic fiber model
- Alterations for EP-mediated uptake
- Mass transport model

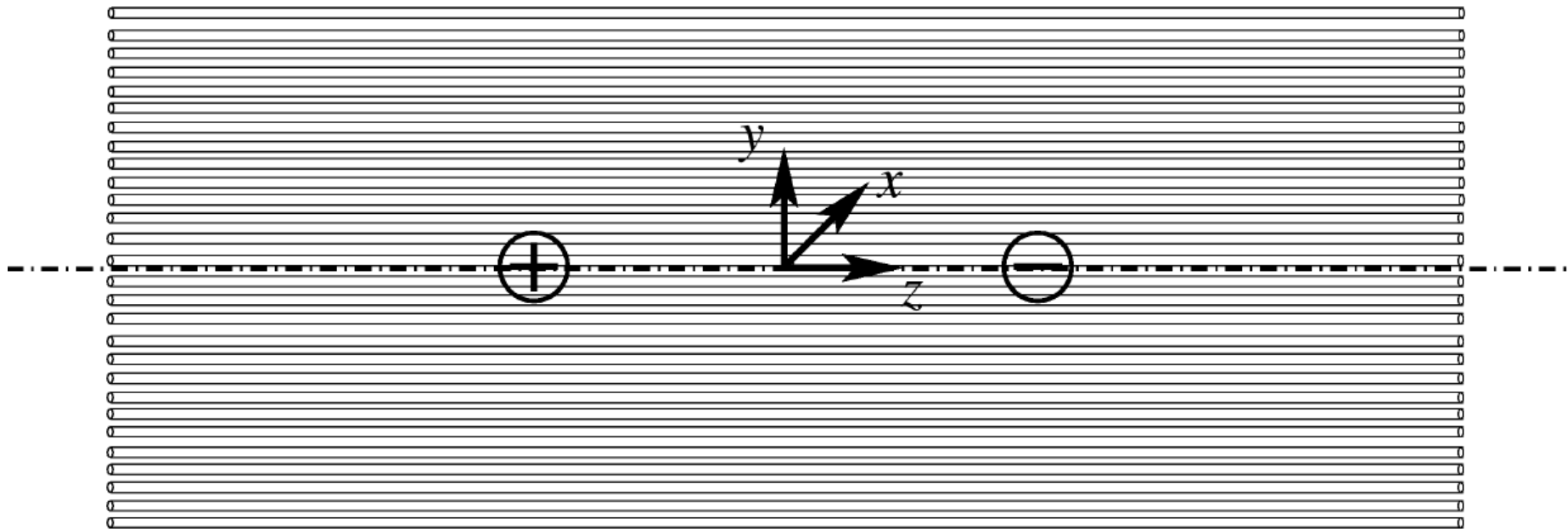
Results

- Single fiber close to electrodes
- Single fiber far from electrodes
- **Model variations and tissue-wide effect**

Conclusions and Connections to LLNL

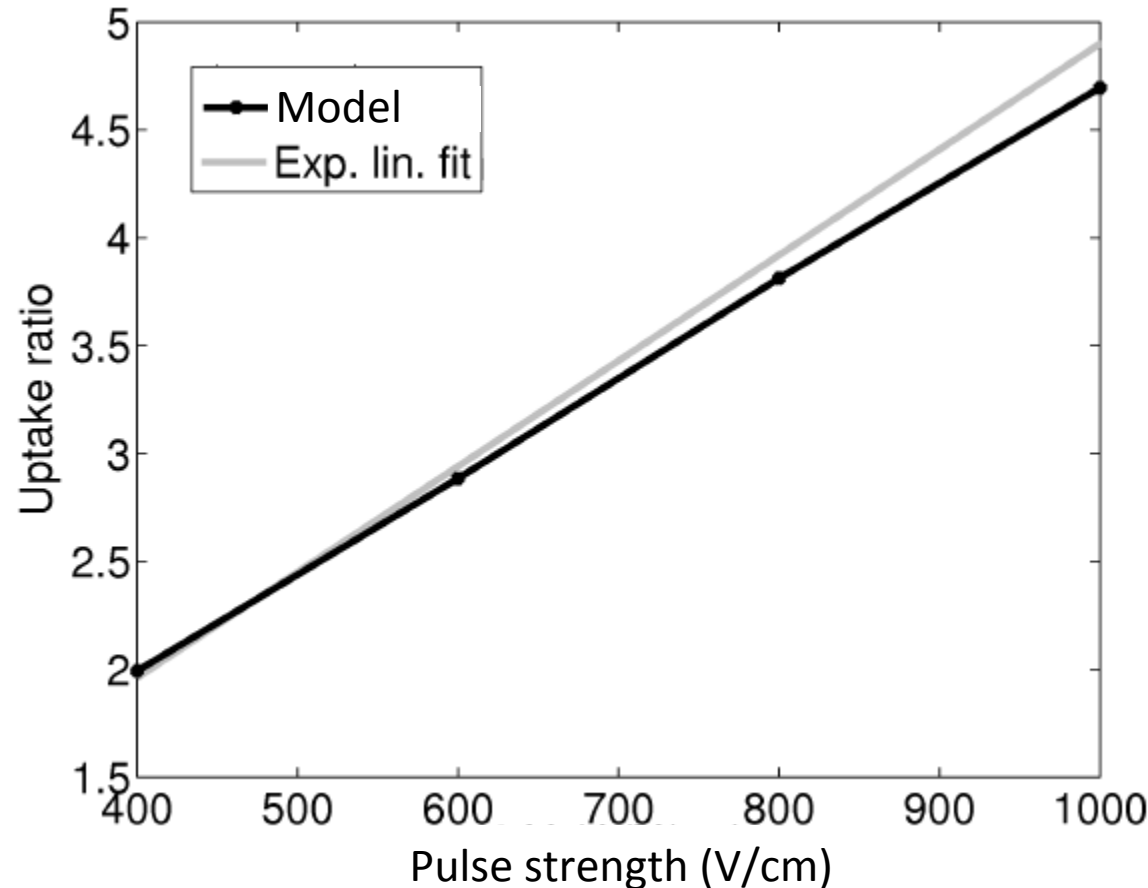
Model Variations

Uptake over entire tissue



Model Variations

Pulse strength: comparison to experiments¹



Relative difference in uptake never exceeds 5%

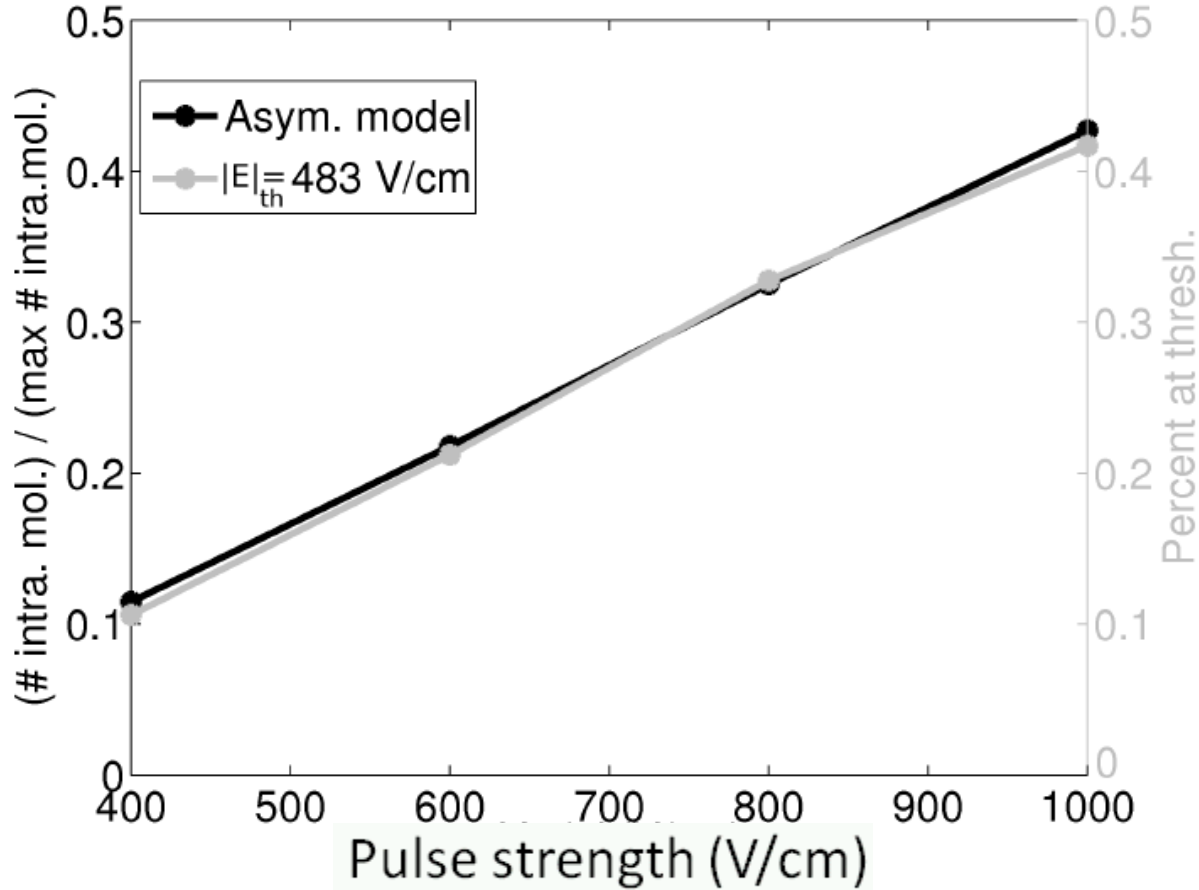
Experimental variability: 6%

Maximum contribution from longitudinal AF in this range of pulse strengths: 10%

¹Grafstrom, G., Engstrom, P., Salford, L.G., Persson, B.R.R. (2006) *Cancer Biother. Radio.*, 21,

Model Variations

Pulse strength: comparison to threshold field model

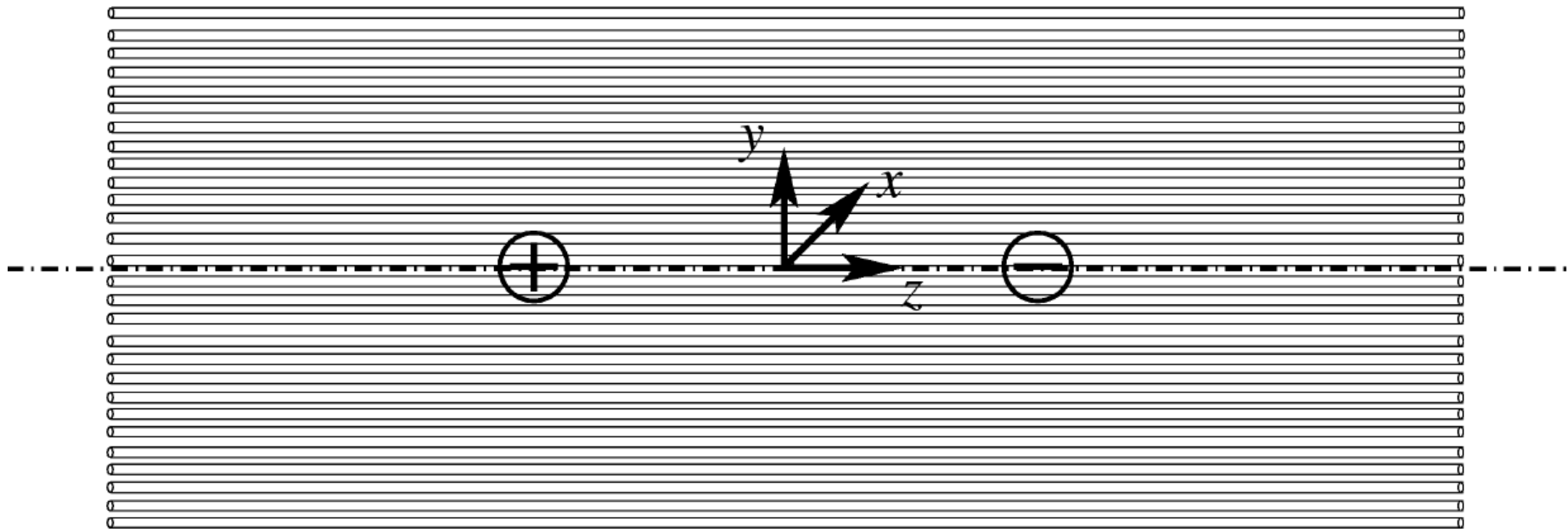


Average relative difference is 3.31%

Maximum relative difference is 7.45% at pulse strength 400 V/cm

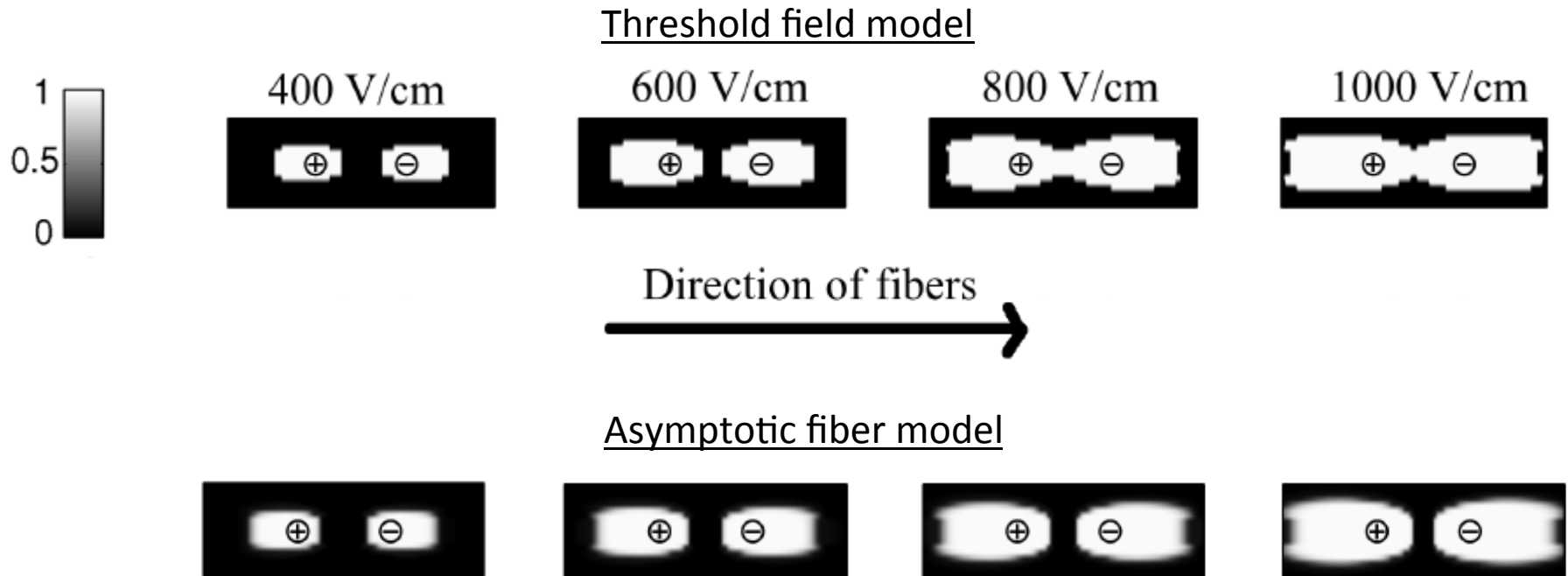
Model Variations

Muscle Tissue View



Model Variations

Comparison to threshold field model



Difference in uptake between electrodes

Relative root-mean-square difference around 20%

Correlation coefficient at most 0.969

Outline

Introduction

- Background
- Published experiments and models

Derivation and Validation of Models

- Asymptotic fiber model
- Alterations for EP-mediated uptake
- Mass transport model

Results

- Single fiber close to electrodes
- Single fiber far from electrodes
- Model variations and tissue-wide effect

Conclusions and Connections to LLNL

Conclusions

Fiber close to electrodes

1. Significant EP occurs on the time scale of transverse charging
2. Charging of the membrane and creation of pores from longitudinal AF is impeded

Fiber far from electrodes

1. Minimal EP occurs on time scale of transverse charging
2. Majority of membrane charging and creation of pores attributed to longitudinal AF

Uptake over entire tissue

- Contribution of longitudinal AF never more than 10%
- 10% on order of experimental variability
- Corroborates previous models in terms of total uptake, but not distribution

Reproduce tissue uptake experiments qualitatively and quantitatively

Model Variations

Pulse strength: comparison to threshold field model

- Other models base electroporation off of field magnitude $|E|$
- Assume electroporating when $|E| > |E|_{th}$
- Formulate tissue conductivity, and corresponding uptake, based on choice of $|E|_{th}$

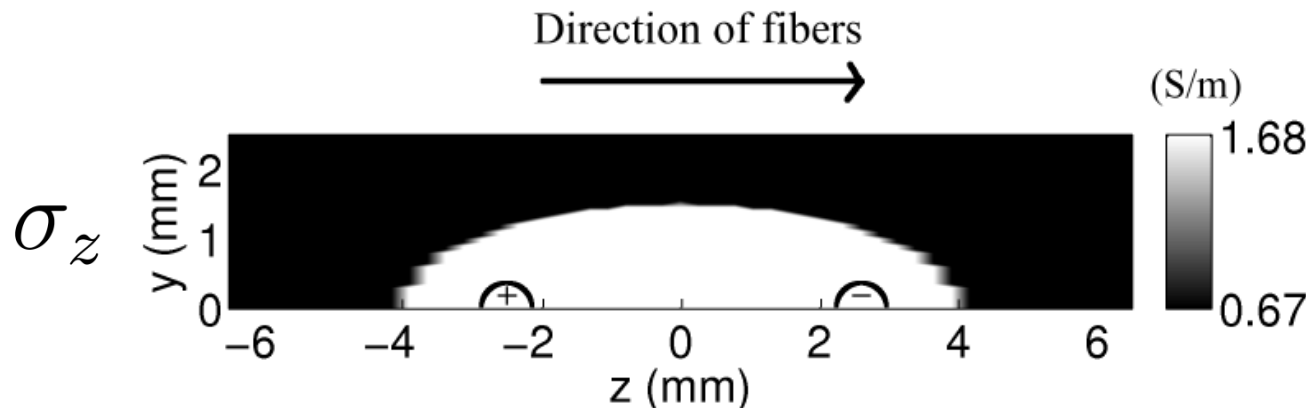
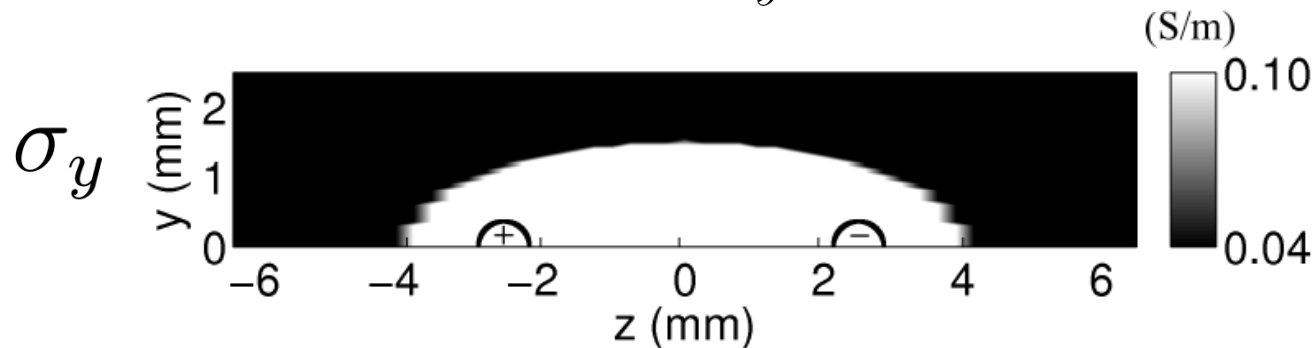
Model Variations

Effect of electroporation on tissue conductivities

Old way: Uniform macroscopic tissue conductivities

New way: Nonuniform macroscopic tissue conductivities

$$\sigma_y \frac{\partial^2 \Psi}{\partial y^2} + \sigma_z \frac{\partial^2 \Psi}{\partial z^2} = 0$$



^aSel, D., Cukjati, D. Batiuskaite, D., Slivnik, T., Mir, L. M. Miklavcic, D. (2005) IEEE T. Bio-Med. Eng., 52, 816-827

^bCorovic, S., Lackovic, I., Sustaric, P., Sustar, T., Rodic, T., Miklavcic, D. (2013) Biomed. Eng. Online, 12:16

Model Variations

Effect of electroporation on tissue conductivities

Simulate two versions solving Ψ numerically

1. Uniform tissue conductivities
2. Nonuniform tissue conductivities

Find numerical transverse AF $\partial \tilde{\Psi} / \partial \rho$ and longitudinal AF $\partial^2 \langle \Psi \rangle / \partial z^2$

Nonuniform version

3.51% increase in uptake
from longitudinal AF

Uniform version

3.77% increase in uptake
from longitudinal AF

Neglecting effect of EP on tissue conductivities does not change
qualitative role of longitudinal AF

Model Variations

Effect of electroporation on tissue conductivities

Total uptake: nonuniform version predicts 8.3% more uptake than uniform version

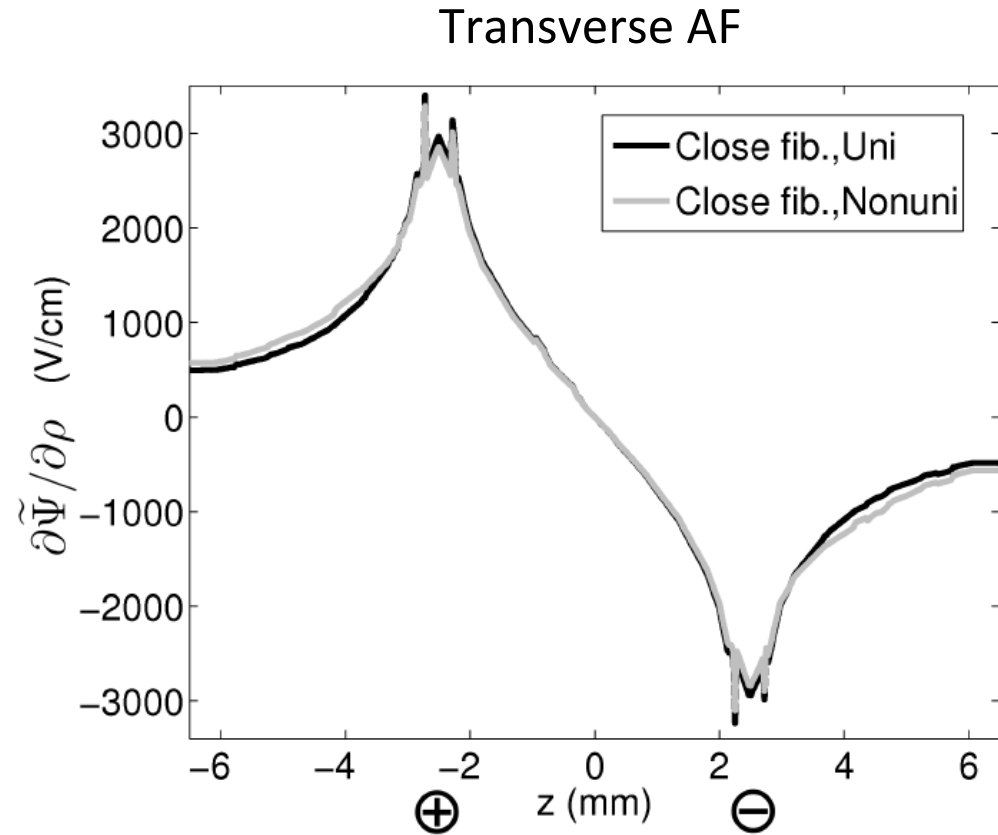
Model Variations

Effect of electroporation on tissue conductivities

Total uptake: nonuniform version predicts 8.3% more uptake than uniform version

Mechanism

- Both AFs larger in magnitude near ends of fiber at $z = \pm 6$ mm
 - Bot AFs smaller in magnitude near electrodes
- “Smears out” AFs



Model Variations

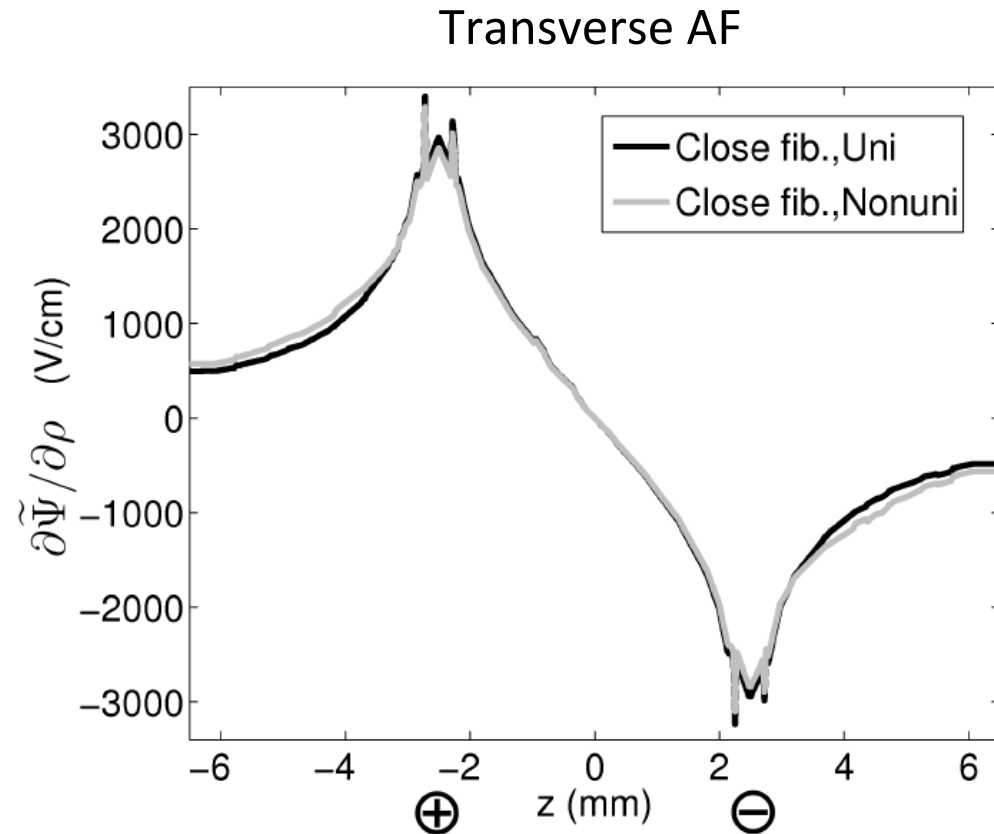
Effect of electroporation on tissue conductivities

Total uptake: nonuniform version predicts 8.3% more uptake than uniform version

Mechanism

- Both AFs larger in magnitude near ends of fiber at $z = \pm 6$ mm
 - Both AFs smaller in magnitude near electrodes
- “Smears out” AFs

30% difference in total tissue uptake using numerical AFs versus analytical AFs



Conclusions

Increase uptake when magnitude of AFs more evenly distributed over tissue

Reproduce tissue uptake experiments qualitatively and quantitatively

Importance of longitudinal AF not sensitive to changes in macroscopic tissue conductivities

Computationally efficient tool for 3D simulation

- New transverse AF
- Transverse and longitudinal electrical properties nerve stimulation

Limitations and Future Directions

Electroporation parameters and dynamics

- Artificial lipid bilayer parameters underestimates EP
- Growing pores

Periodicity of fibers

- Fiber bending
- Between tendons
- Site of electrodes

Limitations and Future Directions

Changes in macroscopic tissue conductivities due to EP

- Numerically solve additional 3D elliptic equation every time step
- Numerical error in AFs

Growing pores and DNA

- Large pores occur on “edges” of electroporated region
- Greater effect from longitudinal AF
 - Mechanism seen near ends of fiber
 - DNA can only pass through large pores

Einstein Relation

$$D = uRT / |z_m| F = u(N_A k_B) T / |z_m| (e N_A) = u k_B T / |z_m| e$$

Electrical mobility is velocity/field, here for intra. space: $u = \mathbf{v} / E$

$$D = \mathbf{v} k_B T / E |z_m| e$$

(don't need $|z_m|$ anymore b/c define it so always positive as v/E)

$$\mathbf{v} = D |z_m| e / k_B T E = -D |z_m| e / k_B T \nabla V$$

(velocity due to electric drift)

Stokes-Einstein Equation Diffusion Coefficient

$D = k_B T / 6\pi\mu a$, μ is viscosity, and a is radius of molecule approx. as spherical

$a = (3 MW / 4\pi\rho N_A)^{1/3}$, where MW is molecular weight, ρ is density, N_A is Avogadro's #

$$a = (3 * 400 \text{ g/mol} / 4\pi 9.932 \times 10^{15} \text{ g/m}^3 * 6.022 \times 10^{23} / \text{mol})^{1/3} = 5.42 \times 10^{-10} \text{ m}$$

$$D = 1.38 \times 10^{-23} \text{ J/K} * 310 \text{ K} / 6\pi * 0.6965 \times 10^{-3} \text{ N s / m}^2 * 5.42597 \times 10^{-10} \text{ m} = 6.0054 \times 10^{-10}$$

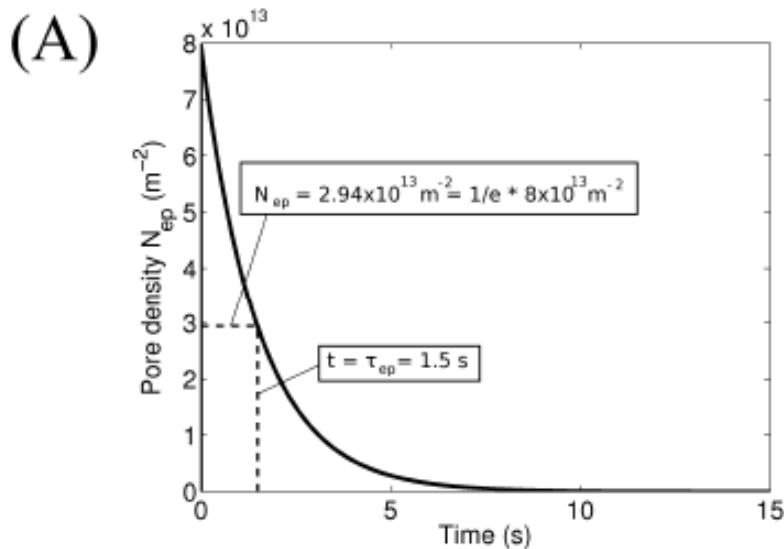
D/3 b/c molecule not in water, it is in intra/extra space, shown to be 1/3 diffusion as in water: $2.01 \times 10^{-10} \text{ m}^2/\text{s}$

Time Scale of Diffusion

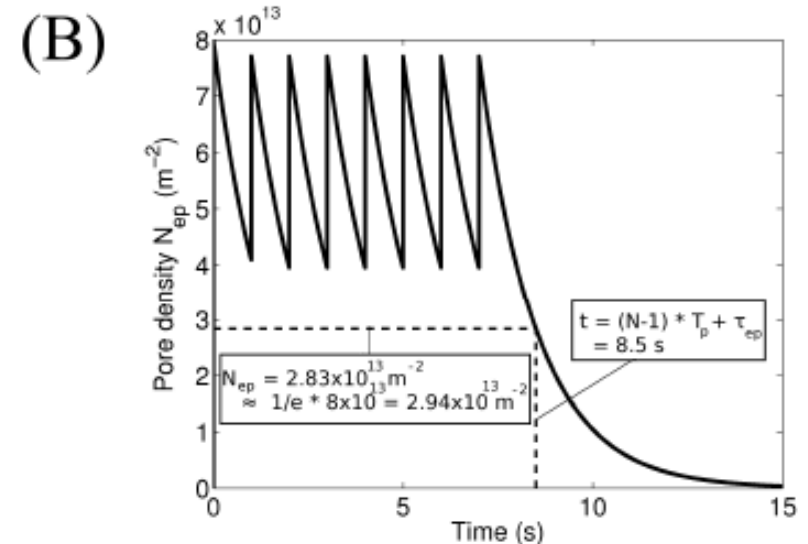
$$\frac{dN_{ep}}{dt} = \alpha e^{(\Phi_m/V_{ep})^2} \left(1 - \frac{N}{N_0 e^{q(\Phi_m/V_{ep})^2}} \right)$$

$$\tau_{ep} = \frac{N_0}{\alpha} e^{(q-1)(\Phi_m/V_{ep})^2} \Big|_{\Phi_m=0} = 1.5 \text{ s}$$

Pore Density
Single Pulse



Pore Density
Multiple Pulses



Radioactivity

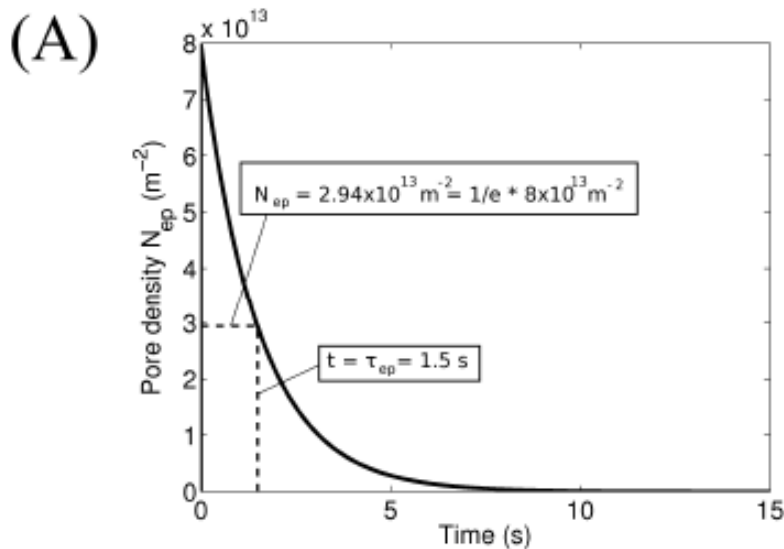
$$A_{Bq} = \frac{m}{m_a} N_A \frac{\ln(2)}{t_{1/2}}$$

Time Scale of Diffusion

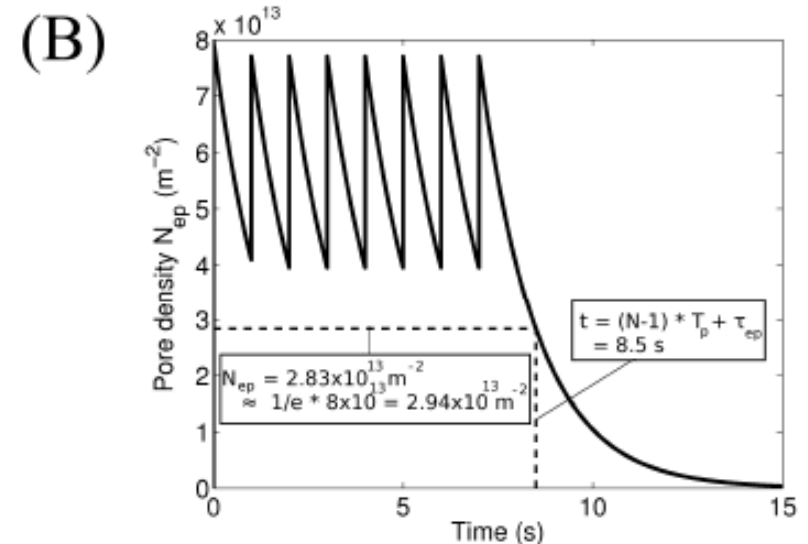
$$\frac{dN_{ep}}{dt} = \alpha e^{(\Phi_m/V_{ep})^2} \left(1 - \frac{N}{N_0 e^{q(\Phi_m/V_{ep})^2}} \right)$$

$$\tau_{ep} = \left. \frac{N_0}{\alpha} e^{(q-1)(\Phi_m/V_{ep})^2} \right|_{\Phi_m=0} = 1.5 \text{ s}$$

Pore Density
Single Pulse



Pore Density
Multiple Pulses



Time and Space Scales in Longitudinal Equation

1D longitudinal Problem for mean potential: Passive membrane

$$\left(1 + \frac{1}{\gamma\mu}\right) C_m \frac{\partial f_m^0}{\partial t} = \frac{\sigma_i a}{2} \frac{\partial^2 f_m^0}{\partial z^2} + \left(1 + \frac{1}{\gamma}\right) \frac{\sigma_i a}{2} \frac{\partial^2 \langle \Psi \rangle}{\partial z^2} - \left(1 + \frac{1}{\gamma\mu}\right) \overline{I_m}$$

Substitute: $\overline{I_m} = \left(\frac{1}{R_m}\right) f_m^0 - \frac{V_{rest}}{R_m}$

$$\underbrace{\{R_m C_m\}}_{\text{Timescale}} \frac{\partial f_m^0}{\partial t} = \underbrace{\left\{ \left(\frac{1}{1 + \frac{1}{\gamma\mu}} \right) \frac{R_m \sigma_i a}{2} \right\}}_{\lambda^2} \frac{\partial^2 f_m^0}{\partial z^2} + \left\{ \left(\frac{\mu(\gamma + 1)}{\gamma\mu + 1} \right) \frac{R_m \sigma_i a}{2} \right\} \frac{\partial^2 \langle \Psi \rangle}{\partial z^2} - f_m^0 - V_{rest}$$

$\text{Timescale} = 10 \text{ ms}$

$\lambda = 1 \text{ mm}$

Lesson: term in parentheses reduces length constant for tightly packed fibers, so that it is reasonable to place no-flux BCs at ends of fiber, which are 4 mm away from respective electrodes for 13 mm fiber and electrodes at +/- 2.5 mm. But for fiber in infinite medium, need longer fiber (24 mm) b/c length constant is $\sqrt{R_m \sigma_i a / 2} = 2.6 \text{ mm}$

Time and Space Scales in Longitudinal Equation

1D longitudinal Problem for mean potential: EPing membrane

$$\left(1 + \frac{1}{\gamma\mu}\right) C_m \frac{\partial f_m^0}{\partial t} = \frac{\sigma_i a}{2} \frac{\partial^2 f_m^0}{\partial z^2} + \left(1 + \frac{1}{\gamma}\right) \frac{\sigma_i a}{2} \frac{\partial^2 \langle \Psi \rangle}{\partial z^2} - \left(1 + \frac{1}{\gamma\mu}\right) \overline{I}_m$$

Substitute: $\overline{I}_m = \left(\frac{R_m + R_p}{R_m R_p}\right) f_m^0 - \frac{V_{rest}}{R_m}$

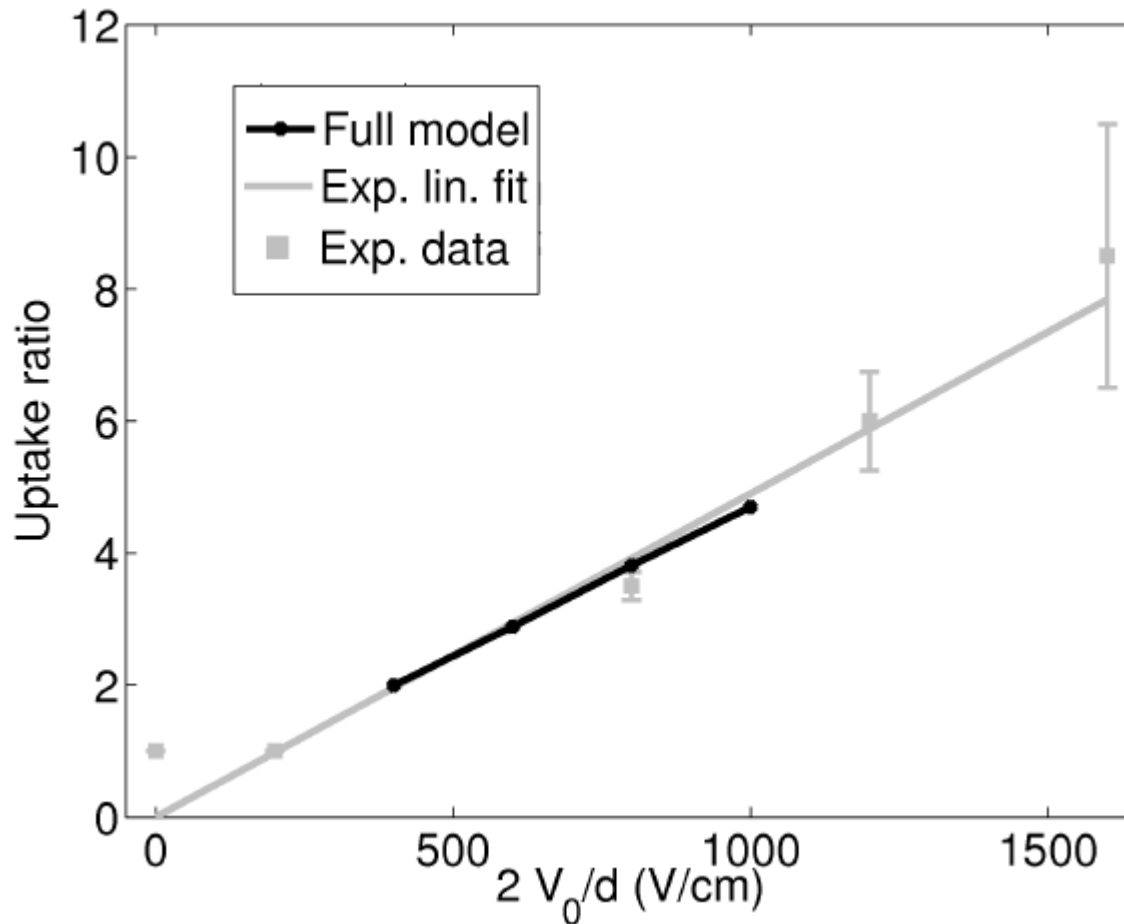
$$\underbrace{\left\{ \frac{R_p R_m C_m}{R_m + R_p} \right\}}_{\text{Timescale}} \frac{\partial f_m^0}{\partial t} = \underbrace{\left\{ \left(\frac{1}{1 + \frac{1}{\gamma\mu}} \right) \frac{R_m R_p}{R_m + R_p} \frac{\sigma_i a}{2} \right\}}_{\lambda^2} \frac{\partial^2 f_m^0}{\partial z^2} + \left\{ \left(\frac{\mu(\gamma + 1)}{\gamma\mu + 1} \right) \frac{R_m R_p}{R_m + R_p} \frac{\sigma_i a}{2} \right\} \frac{\partial^2 \langle \Psi \rangle}{\partial z^2} - f_m^0 - \frac{R_m R_p}{R_m (R_m + R_p)} V_{rest}$$

$$\text{Timescale} = 1 \mu\text{s}$$

$$\lambda = 1 \times 10^{-11} \text{ m}$$

Model Variations

Pulse strength: comparison to experiments



Linear correlation coefficient

- Model: 0.99996
- Experiments: 0.98

Relative difference in uptake
never exceeds 5%

Experimental variability: 6%

Maximum contribution from
longitudinal AF: 10%

Model Variations

Effect of electroporation on tissue conductivities

Simulate two versions solving Ψ numerically

1. Uniform tissue conductivities
2. Nonuniform tissue conductivities – not change σ_{maz} , σ_{may} 2.5x larger still

Find numerical transverse AF $\partial \tilde{\Psi} / \partial \rho$ and longitudinal AF $\partial^2 \langle \Psi \rangle / \partial z^2$

Nonuniform version

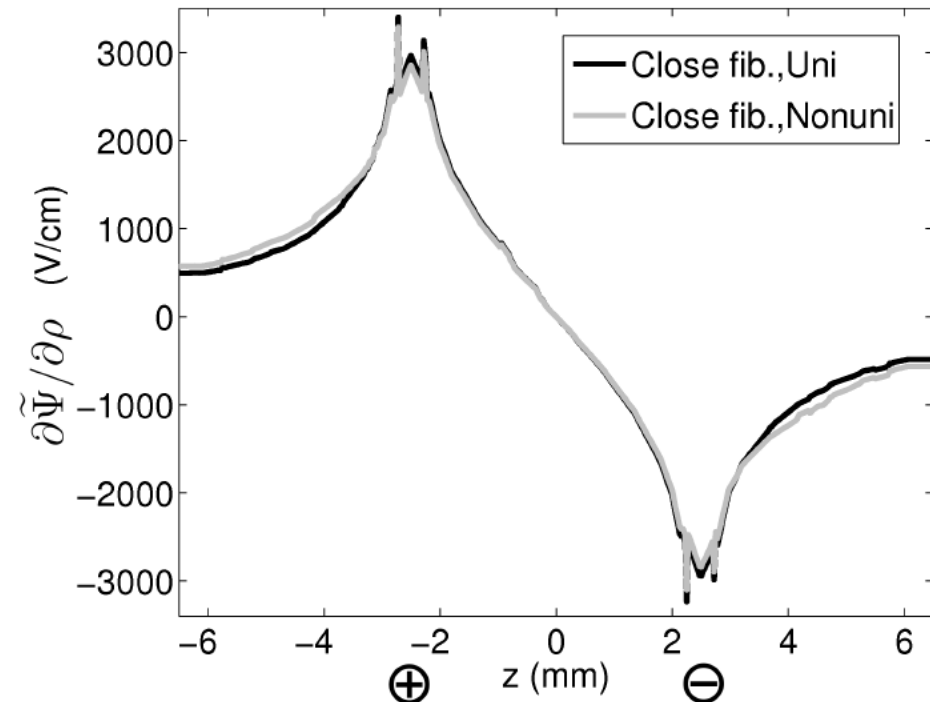
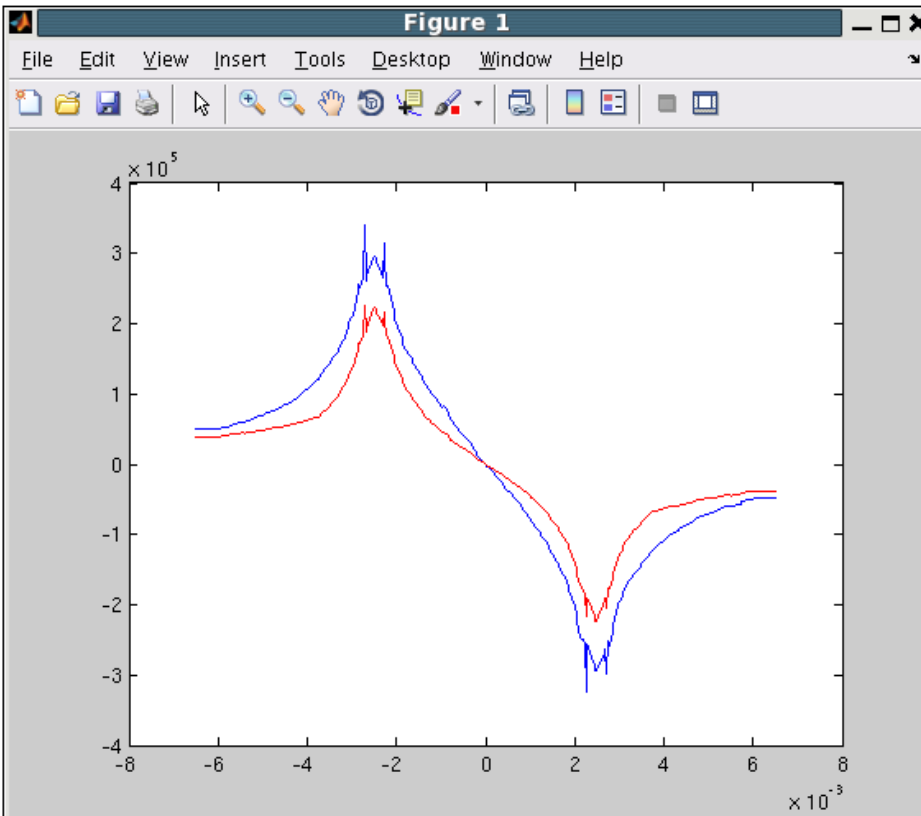
6.7% increase in uptake
from longitudinal AF

Uniform version

3.77% increase in uptake
from longitudinal AF

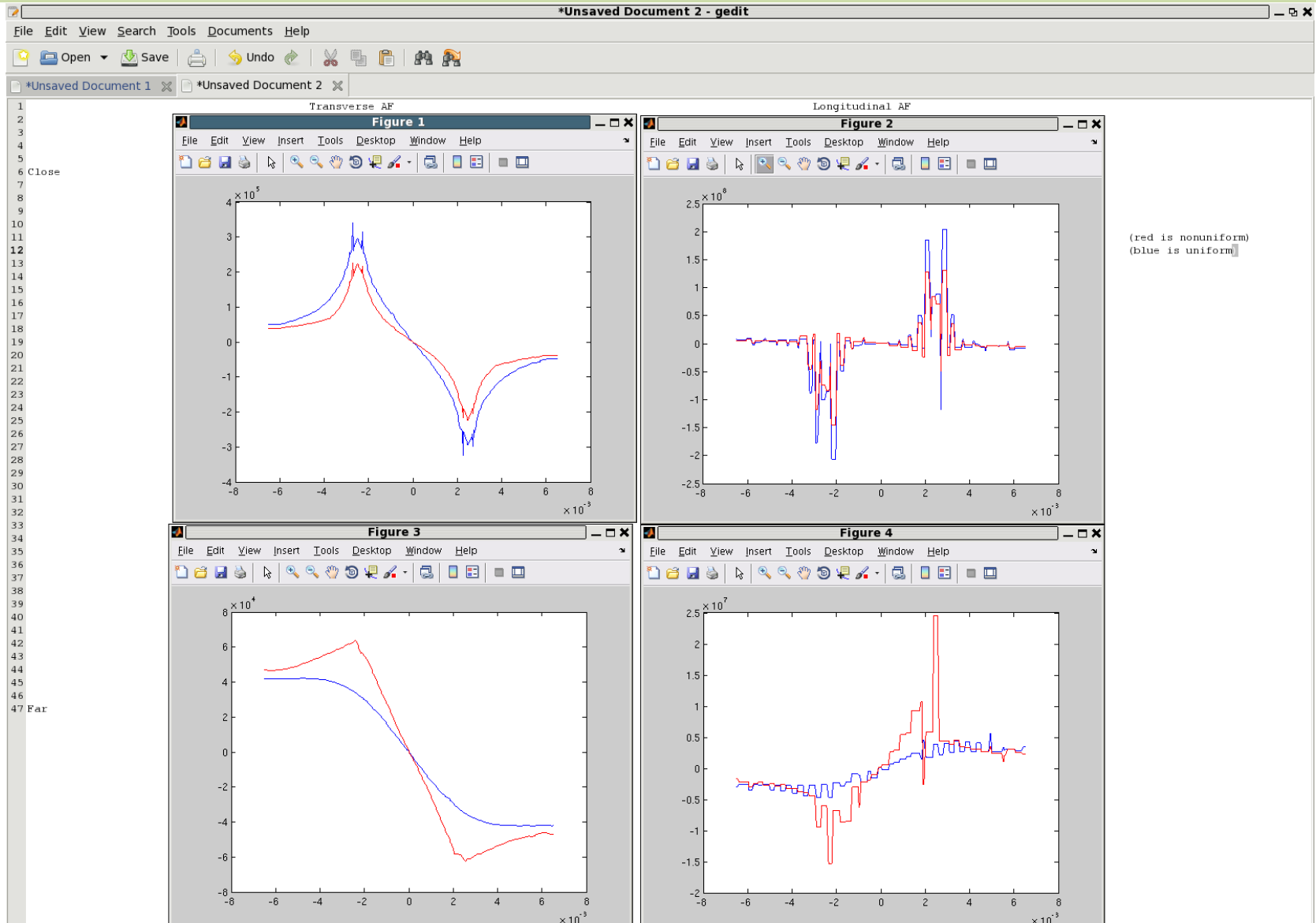
Neglecting effect of EP on tissue conductivities does not change
qualitative role of longitudinal AF much

Model Variations

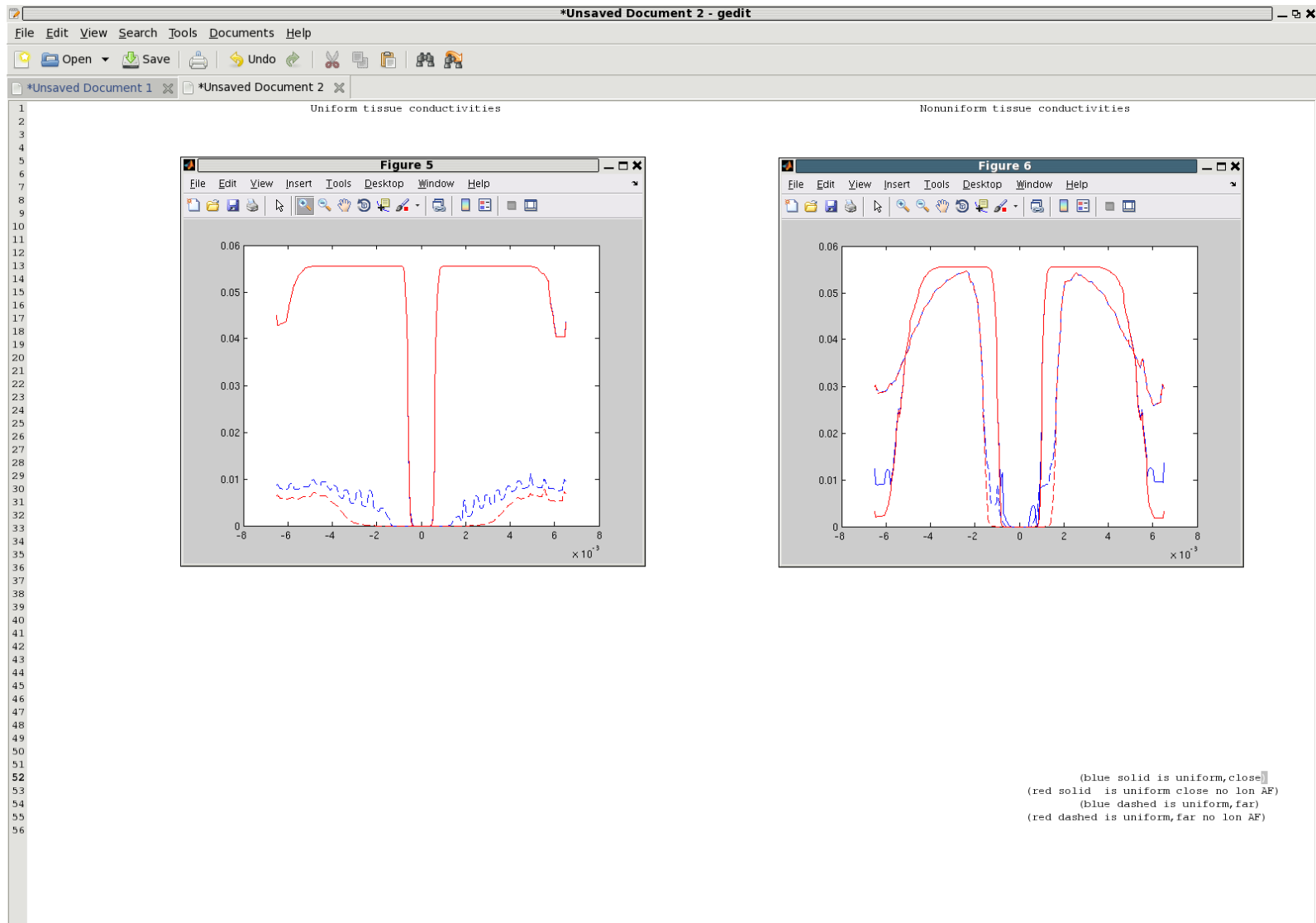


Reason is, for fiber close to electrodes, transverse AF is smaller near ends of fiber, which allows lon. AF to charge and EP membrane more. Whereas when both conductivities increase the trans. AF is larger near ends, which prevents charging via lon. AF

Model Variations



Model Variations



Broader Significance

Current models of neurons

1. Solely longitudinal charging of membrane
 - Not as accurate for close fibers
2. Longitudinal charging, with constant term for transverse charging
 - Not as accurate for close fibers
3. Solve 3D BVP for longitudinal and transverse charging
 - Computationally expensive

Advantages of asymptotic fiber model

- Less expensive
- Accounts for **dynamic** effect of combined transverse and longitudinal charging

Acknowledgements

Thank you

Dr. Wanda K. Neu

Dr. Roger C. Barr

Dr. Thomas J. McIntosh

Dr. Thomas P. Witelski

Dr. Fan Yuan

Other colleagues and mentors

Dr. Margot Bowen

Dr. James Esterline

Dr. Mazella B. Fuller

Dr. Ann Pitruzzello

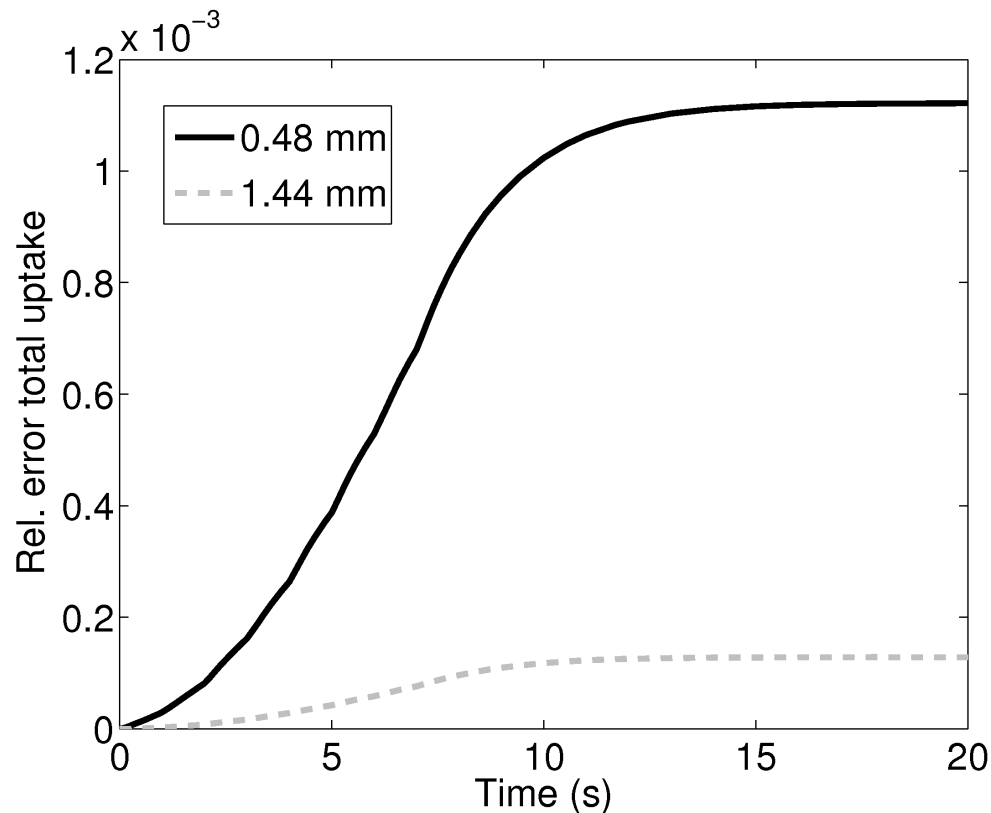
Caroline Ring

Dr. John Wambaugh

Description of Model Components

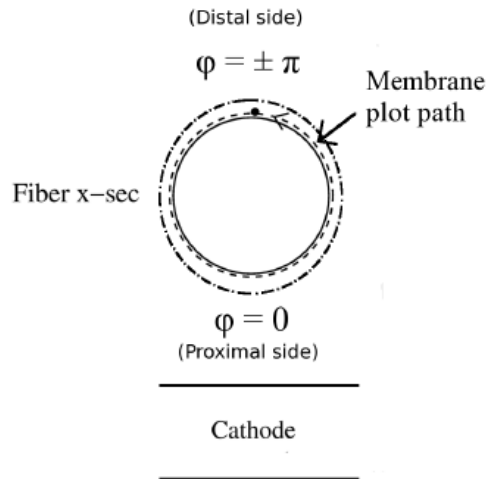
4. Mass transport

- Justification for using series of longitudinally-independent problems



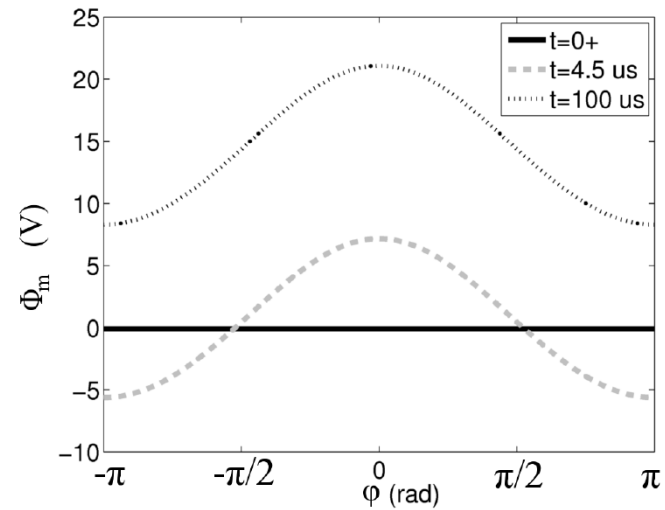
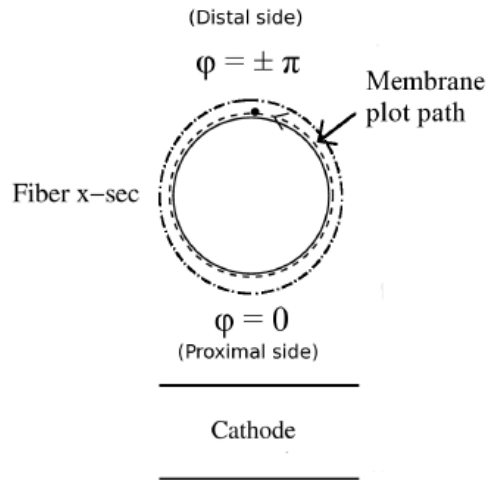
Passive Versus Electroporating

Circumferential Transmembrane Potential



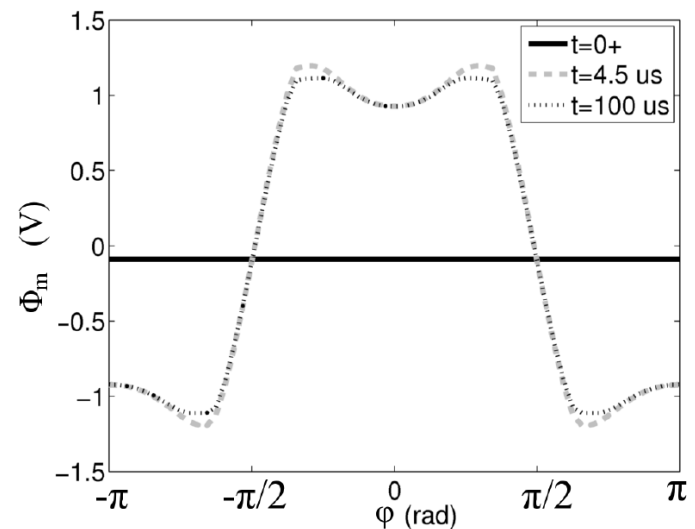
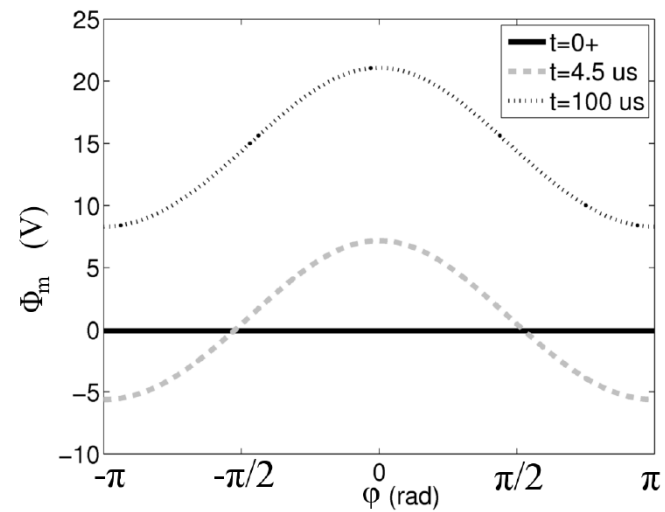
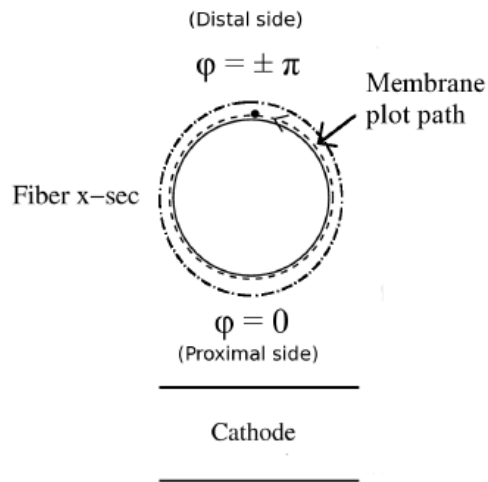
Passive Versus Electroporating

Circumferential Transmembrane Potential



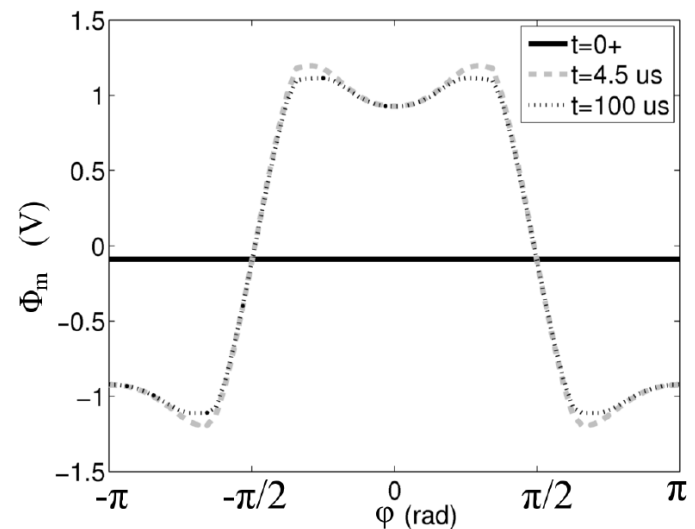
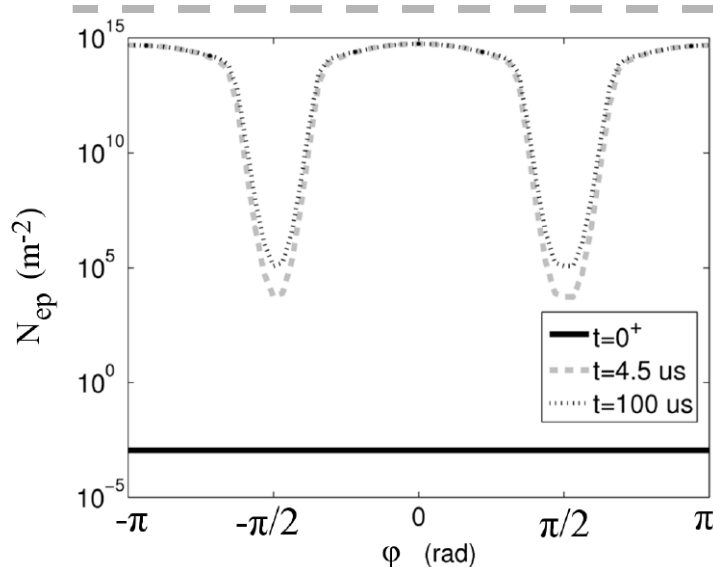
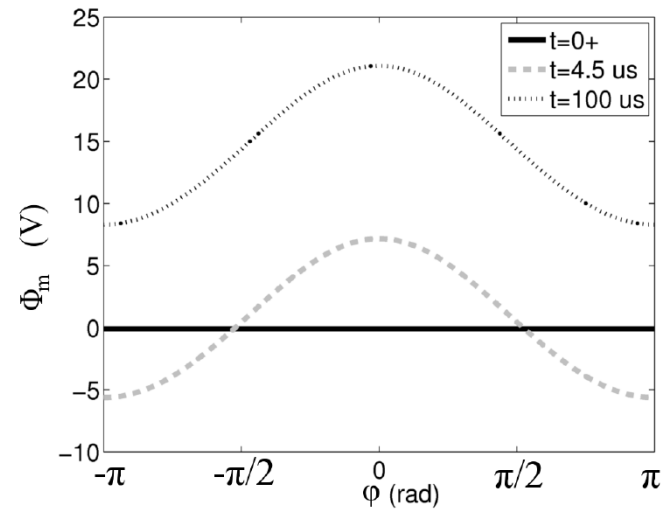
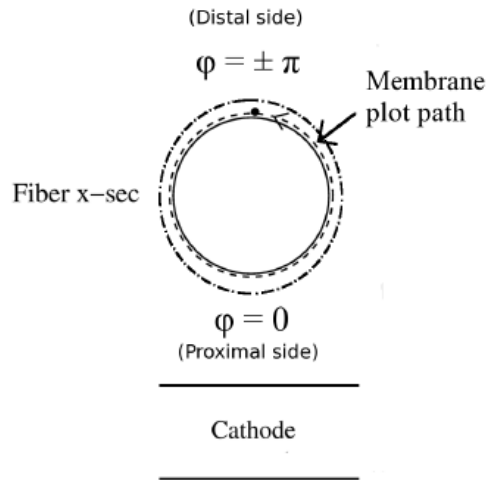
Passive Versus Electroporating

Circumferential Transmembrane Potential



Passive Versus Electroporating

Circumferential Transmembrane Potential



Broader Significance

Difficulties in modeling DNA uptake

- DNA may have two order magnitude small diffusion constant, so may not be able to assume well-mixed in transverse direction
- DNA may form complexes with the membrane before crossing membrane
- May involve endocytosis

Broader Significance

Extension to DNA uptake

- pDNA shown to move micrometers during pulses, pore density varies longitudinally on order of millimeters
- Series of longitudinally independent problems may still hold validity
- N_{lep} along length of fibers determines uptake across membrane
- N_{lep} found from asymptotic fiber model, independent of uptake,
- Uptake across membrane may be similar, but just scaled by slower diffusion coefficient $D \ll D_0$, retaining same dependency on longitudinal charging

$$\partial c_{li} / \partial t \propto D(c_{le} - c_{li}) N_{lep}$$

$$\partial c_{le} / \partial t \propto D(c_{li} - c_{le}) N_{lep}$$

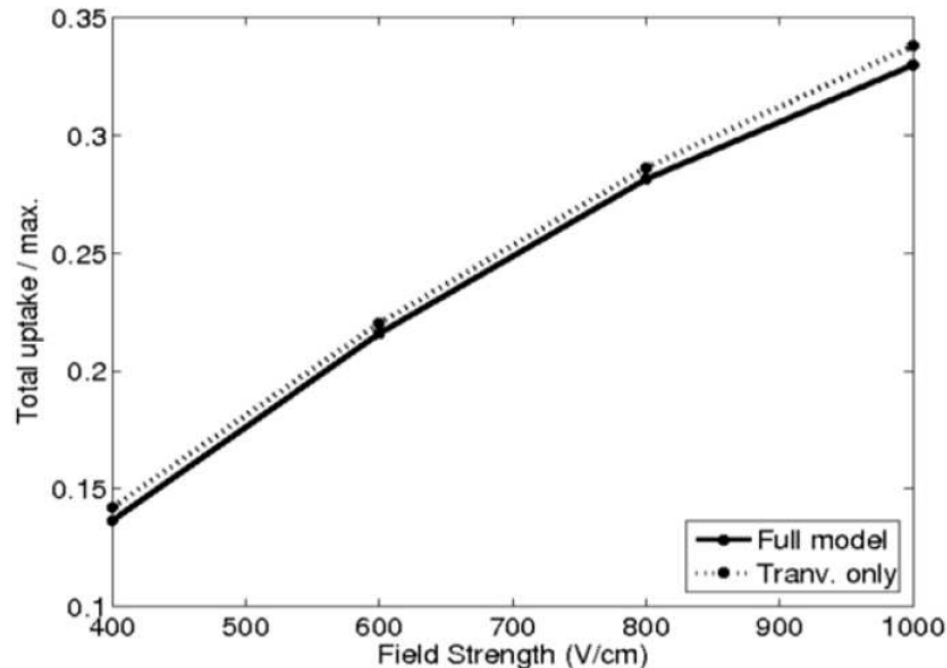
Proposed Additions

I. Error from neglecting longitudinal AF in tissue-wide uptake as vary

- 1) Pulse strength
- 2) Pulse duration

Interesting – no further membrane charging for fibers close to electrodes

- 3) Electrode orientation (longitudinal versus transverse)



Proposed Additions

- II. Compare tissue-wide uptake results in I. to results from experiments and other models of electroporation-mediated molecular uptake

